



*Institute of Fundamental Technological Research
Polish Academy of Sciences*

Synergy of Experiments and Computer Simulations in Research of Turbulent Convection

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Contributions:

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Introduction

- Synergy in scientific research has long been recognized and practiced
- Yet, scientist tend more and more to specialize!
- Rapid developments of computers created a particular gap between experiments and computer simulations
- “Experiments will become obsolete and wind tunnels will be turned into storages of computer outputs”
- “Computer simulations? GiGo!” (“Garbage in, garbage out!”)
- Yet, tremendous advancements in both experimental and simulation/modelling techniques and mutual feedback, synergic inspiration and incentives!

Examples of synergy in three research problems:

1. Thermal convection over horizontal and sloped surfaces in a broad range of conditions including the *extreme* ones
 - Experiments for $Ra=10^8-10^9$
 - DNS for $Ra=10^5-10^8$; LES for $Ra=10^6-10^9$
 - VLES/T-RANS for $Ra=10^6-2 \times 10^{16}$
2. Impinging flows and heat transfer at *higher* Re numbers
 - Single impinging round jets
 - Experiments, RANS and LES, $Re=20.000$
 - Multiple impinging jets
 - Experiments and RANS
 - Single impinging round jet on a cube in cross-flow
 - Experiments, RANS and LES
1. Fluid magnetic dynamo: Hybrid DNS/RANS Computer simulations and interaction with experiments in Riga (Latvia) and Dresden (G)

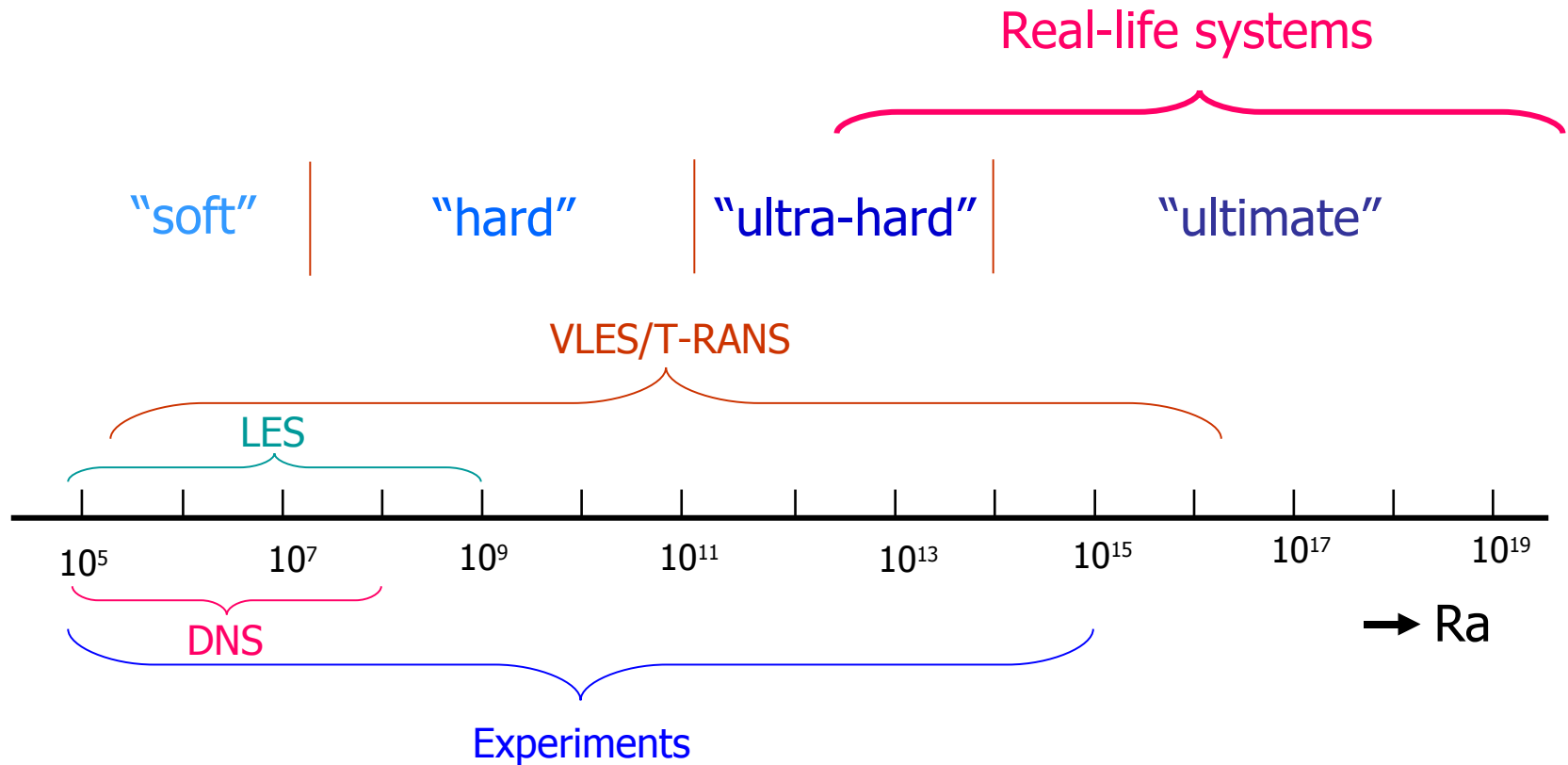
1. Thermal convection

Thermal convection from horizontal surfaces

(Rayleigh-Bénard (R-B) and related problems

- R-B convection = a paradigm of thermal convection; contains most events, structures and features of real-large-scale situations in environmental, geo-, terrestrial and technological systems
- Despite long research, still burdened with controversies:
 - “soft”, “hard”, “ultra-hard” turbulence;
 - $Nu_{\infty} Ra^n$, “n” from 2/7 (1/3?) ($10^7 < Ra < 10^{11}$) to 1/2 when $Ra \rightarrow \infty$
 - scaling of flow properties in various regions and regimes;
 - existence and definition of “wind”, plumes, thermals,..
 - convective-cells and plume structure formation, ordering, ..
 - long-term oscillations, flow reversal, causes-consequences

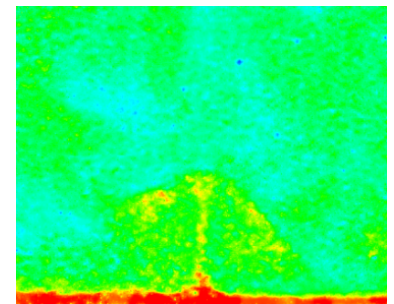
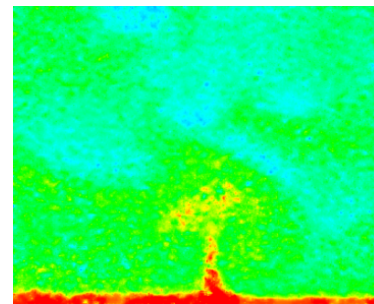
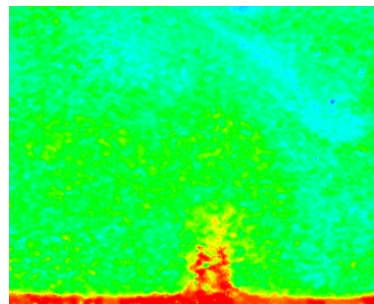
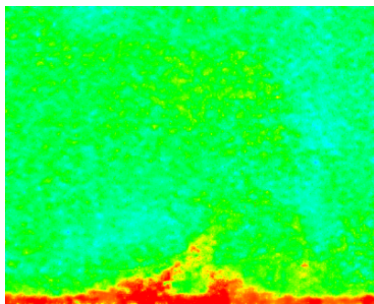
Thermal convection: problems and solutions



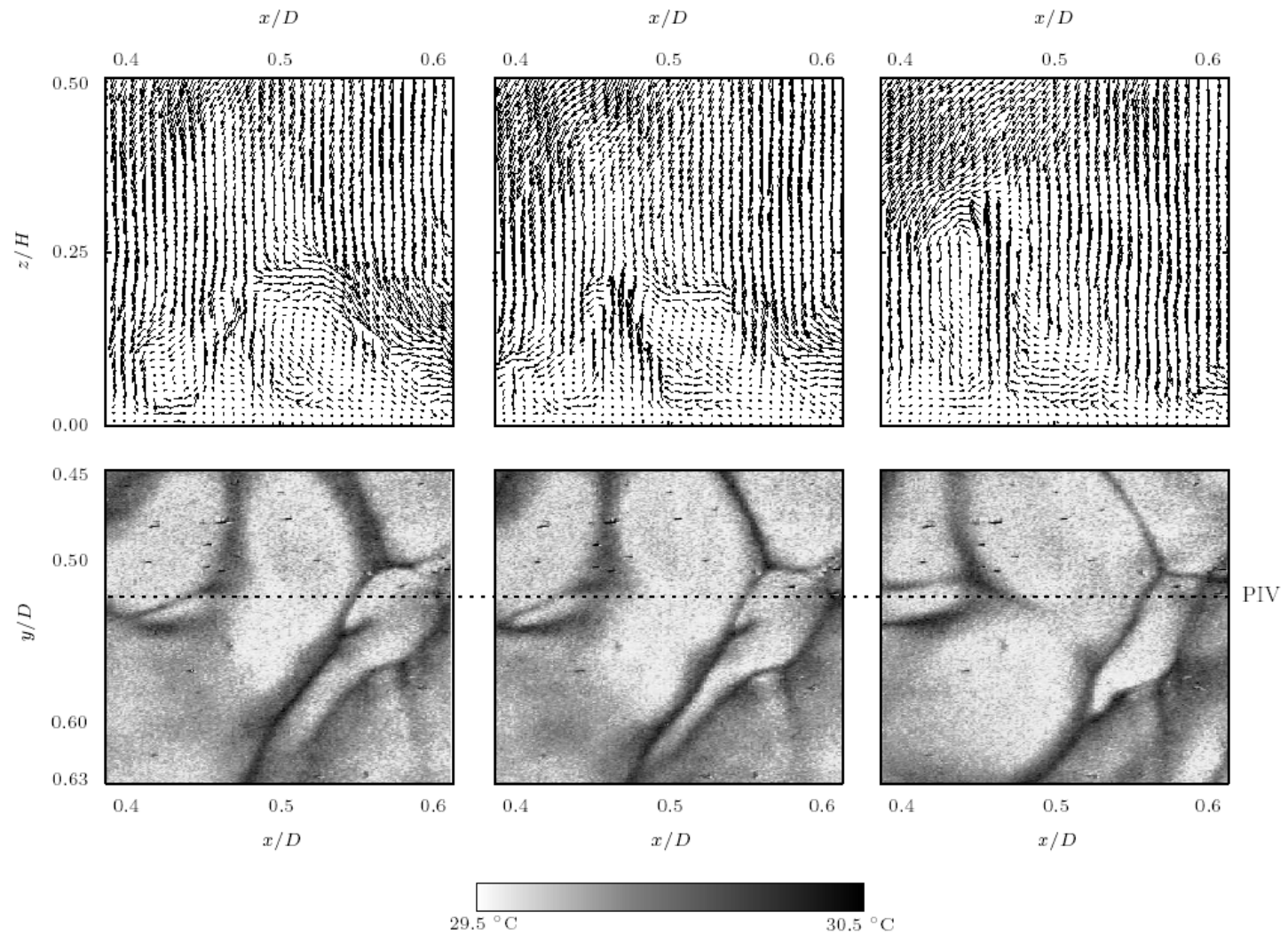
Achievements and limitations in R-B experiments

- Until recently, only point-measurements (especially at high Ra)
- PIV, PTF, LIF, LC brought much advancement, (**almost** all data) but still confined to one-plane, limited domains and single-fields
- 3-D instantaneous field essential for capturing structure and full dynamics: desirable simultaneous application of 3D PIV of PTV (holographic) with suspended LC, thermography and/or spectrometry
- Problems become more challenging with an increase in Ra!

Formation and evolution of thermal plume, suspended Liquid Crystals, $Ra=10^8$, $Pr=7.0$, 4:4:1 domain, (*Verdoold, Tummers, Hanjalic et al. 2004*)



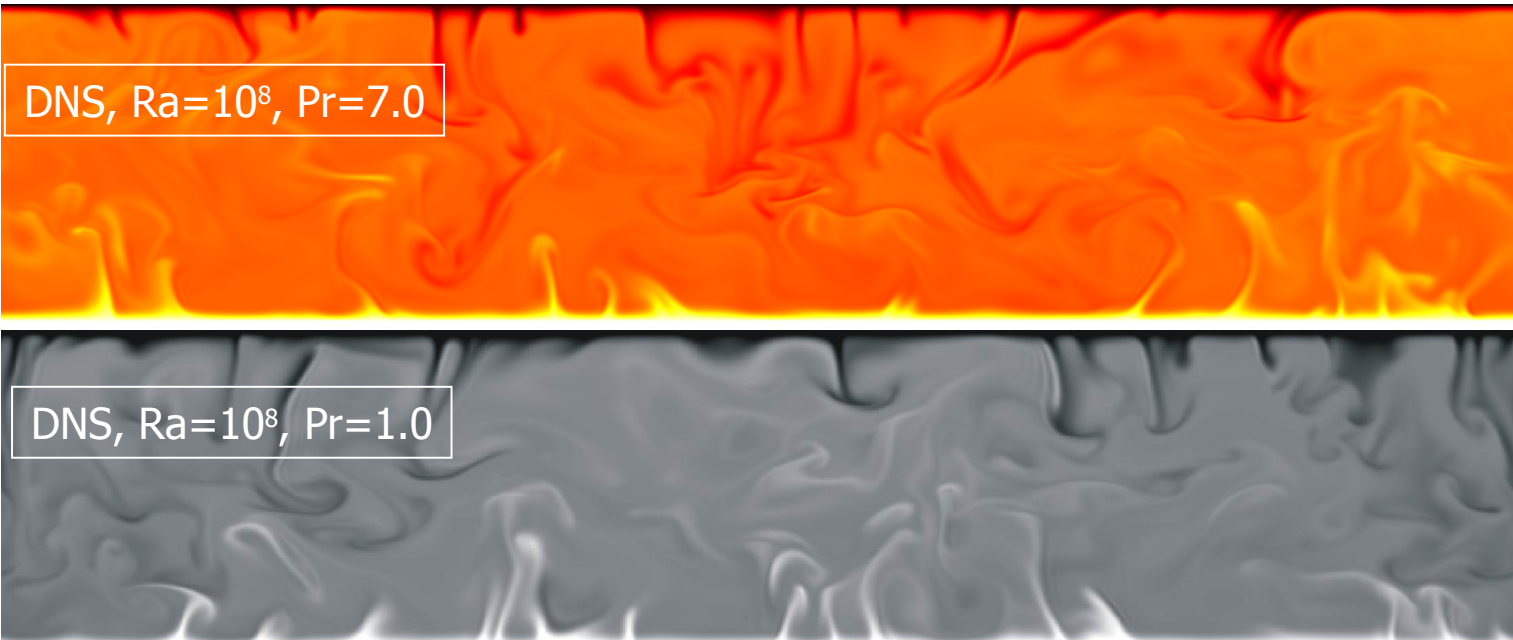
Synchronised snapshot of PIV in x-y and LCT in a near-wall x-z plane in R-B convection at 10 sec intervals ($Ra=1.3 \times 10^8$, $Pr=7.0$) (*Verdool et al. 2004*)



Achievements and limitations in DNS of R-B convection

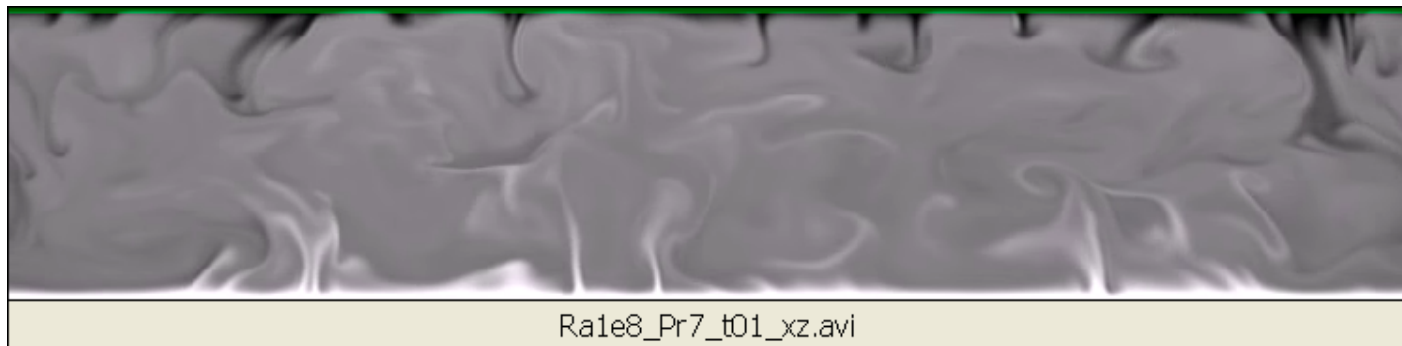
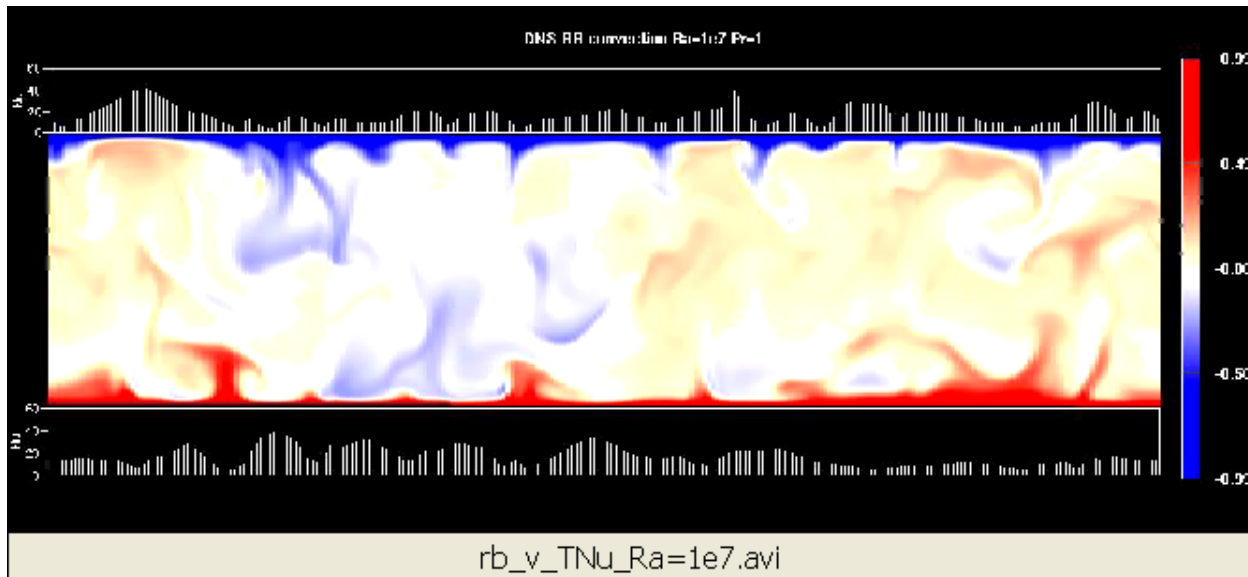
- Computer simulations (DNS, LES): tremendous potential, make it possible to collect **all** information needed, (some are still inaccessible to experiments) , but only for low Ra's!

Recent: $Ra=1.1 \times 10^8$, $Pr=7$, grid: $768 \times 768 \times 320$ (~ 188 million!) on 192 processor of TERAS, (~ 22 hours per processor for **one turn-over time, ~ 55 sec real time**) (*Van Reeuwijk, Jonker, Hanjalic, 2005*)



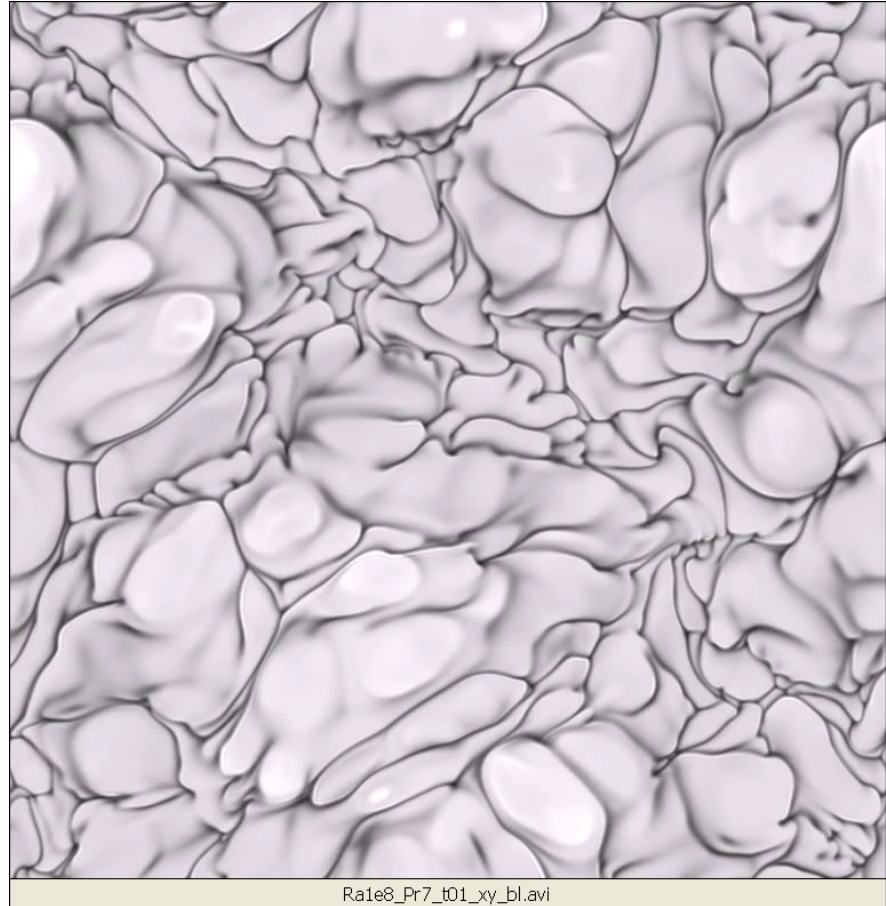
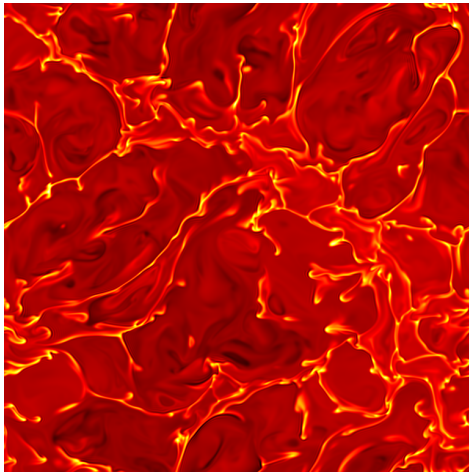
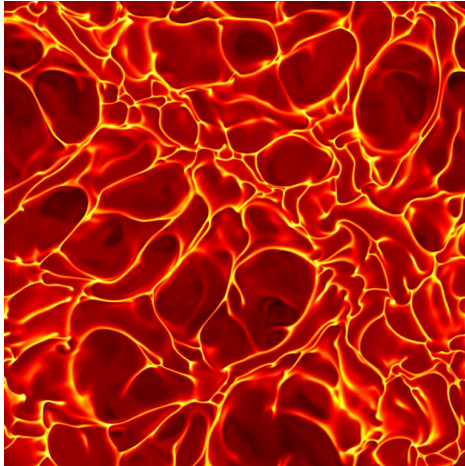
Recent achievements in DNS of R-B convection

$Ra=1.1 \times 10^8$, $Pr=7$, grid: $768 \times 768 \times 320$, Finite volume + spectral integration, grid clustered in near-wall regions, (*Van Reeuwijk, Jonker, Hanjalic, 2005*)



Recent achievements in DNS of R-B convection

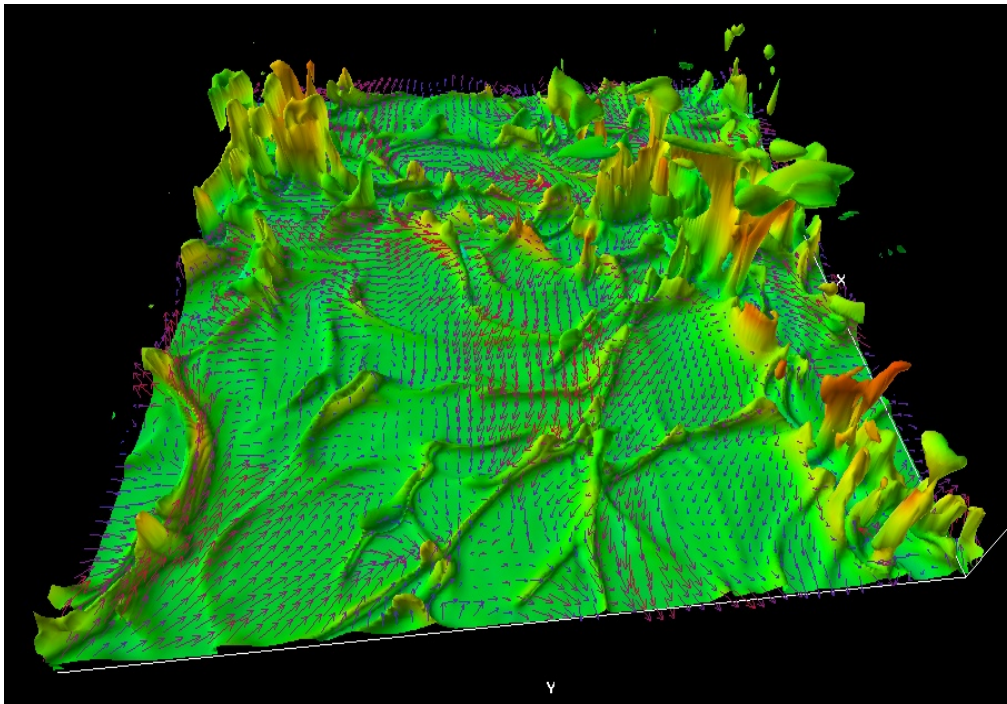
$Ra=1.1 \times 10^8$, $Pr=7$, grid: $768 \times 768 \times 320$, Finite volume + spectral integration, grid clustered in near-wall regions, (*Van Reeuwijk, Jonker, Hanjalic, 2005*)



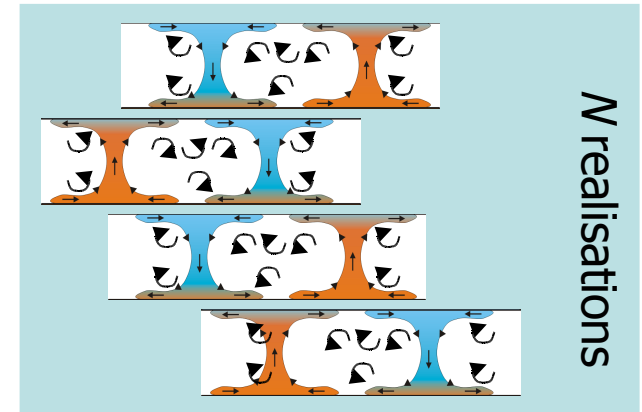
Some outcome of DNS: Identification of "Wind"

Triple decomposition:

$$\text{Flow} = \underbrace{\text{'Wind' + Plumes}}_{\text{Deterministic}} + \underbrace{\text{Fluctuations}}_{\text{Stochastic}}$$



Classic ensemble averaging:



$$\bar{\mathbf{u}}(\vec{\mathbf{x}}) = \frac{1}{N} \sum_N \mathbf{u}_n(\vec{\mathbf{x}} - \vec{\mathbf{d}}_n)$$

$$\boxed{\bar{u}_i = 0} = \left[\begin{array}{c} \text{[Diagram 1]} + \text{[Diagram 2]} \end{array} \right] +$$

$$\downarrow \text{g}$$

$$\left[\begin{array}{c} \text{[Diagram 3]} + \text{[Diagram 4]} \end{array} \right] +$$

$$\dots$$

'Wind': $\bar{u} = \bar{v} = \bar{w} = 0$

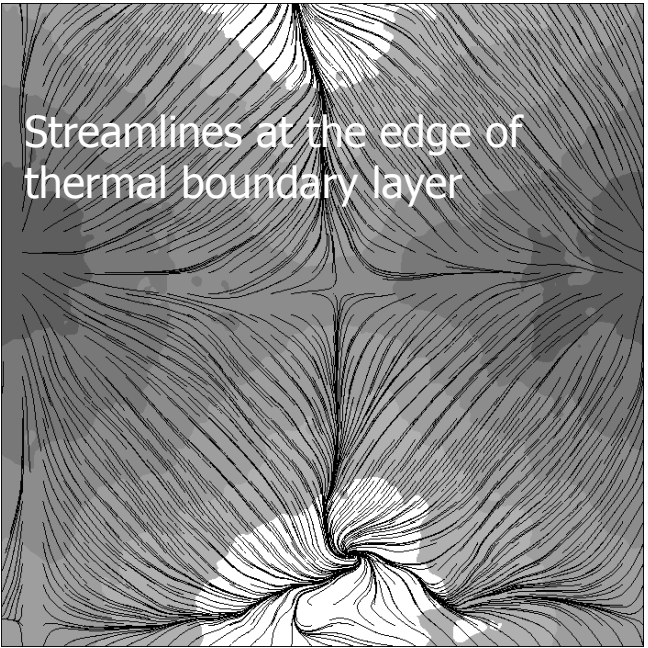
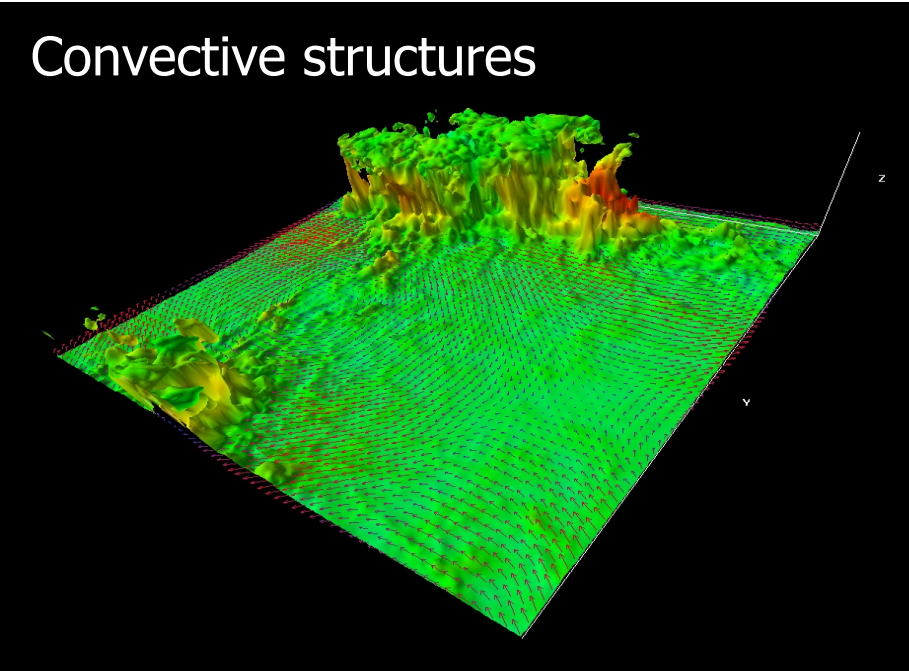
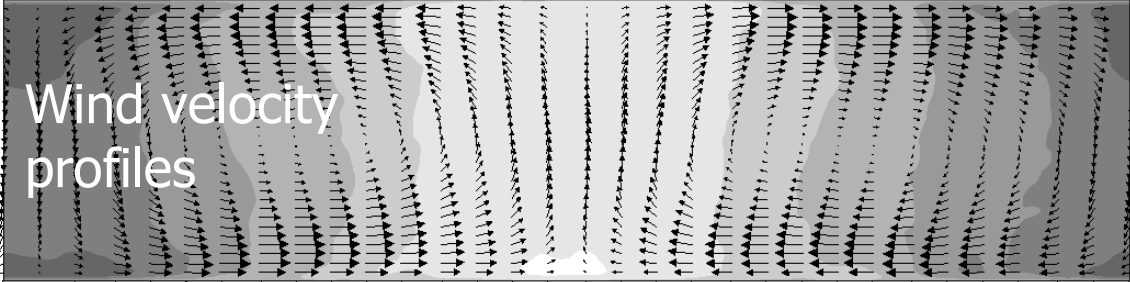
Symmetry-accounting ('conditional') ensemble averaging

Using DNS data

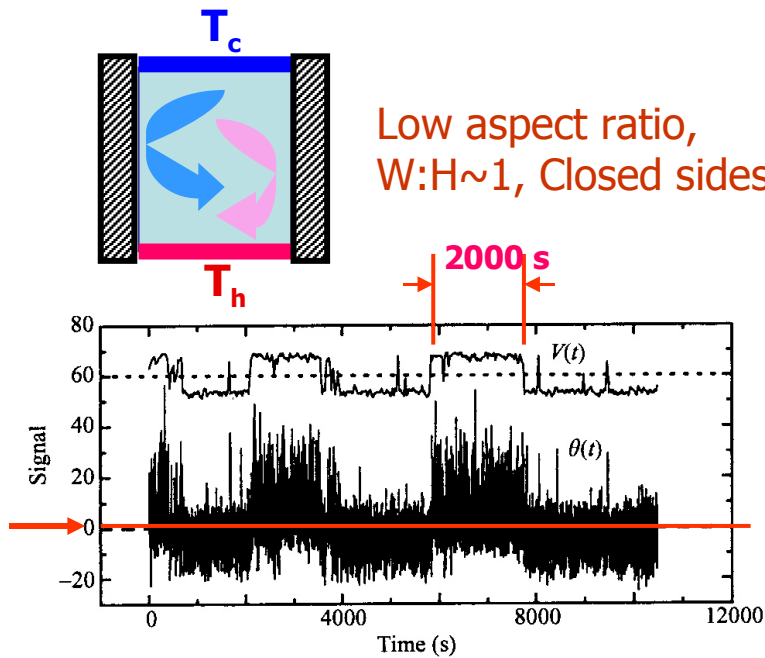
Van Reeuwijk, Jonker, Hanjalic, Phys.Fluids, 2005

$$\bar{u}(\vec{x}) = \frac{1}{N} \sum_N u_n(\vec{x} - \vec{d}_n) \mathcal{S}^{(\alpha)}$$

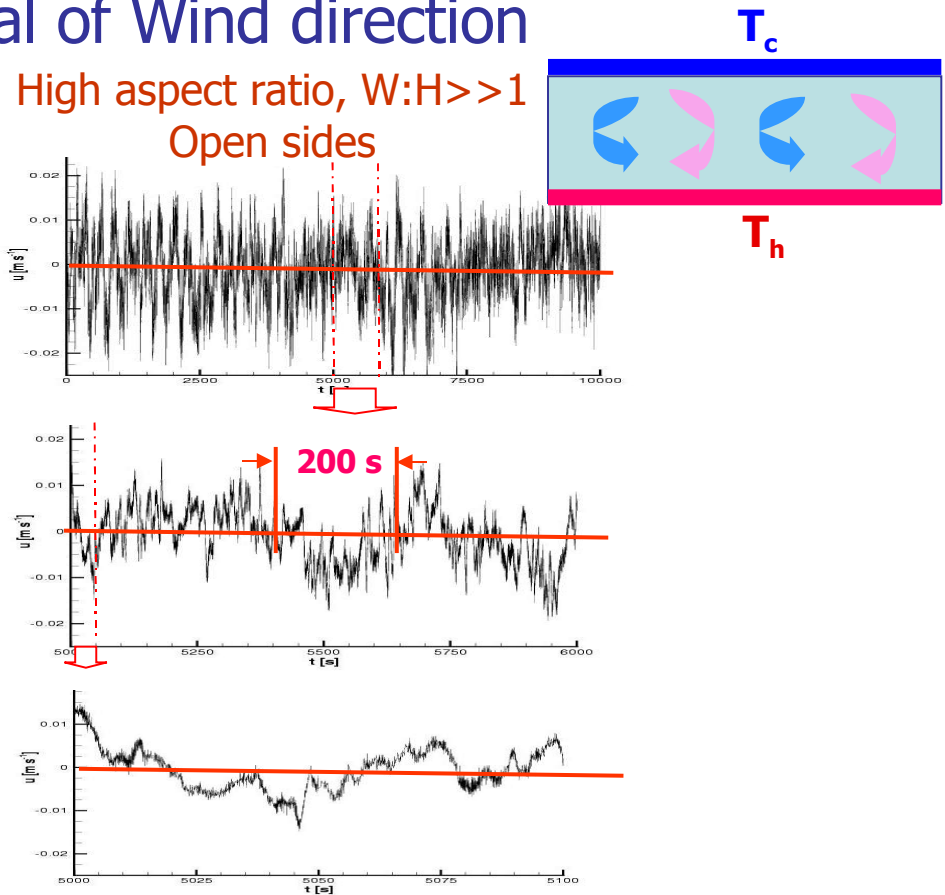
$\mathcal{S}^{(\alpha)}$ = symmetry operator



Some controversies related to long-term oscillations and sudden reversal of Wind direction



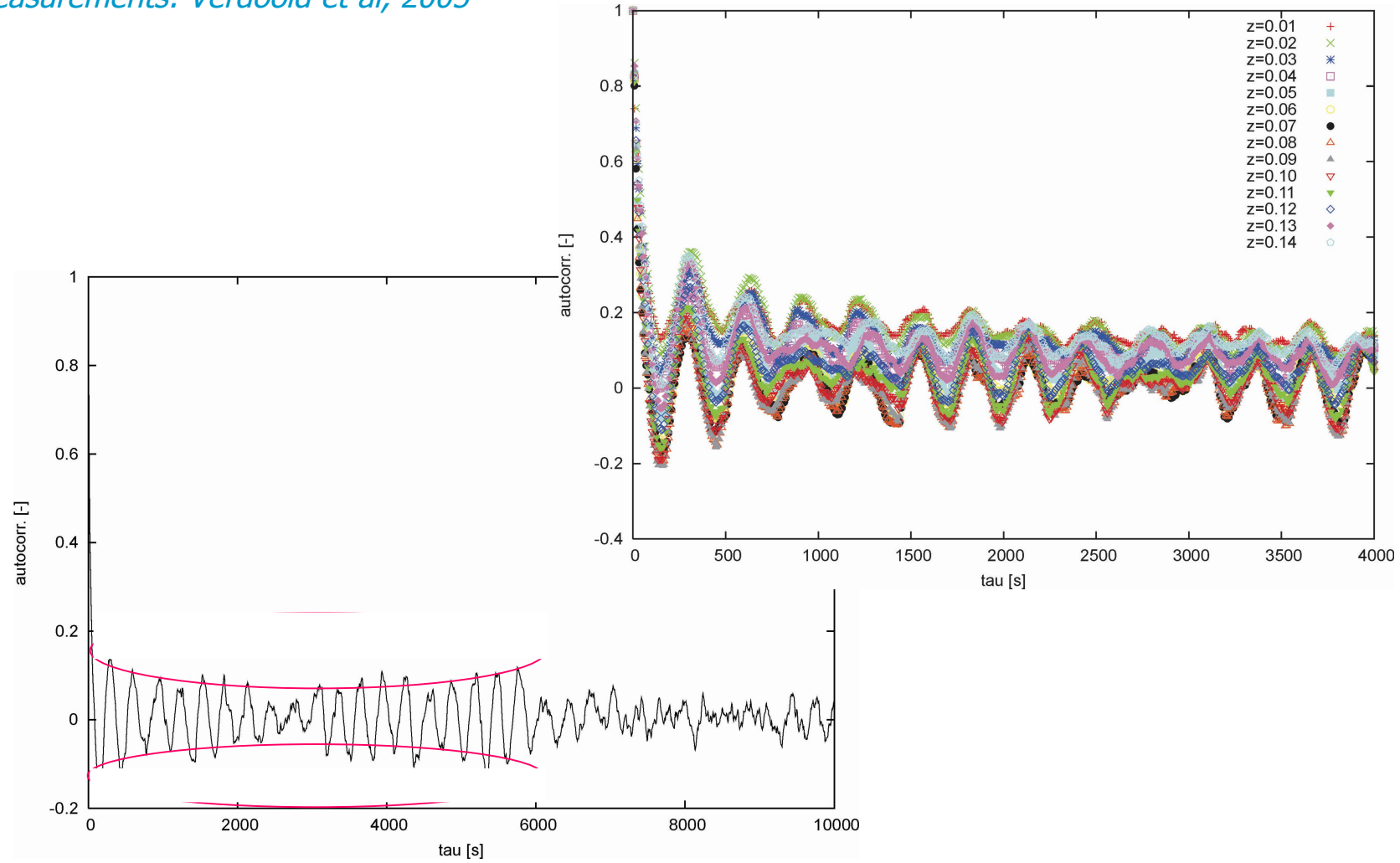
Wind velocity (smoothed) and temperature recordings in the experiment with cryogenic helium in a cylindrical enclosure 50×50 cm at $Ra = 1.5 \times 10^{11}$ (Niemela et al. 2001)



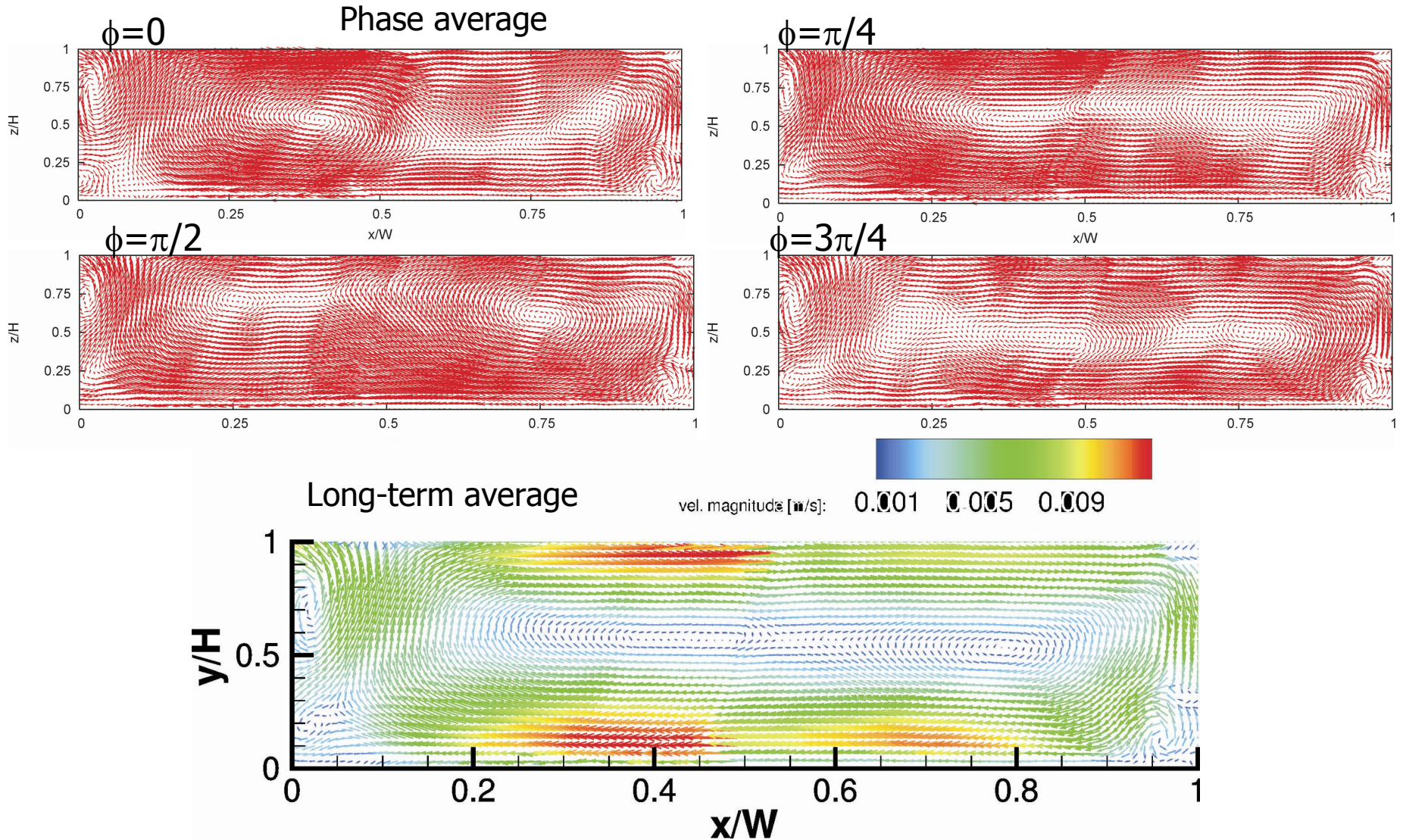
LDA long-term velocity recording at 33 mm from the bottom wall in the centre of a $60 \times 60 \times 15,5$ cm R-B experiment with water, $Ra \approx 10^9$ (Verdoold et al. 2004/05)

Velocity autocorrelation in R-B convection

Measurements: Verdoold et al, 2005

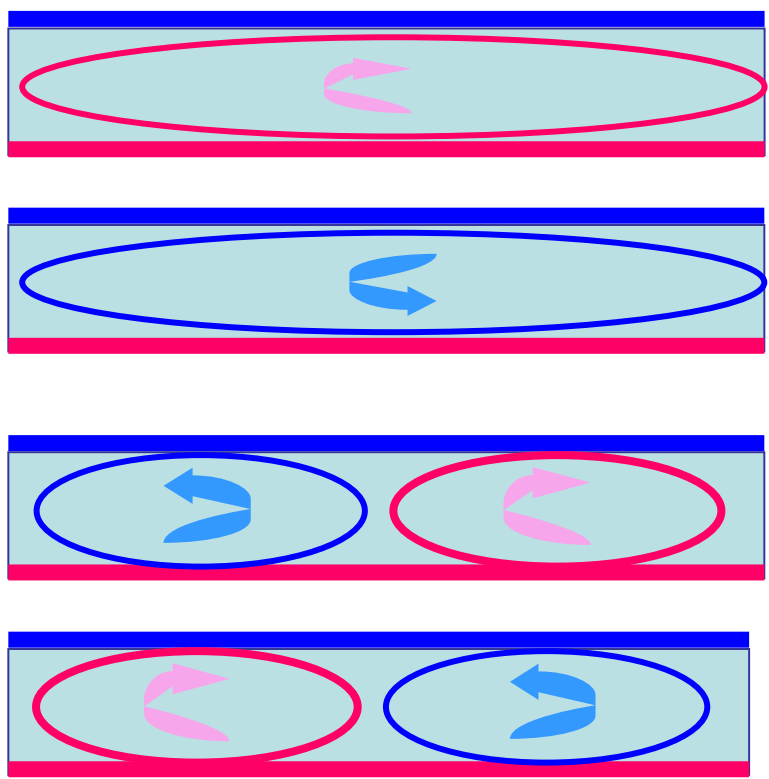


Convective patterns in R-B convection

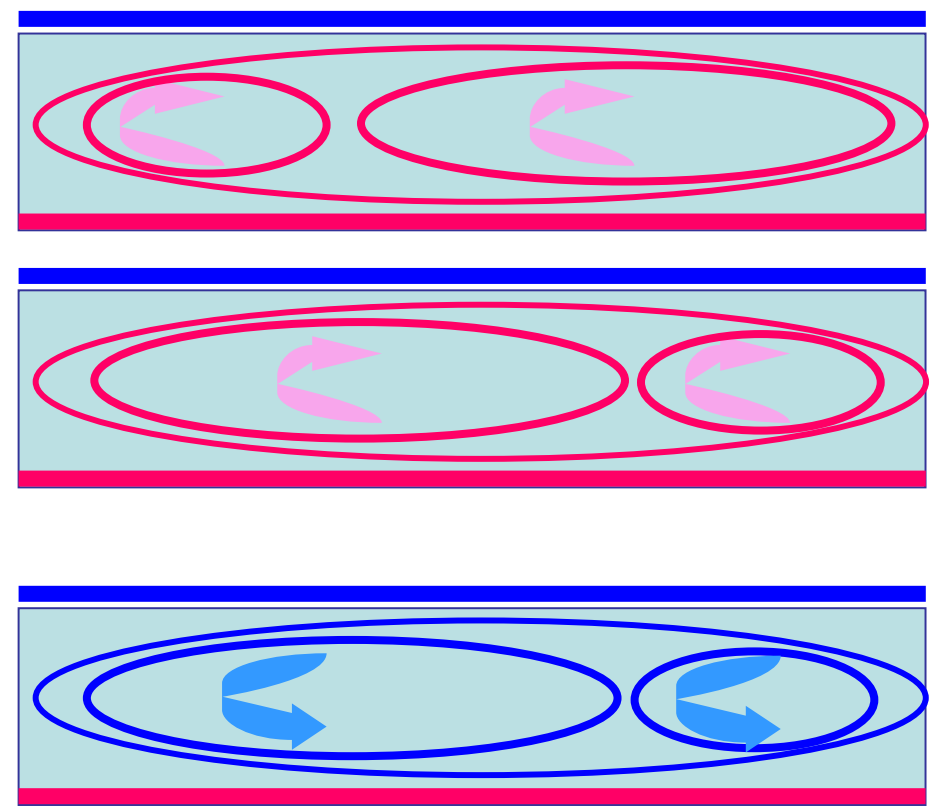


Convective patterns in R-B convection

Classic



Possible long-term periodic scenario

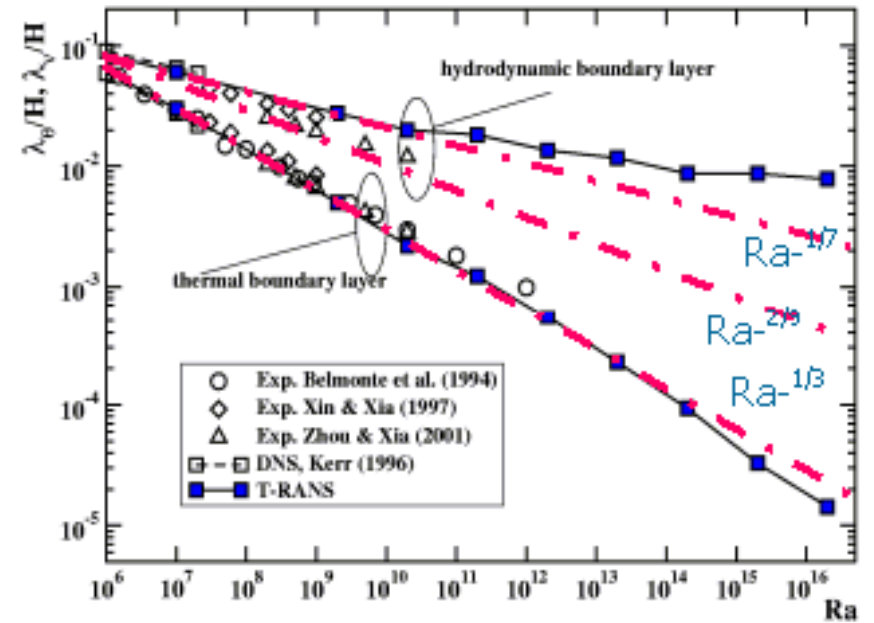
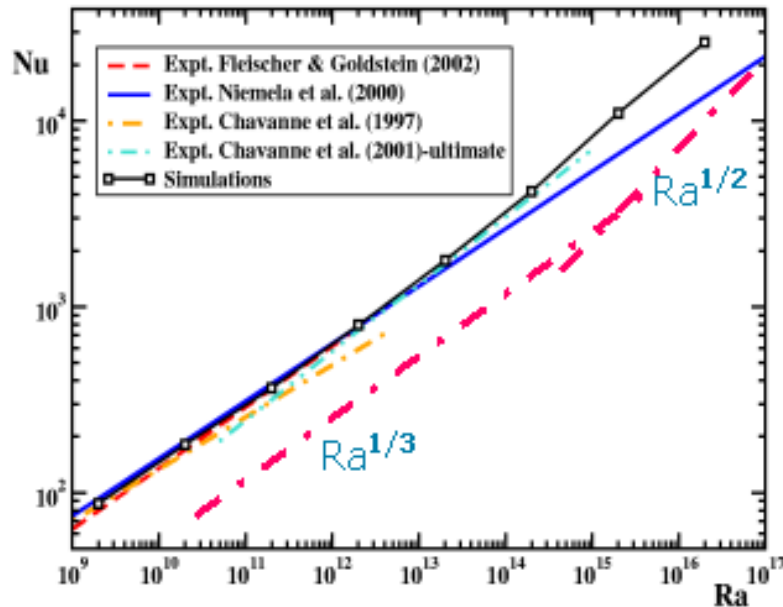


High Ra challenges in R-B convection: Transient RANS

(Kenjeres & Hanjalic, Phys Rev. E, 2002)

➤ $Nu \propto Ra^{1/3}$ for $Ra < 10^{12}$ ($Pr \ll 1$)

➤ $Nu \propto Ra^{1/2}$ for $Ra \rightarrow \infty$



➤ $\lambda_v/H \propto Ra^{-1/7}$

➤ $\lambda_\theta/H \propto Ra^{-1/3}$

A priori test of different models in generic flows

The AFM of Kenjeres *et al* (2004) in conjunction with the $v_2-f-\theta^2$ model reproduces best the heat flux components in both generic cases of natural convection: vertical and horizontal plane channels with $\nabla T \perp \vec{g}$ and $\nabla T \parallel \vec{g}$ respectively.

SGDH

$$\overline{\theta u_i} = -C_\theta \tau k \frac{\partial T}{\partial x_i}$$

GGDH

$$\overline{\theta u_i} = -C_\theta \tau \overline{u_i u_j} \frac{\partial T}{\partial x_j}$$

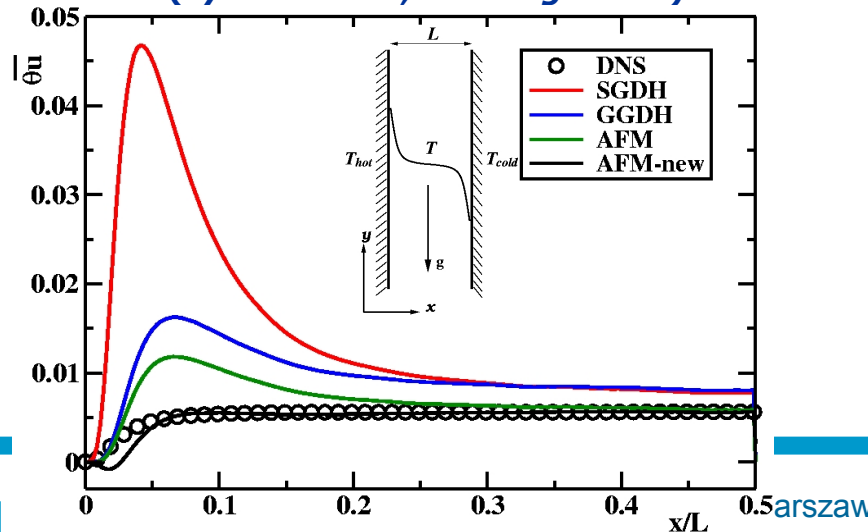
AFM

$$\overline{\theta u_i} = -C_\theta \tau \left[\overline{u_i u_j} \frac{\partial T}{\partial x_j} + \overline{\theta u_j} \frac{\partial U_i}{\partial x_j} + \beta g_i \overline{\theta^2} \right]$$

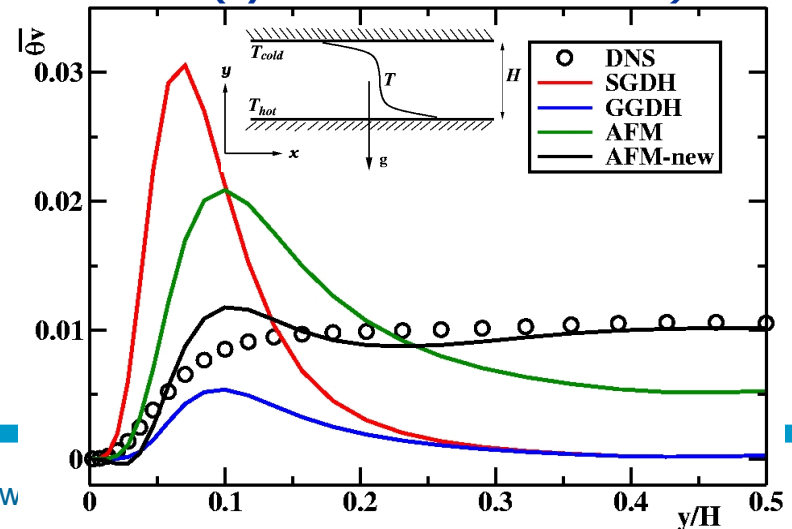
AFM-new

$$\overline{\theta u_i} = -C_\theta \tau \left[\overline{u_i u_j} \frac{\partial T}{\partial x_j} + \overline{\theta u_j} \frac{\partial U_i}{\partial x_j} + \beta g_i \overline{\theta^2} \right] + 1.5 a_{ij} \overline{\theta u_j} \quad \tau = \max \left[\frac{k}{\varepsilon}, C_\mu \left(\frac{\nu}{\varepsilon} \right)^{1/2} \right]; \quad C_\theta = 0.3$$

Wall normal heat flux in a side heated vertical channel; $Ra=5 \times 10^6$, $Pr=0.71$ (Symbols: DNS, Versteegh 1998)



Wall normal heat flux in a heated-from-below horizontal channel (R-B-convection) $Ra=6.3 \times 10^5$, $Pr=0.71$ (Symbols: DNS Woerner 1994)



T-RANS Equations and subscale models:

$$\frac{\partial \langle U_i \rangle}{\partial t} + \langle U_j \rangle \frac{\partial \langle U_i \rangle}{\partial x_j} = \frac{\partial}{\partial x_j} \left(\nu \frac{\partial \langle U_i \rangle}{\partial x_j} - \tau_{ij} \right) - \frac{1}{\rho} \frac{\partial (\langle P \rangle - P_{\text{ref}})}{\partial x_i} + \beta g_i (\langle T \rangle - T_{\text{ref}})$$

$$\frac{\partial \langle T \rangle}{\partial t} + \langle U_j \rangle \frac{\partial \langle T \rangle}{\partial x_j} = \frac{\partial}{\partial x_j} \left(\frac{\nu}{Pr} \frac{\partial \langle T \rangle}{\partial x_j} - \tau_{\theta j} \right)$$

$$\frac{\partial \langle C \rangle}{\partial t} + \langle U_j \rangle \frac{\partial \langle C \rangle}{\partial x_j} = \frac{\partial}{\partial x_j} \left(\frac{\nu}{Sc} \frac{\partial \langle C \rangle}{\partial x_j} - \tau_{cj} \right)$$

assuming weak equilibrium
 $(D/Dt - Diff) \overline{\varphi u_i} = 0$

Subscale ASM/AFM/ACM

+final closure: 3eqn. model

$$\frac{D \langle k \rangle}{Dt} = D_k + P_k + G_k - \langle \varepsilon \rangle$$

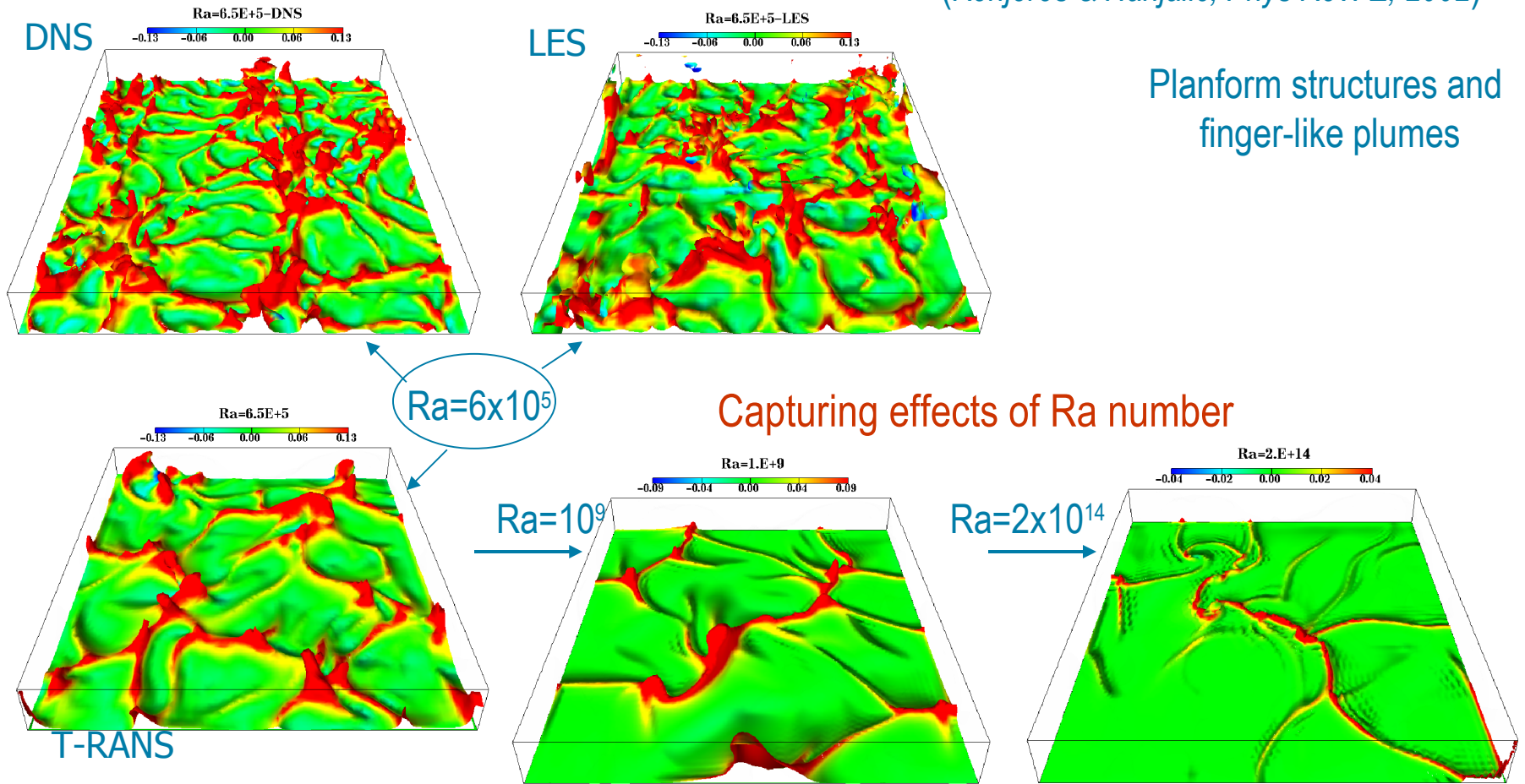
$$\frac{D \langle \varepsilon \rangle}{Dt} = D_\varepsilon + P_{\varepsilon 1} + P_{\varepsilon 2} + G_\varepsilon - Y$$

$$\frac{D \langle \theta^2 \rangle}{Dt} = D_\theta + P_\theta - \langle \varepsilon_\theta \rangle$$

$$\begin{aligned} \tau_{\theta i} &= -C_\phi \frac{\langle k \rangle}{\langle \varepsilon \rangle} \left[\tau_{ij} \frac{\partial \langle T \rangle}{\partial x_j} + \xi \tau_{\theta j} \frac{\partial \langle U_i \rangle}{\partial x_j} + \eta \beta g_i \langle \theta^2 \rangle \right] \\ \tau_{ci} &= -C_\phi \frac{\langle k \rangle}{\langle \varepsilon \rangle} \left[\tau_{ij} \frac{\partial \langle C \rangle}{\partial x_j} + \xi \tau_{cj} \frac{\partial \langle U_i \rangle}{\partial x_j} \right] \\ \tau_{ij} &= -\nu_t \left(\frac{\partial \langle U_i \rangle}{\partial x_j} + \frac{\partial \langle U_j \rangle}{\partial x_i} \right) + \frac{2}{3} \langle k \rangle \delta_{ij} + C \frac{k}{\varepsilon} \beta g_i \overline{\theta u_j} \end{aligned}$$

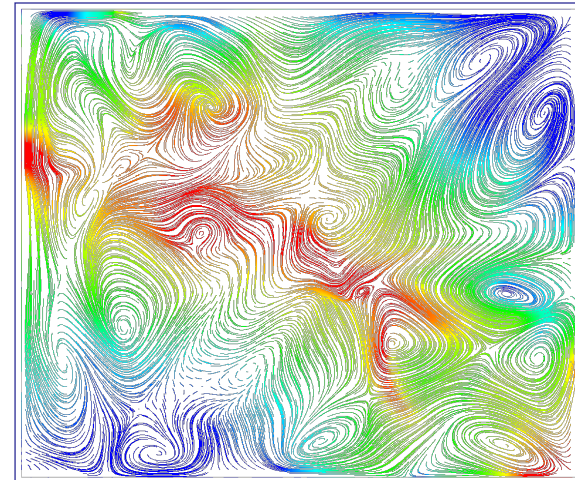
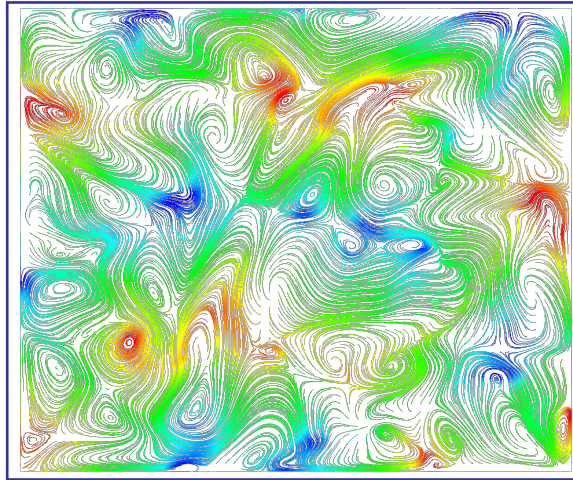
Meeting the high Ra challenges: Transient RANS

Comparison of DNS, LES and T-RANS for $Ra=6 \times 10^5$ and T-RANS “extrapolation”
(Kenjeres & Hanjalic, *Phys Rev. E*, 2002)

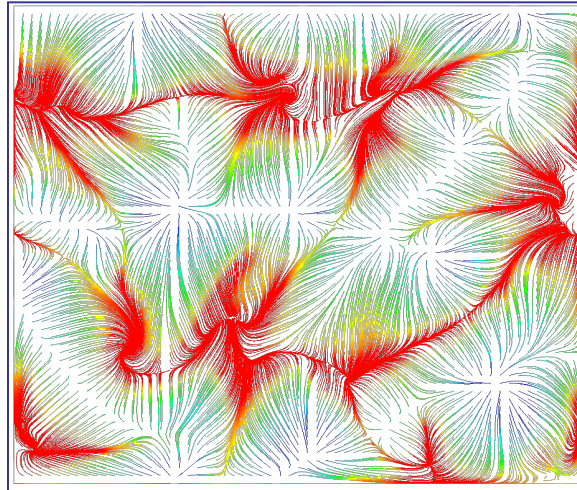


T-RANS of R-B: Temperature colored instantaneous trajectories

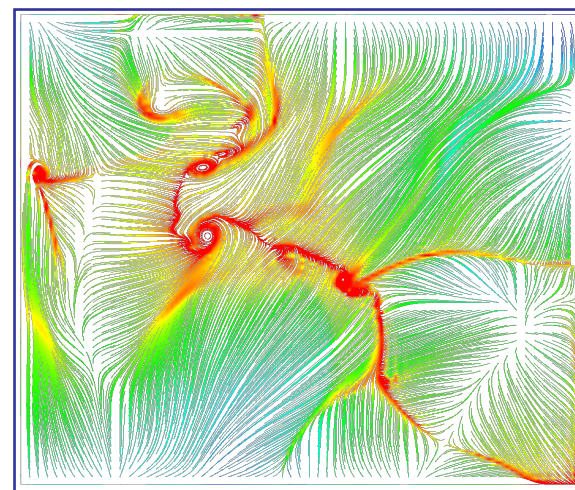
central horizontal
plane ($z/D=0.5$)



inside thermal
boundary layer
($z/D=0.05$):



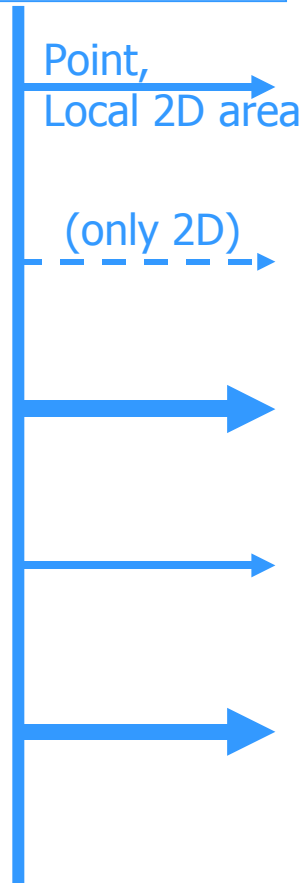
$Ra=6.5 \times 10^5$



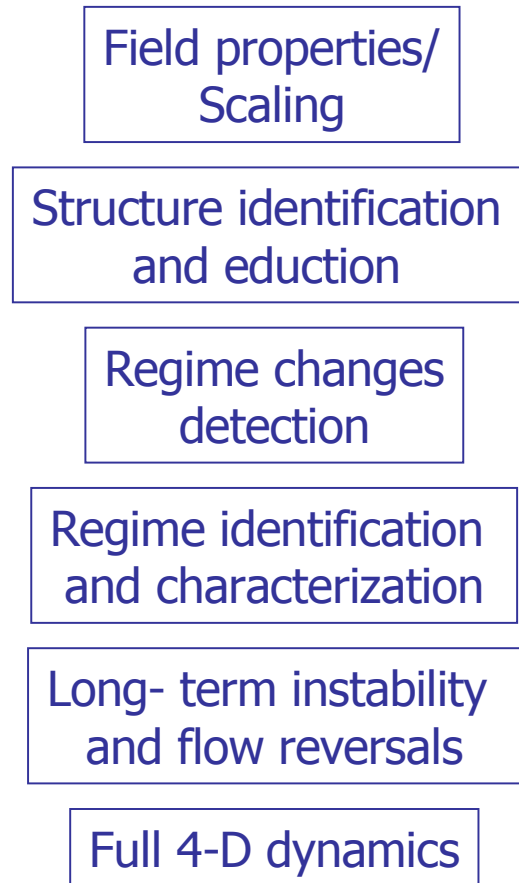
$Ra=2 \times 10^{14}$

Complementarity & Synergy in R-B

Experiments

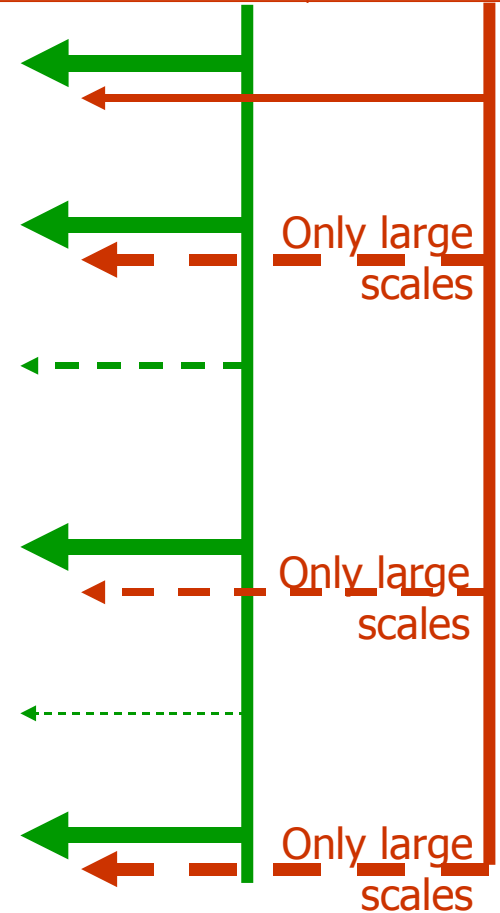


Phenomena, Parameters



Computer Simulations

DNS/LES $Ra < 10^8 / 10^9$	Modelling $Ra < 10^{17}$
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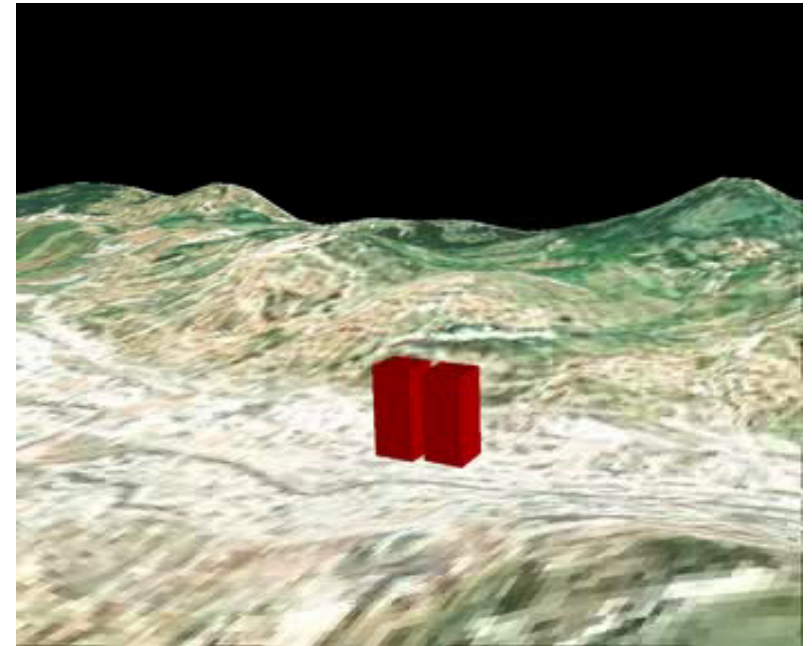
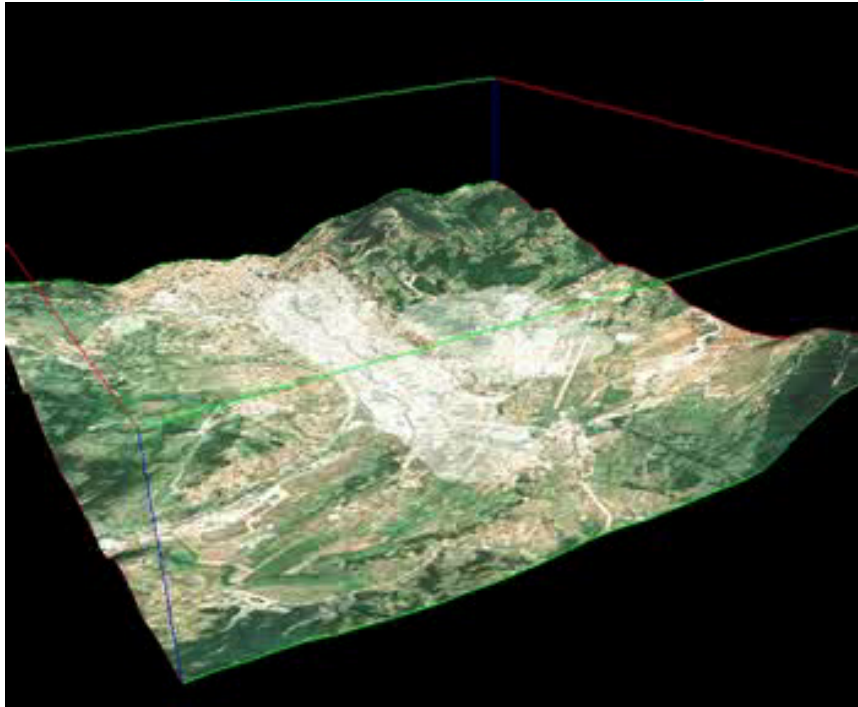
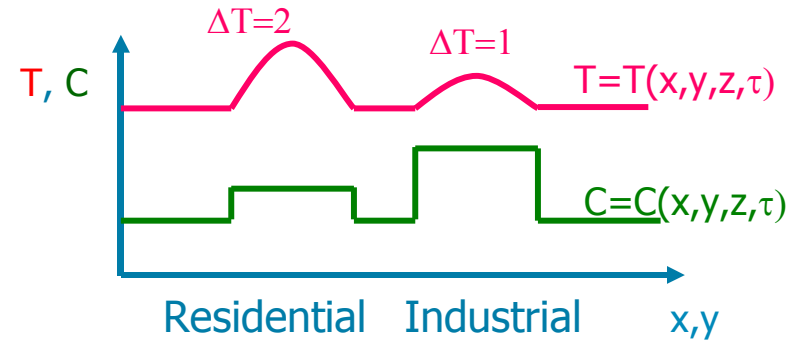
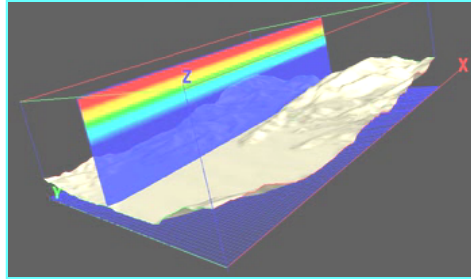


Synergy of experiments and simulations of R-B: A summary

- Experiments proved invaluable in detecting some new physics, e.g.:
 - long-term oscillation and the phenomenon of sudden or gradual reversal of flow direction ($\tau=200-2000$ sec);
 - detecting a change of regimes, etc., but
 - limited to point- or (local) plane measurements, and usually **only a single field (velocity or temperature or..)**!
- Computer simulations are uncontested in providing 3-D time dynamics (4D field) and subtle flow and structural details, but
 - DNS and LES very demanding on computational resources (**only low Ra's and short real times!**)
 - VLES/T-RANS (hybrid RANS/LES): the only viable tools for very high Ra's, but burdened with modelling approximations!
- "Together, we win!"

T-RANS of pollutant dispersion in a town valley

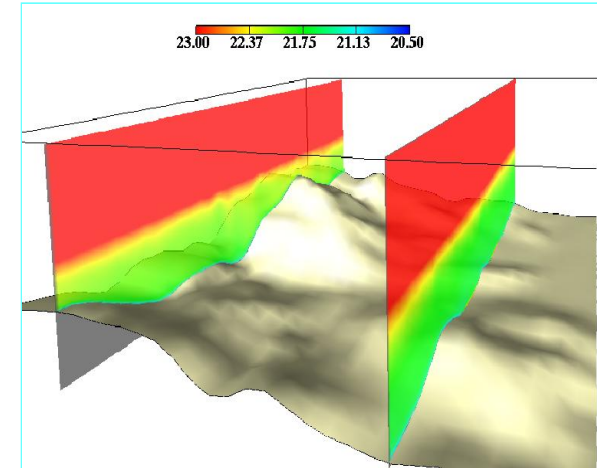
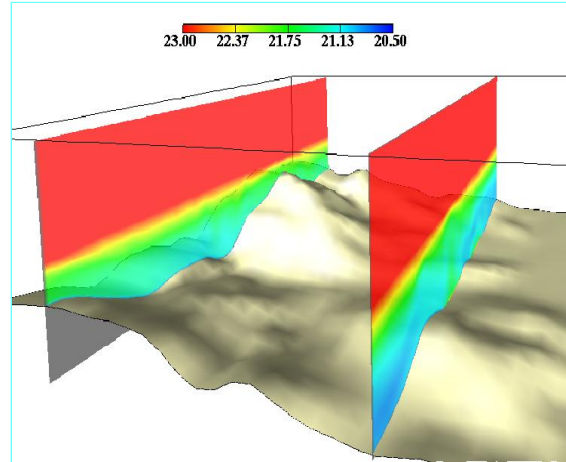
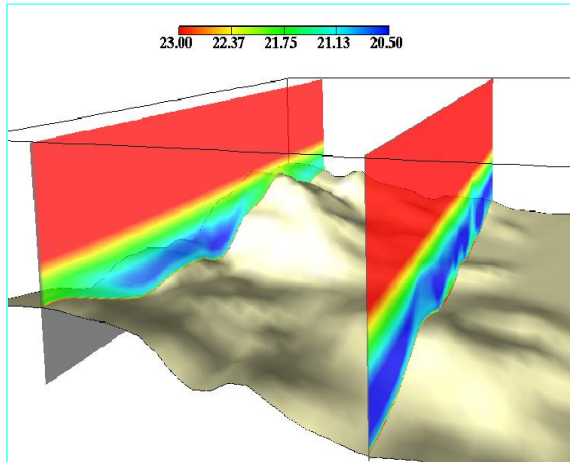
Diurnal cycles for a windless period capped by an inversion layer, with imposed ground temperature and emission scenarios



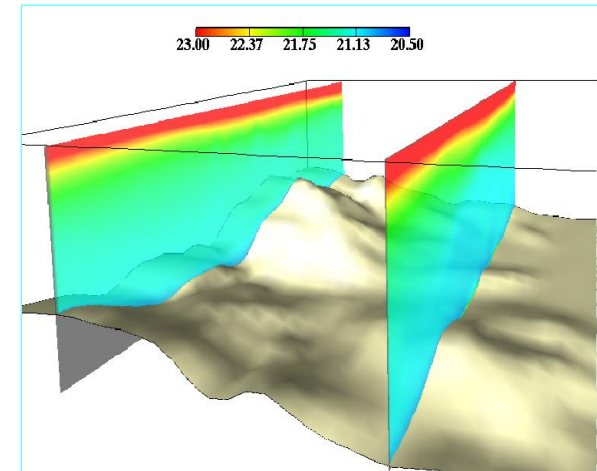
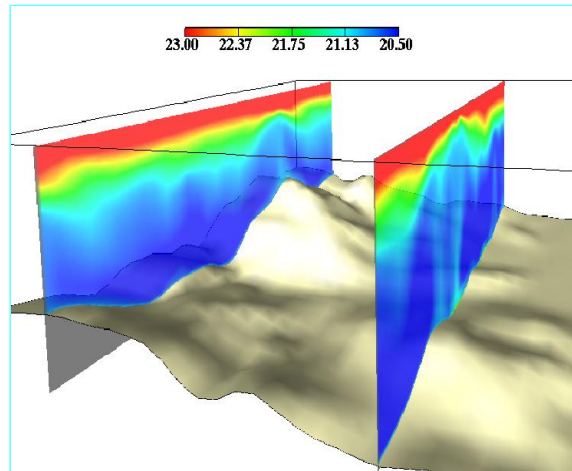
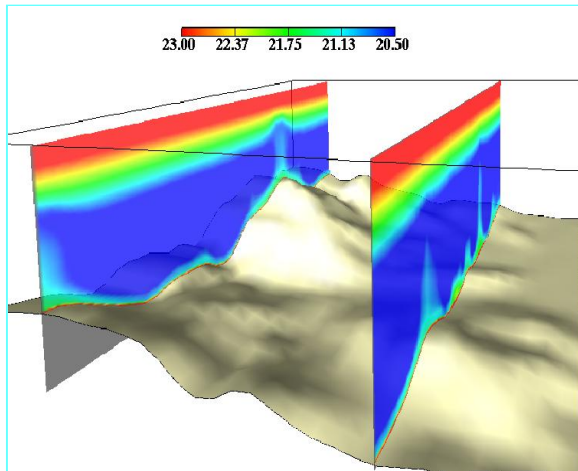
Time evolution of the potential temperature

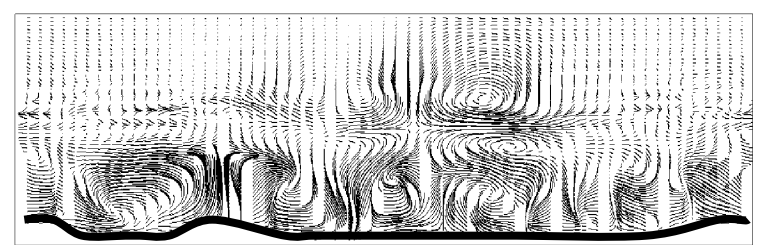
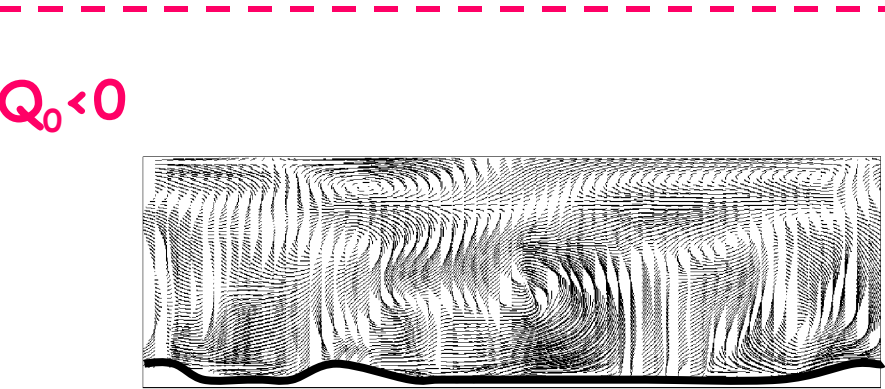
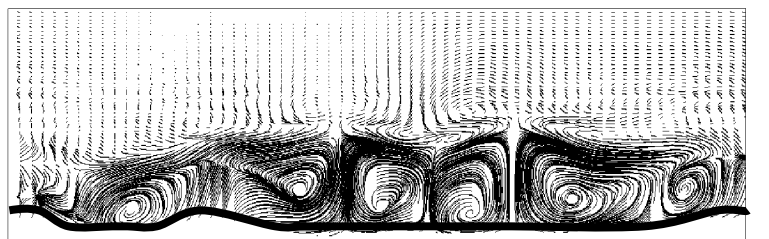
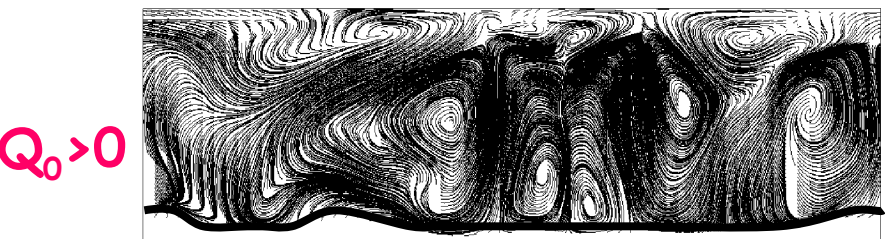
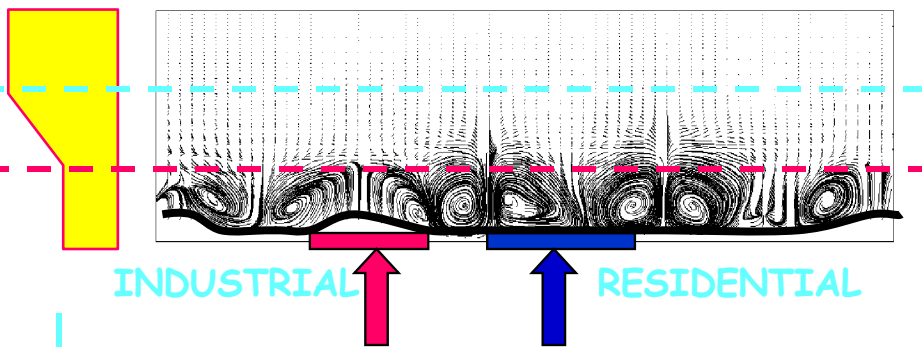
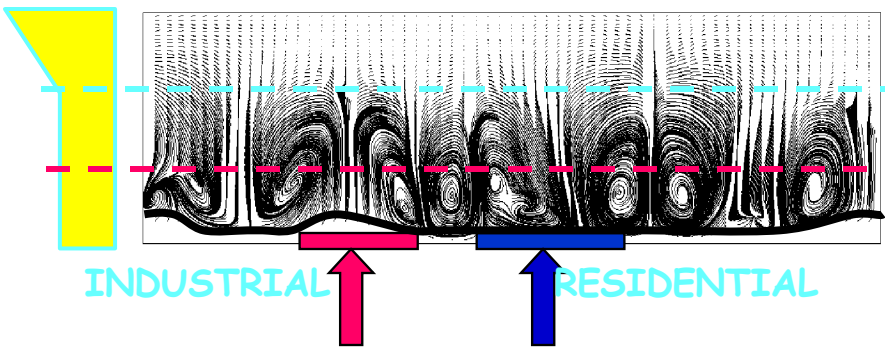
→ Time

Strong stratification



Weak stratification

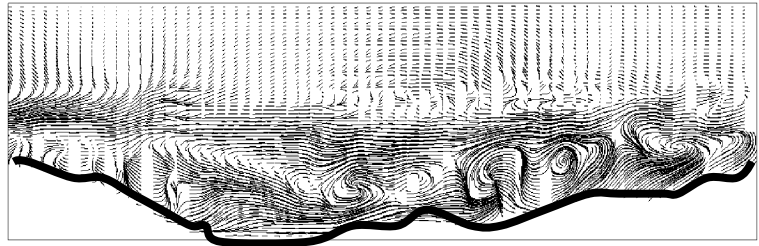
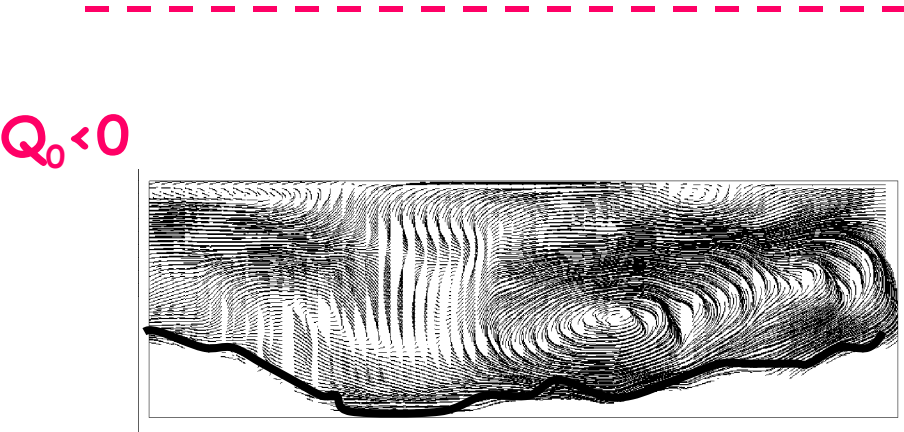
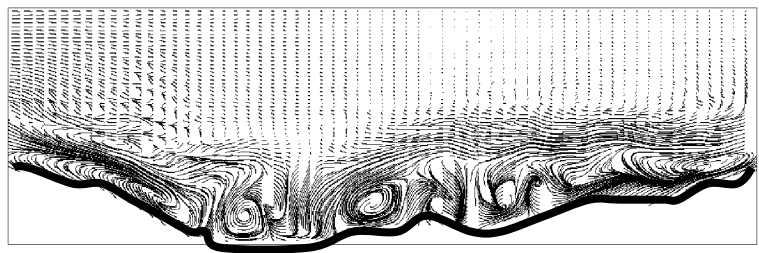
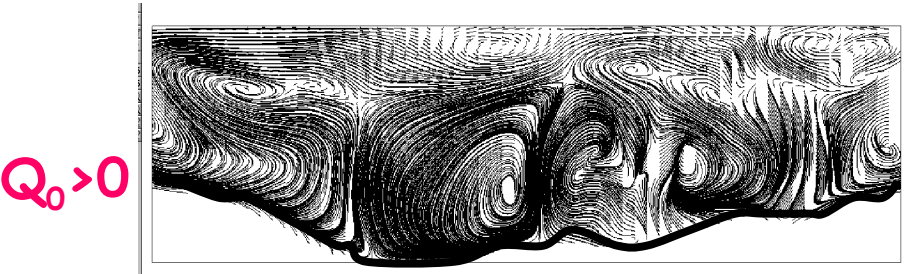
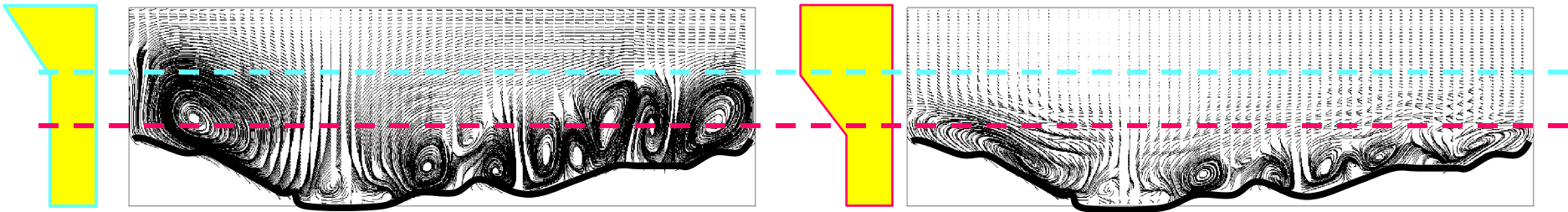




TIME

CASE (I): weak stratification

CASE (II): strong stratification

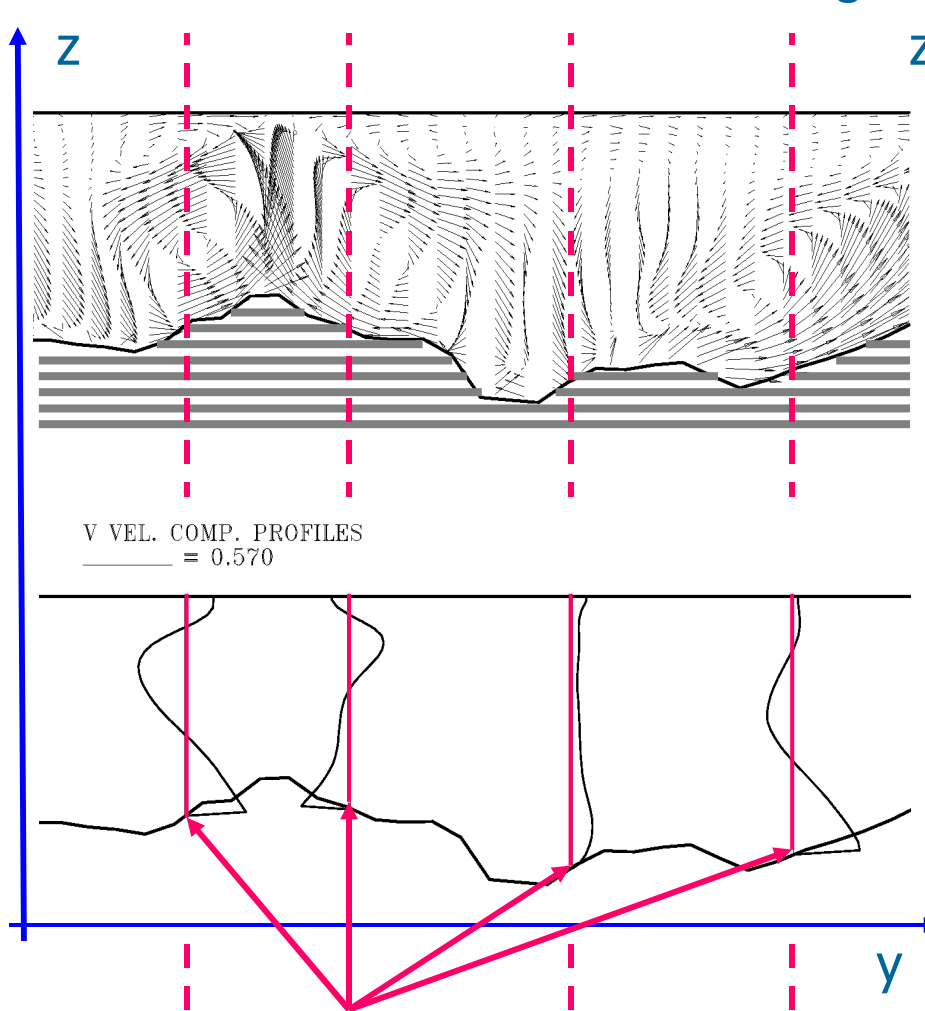


TIME

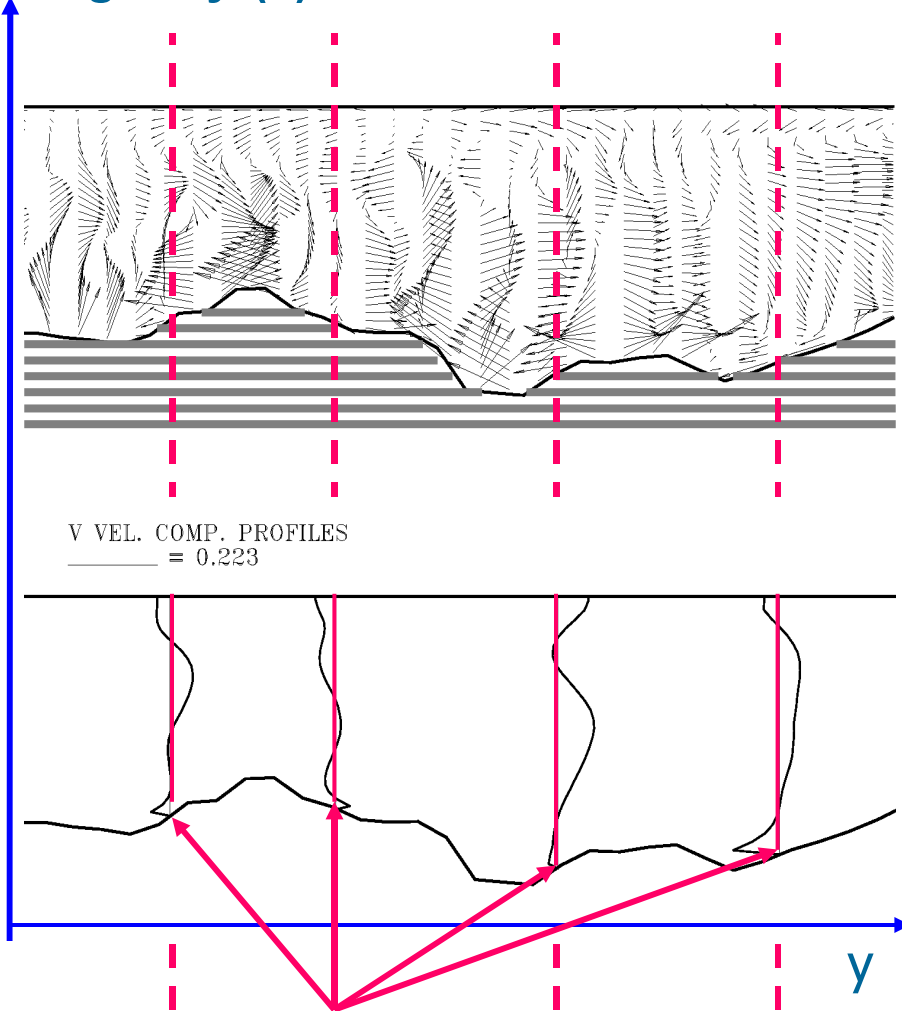
CASE (I): weak INITIAL STRATIFICATION CASE (II): strong

Instantaneous trajectories in vertical plane over hilly terrain

Velocity vectors and horizontal velocity component profiles 2hrs after onset of heating/cooling, day (II), weak stratification



up-slope (anabatic) inertial flow

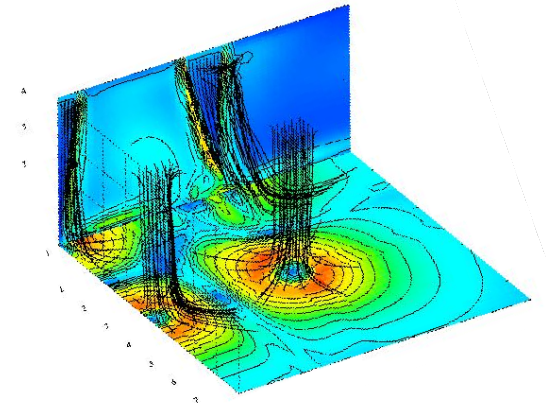
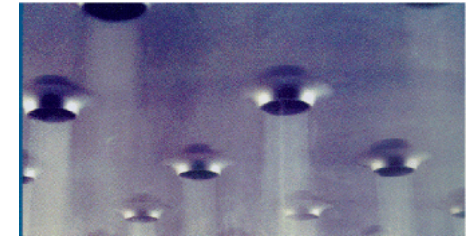


down-slope inertial flow

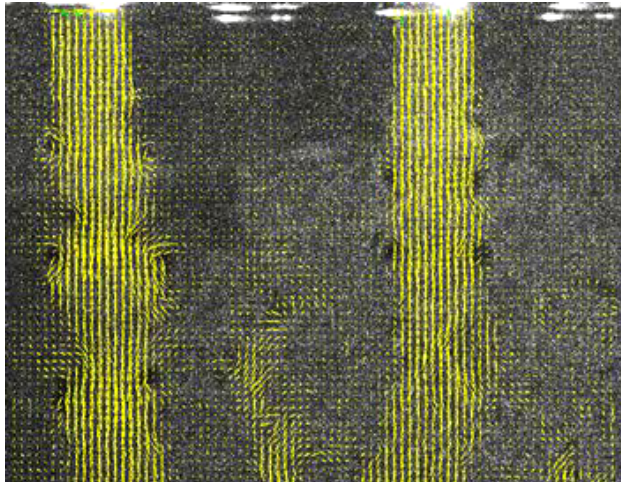
1. Impinging flows and heat transfer at *higher* Re numbers

Impinging flows and heat transfer

- Impinging jets: one of the most frequent configuration for efficient heating and cooling of solid surfaces
- Optimum performances depend on a number of parameters and no unique criteria exist
- In addition to technological interest, Impinging jets contain a number of interesting physical events and phenomena
- Challenges: identification of flow and turbulence structure, their interaction with heated surface (thermal imprints), heat transfer mechanism and its control
- Most studies confined to a single jet at relatively low Re numbers, but extra effects in multiple jets (jet-jet interaction, wall-jets collision, ejection fountains, embedded vortices, jets in cross-flows,...



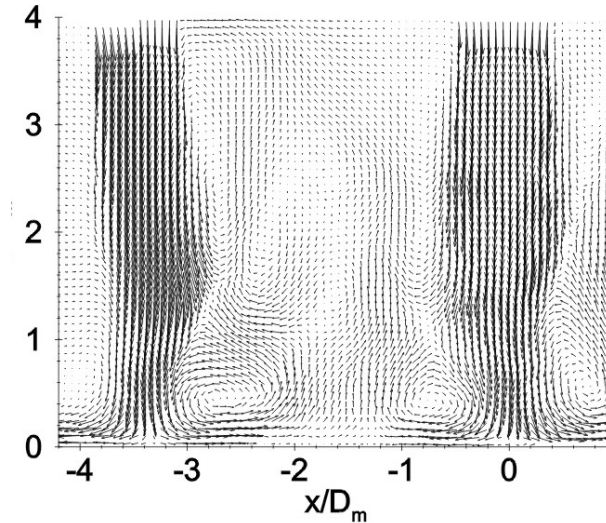
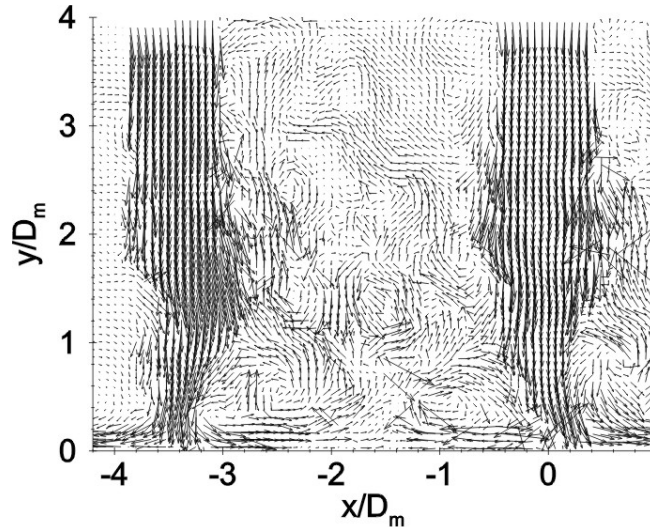
Impinging jets: potential and limitations of experiments



PIV of multiple round jets
impinging on a flat surface

(Geers, Tummers, Hanjalic, *Exp. Fluids*, 2004)

Original (left) and POD filtered (left)
snapshots.

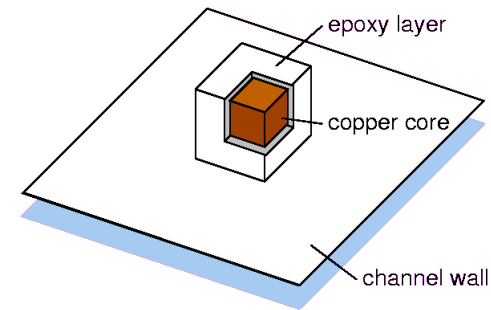


Computer simulations of Impinging Flows

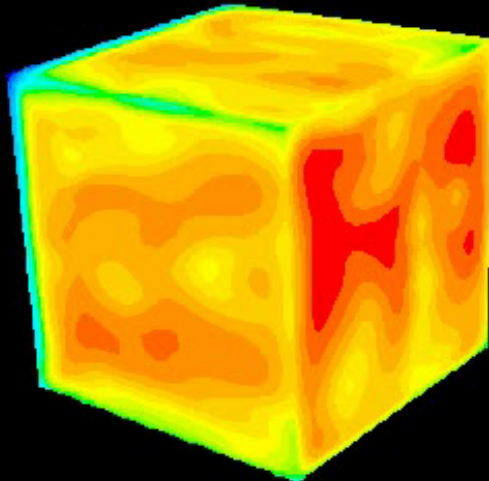
LES Niche: low Re's, separated flows
(electronics cooling, gas turbine blades,...)

Heat transfer on a multi-layered wall-mounted cube in a matrix ($Re \sim 10^4$)

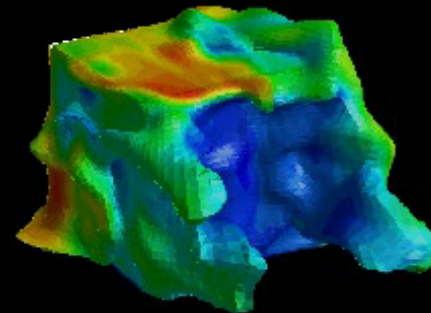
(Nićeno and Hanjalić, 2001, 2002)



Surface temperature



Thermal plume (surface of $T = \text{const}$ coloured by fluid velocity)



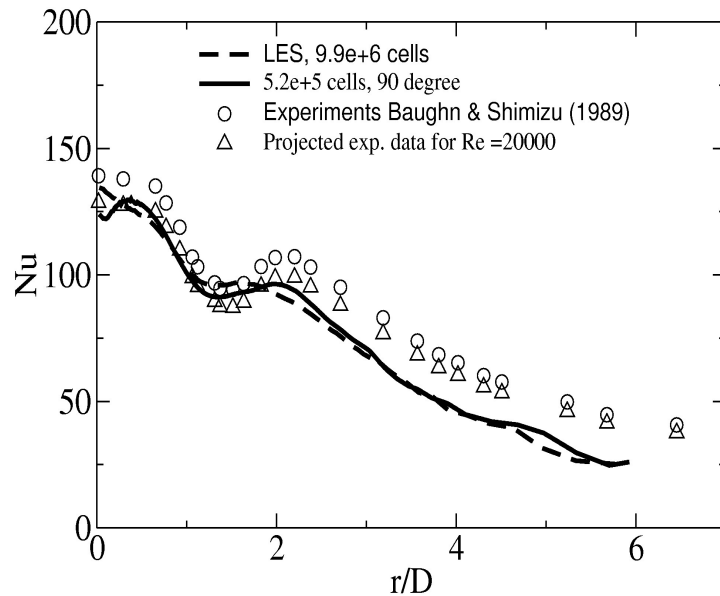
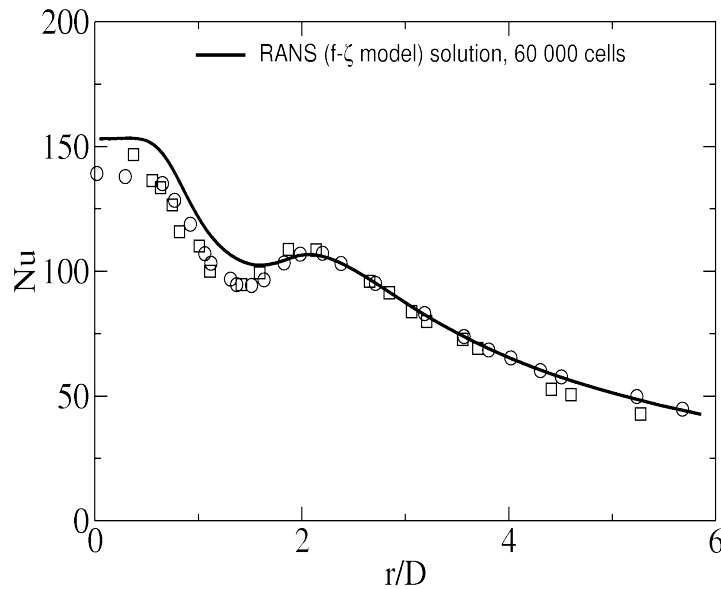
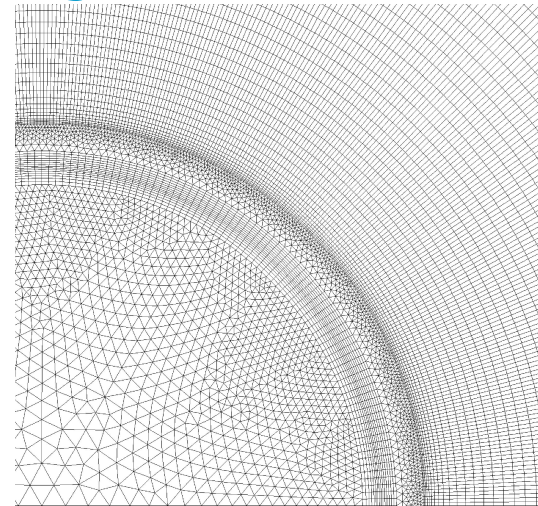
Computer simulations of Impinging Flows

Real challenge: attached impinging flows

A single round impinging jet

$Re=20.000$, $H/D=2$, unstructured grid

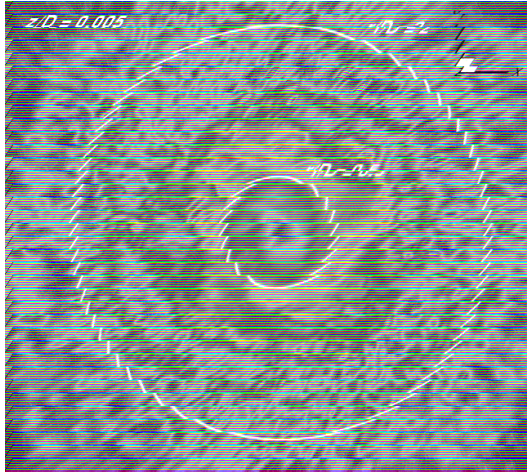
(Hadziabdic and Hanjalić, 2004/05)



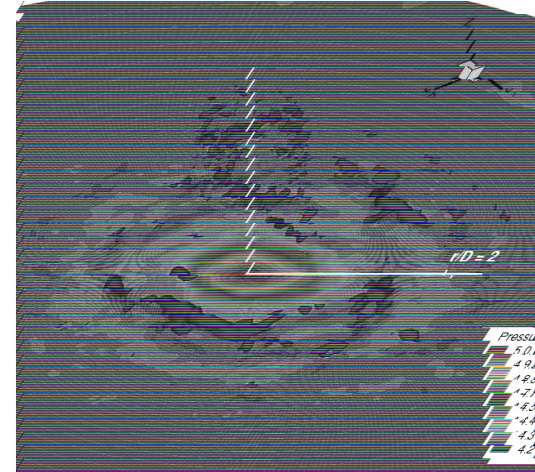
Computer simulations of Impinging Jets

Re=20.000, H/D=2 (*Hadziabdic and Hanjalić, 2004/5*)

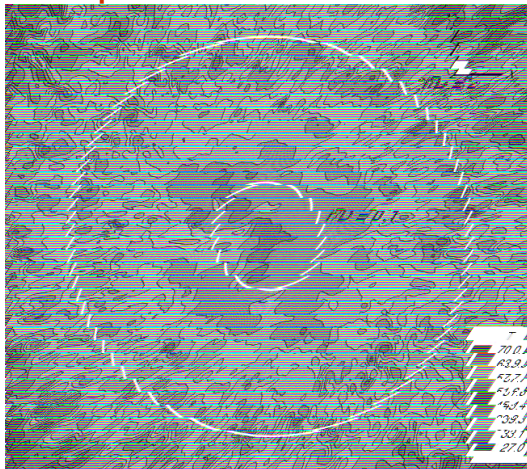
Pressure field: top view



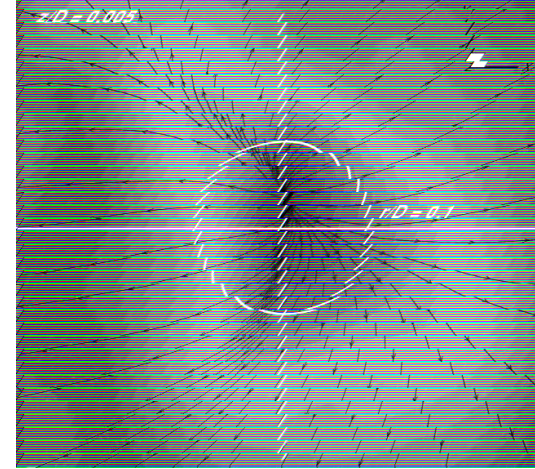
Pressure field: side view



Temperature field

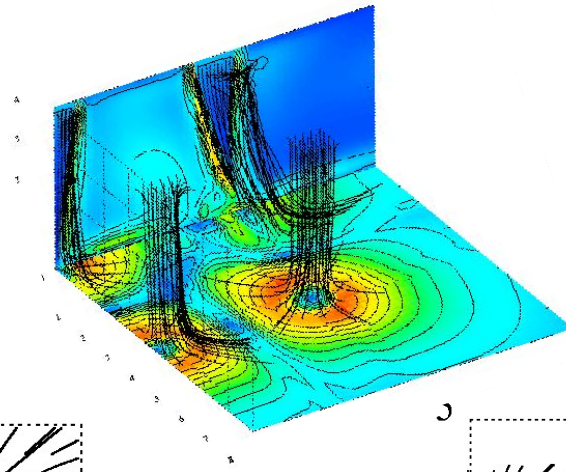


Stagnation point meandering

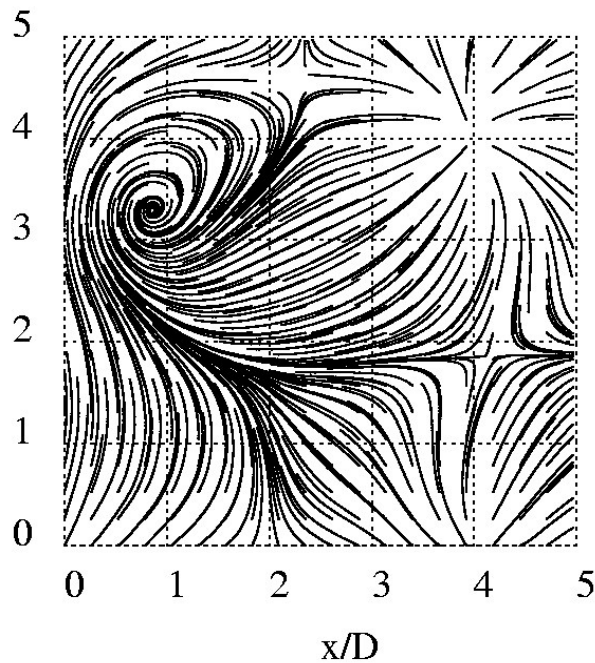


Some RANS-detected anomalies in multiple-impinging jet

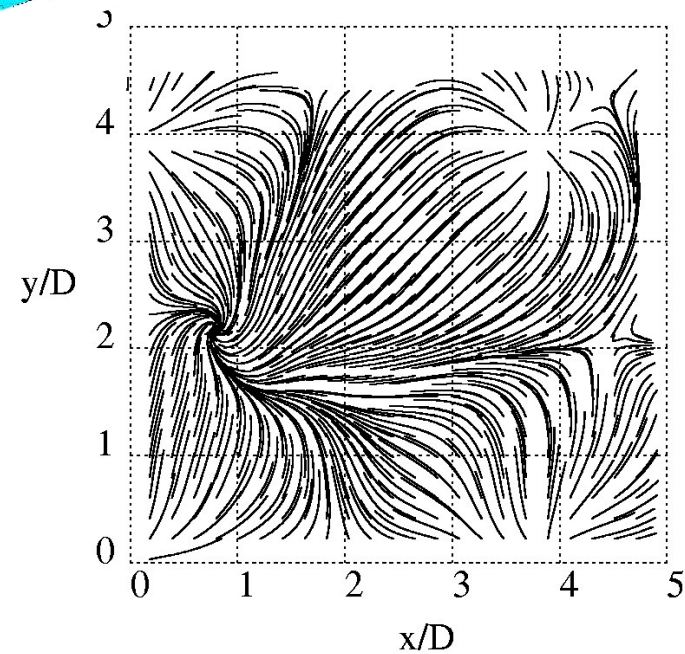
Inherent asymmetry,
bifurcation, embedded
vortices



RANS computations

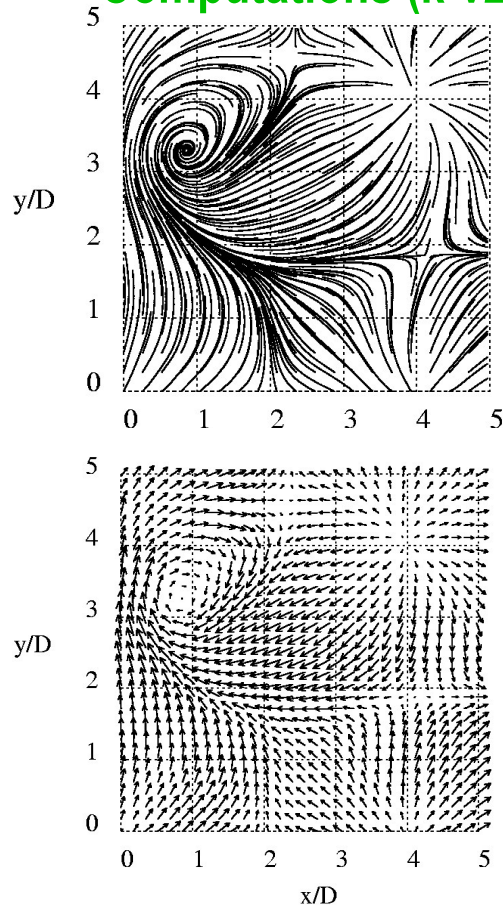


PIV measurements

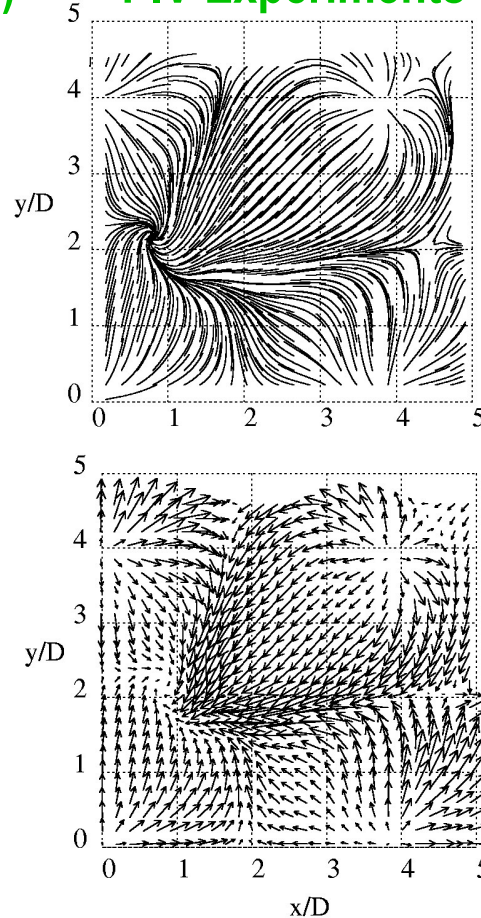


Averaged streamlines and velocity vectors for the square jet arrangement at $y/D=0.54$ above the plate *(Geers, Tummers, Hanjalic, Exp. Fluids, 2004; Thielen, Jonker & Hanjalic, IJHFF, 2003)*

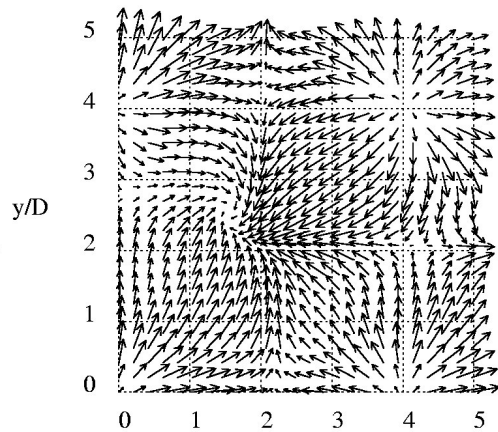
Computations (k-v2-f)



PIV Experiments



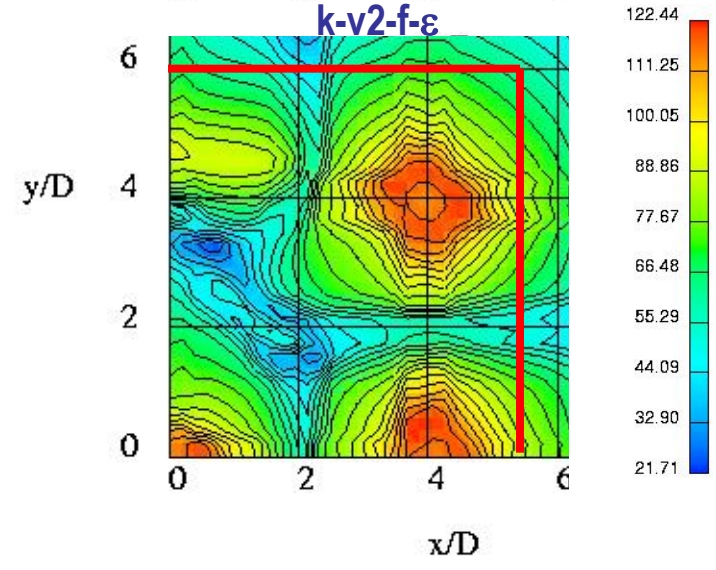
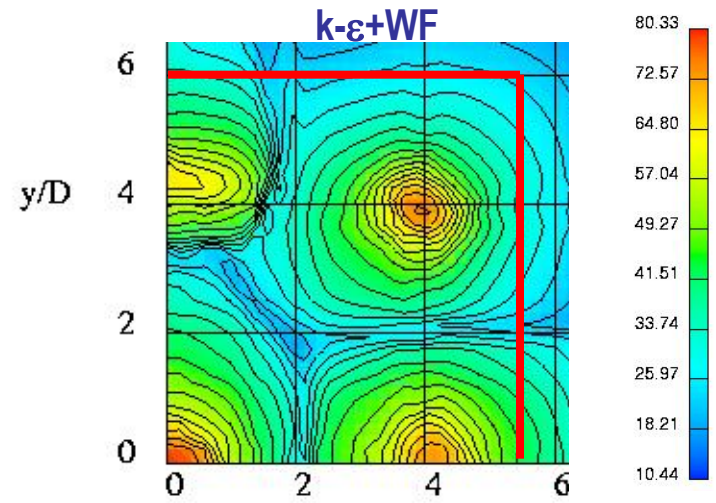
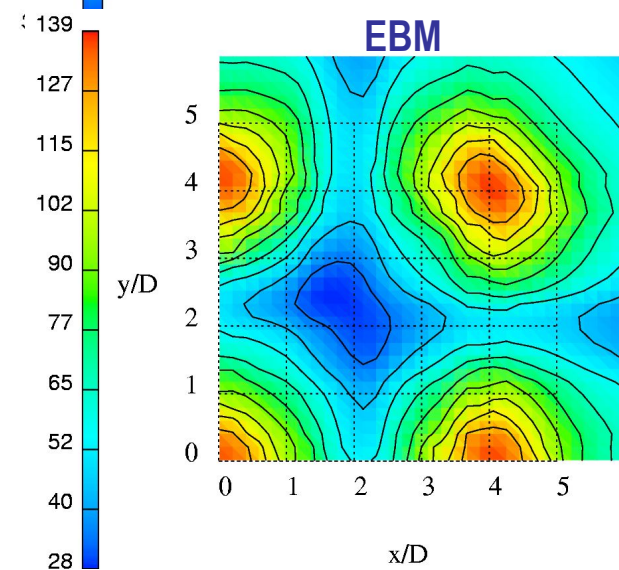
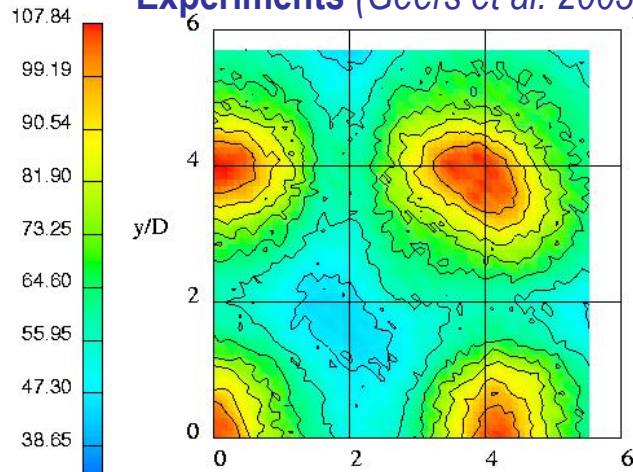
Computations (EBM)



Computed Nusselt number for square jet arrangement

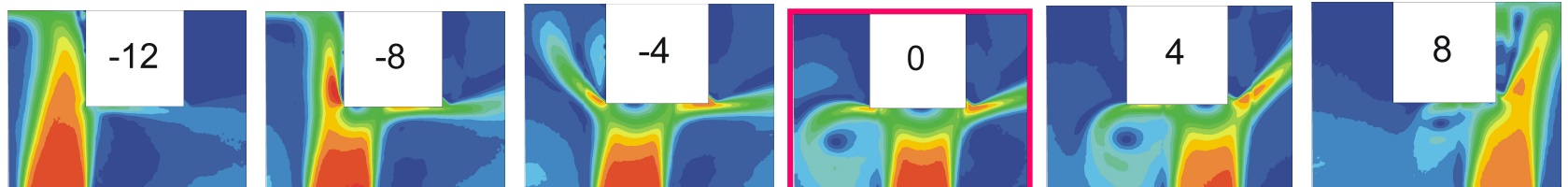
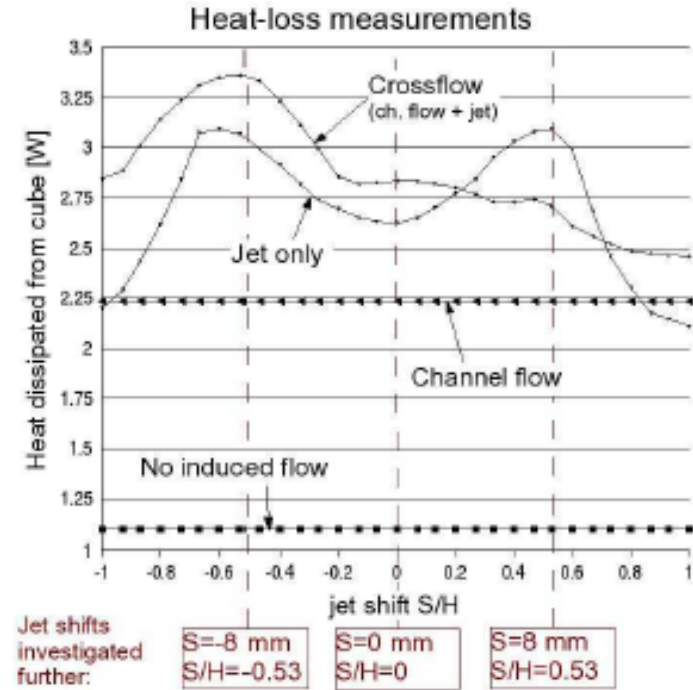
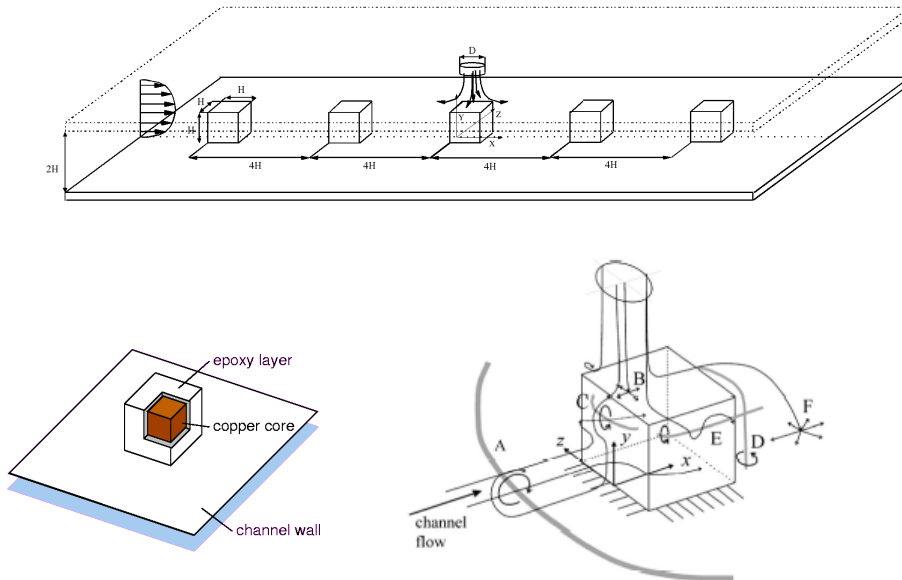
(Thielen, Hanjalic, Jonker, Manceau, IJHMT'03)

Experiments (Geers et al. 2003)



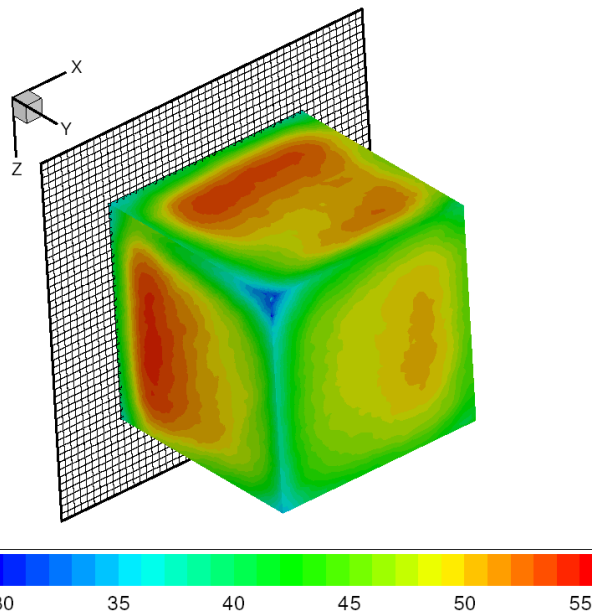
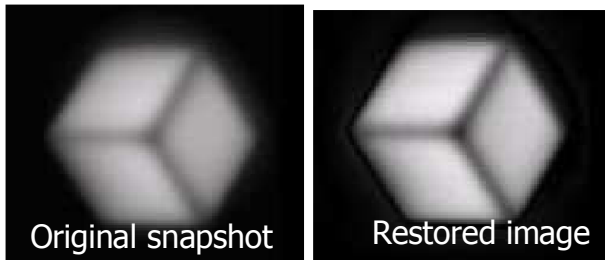
Impinged cube in a cross-flow

(Flikweert et al., 2005)



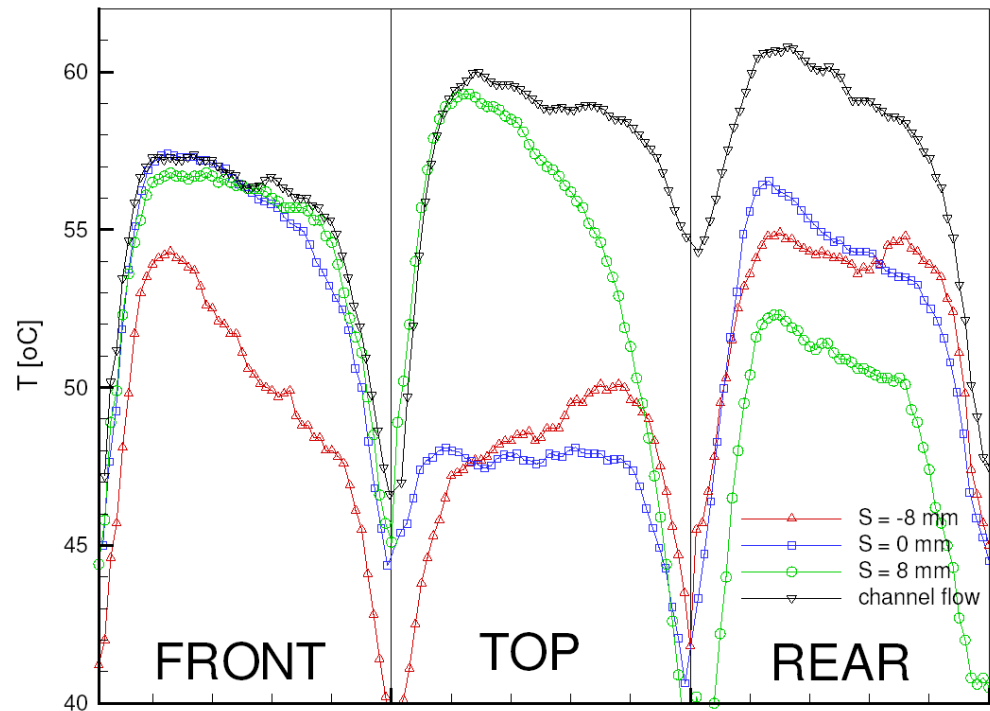
Impinged cube in a cross-flow: Surface temperature

Infrared Thermography (Flikweert et al., 2005)

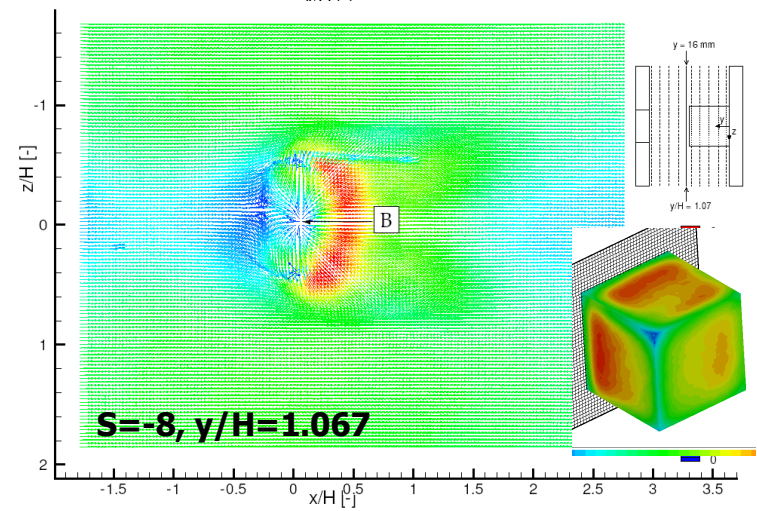
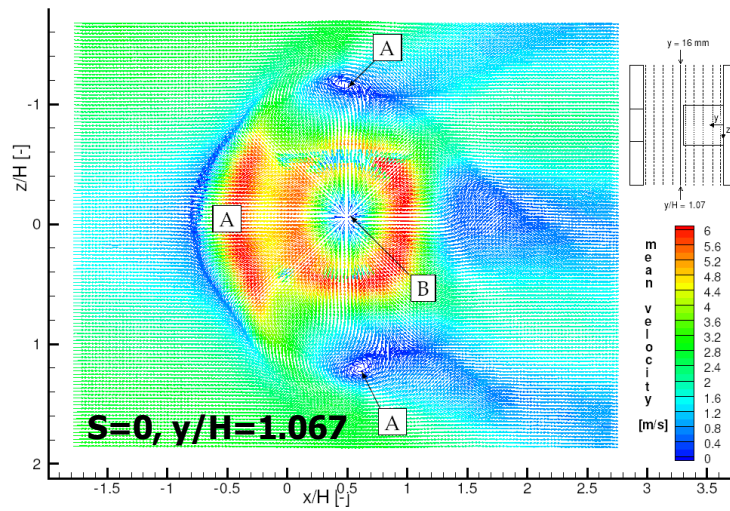
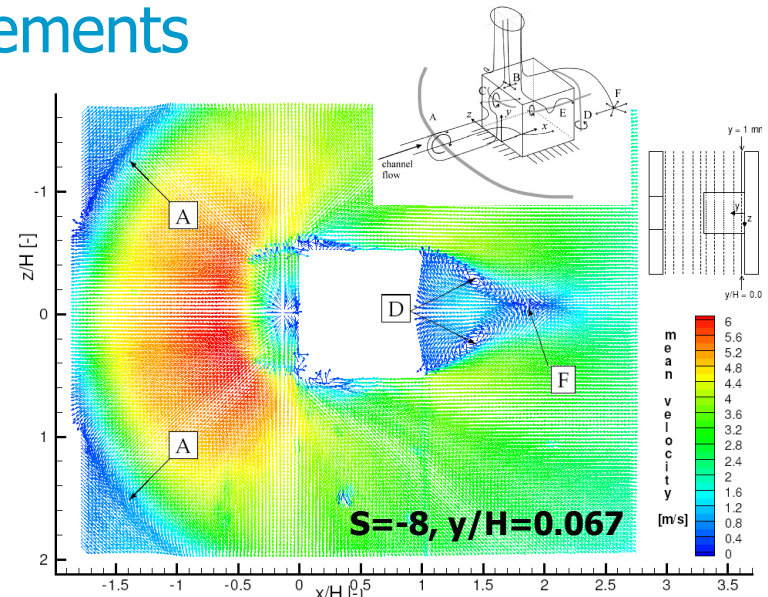
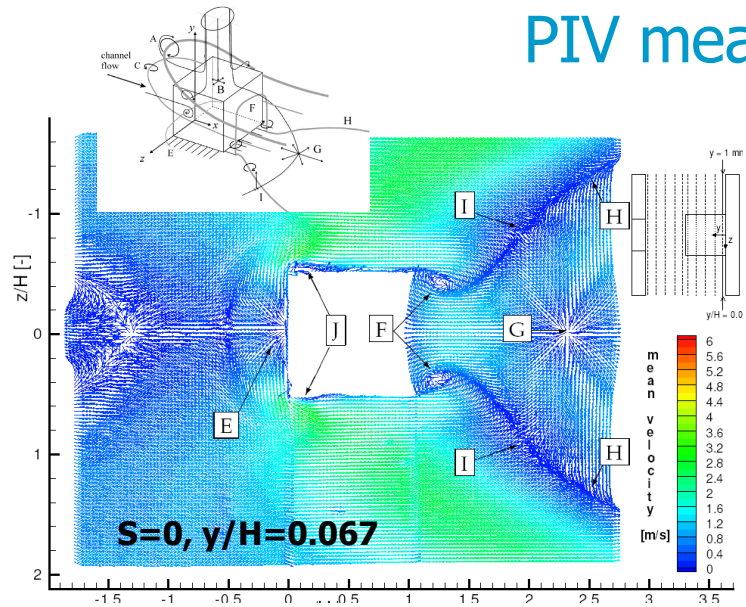


Challenges:

- optical accessibility
- image deformation and degradation
- low time resolution (only time average)

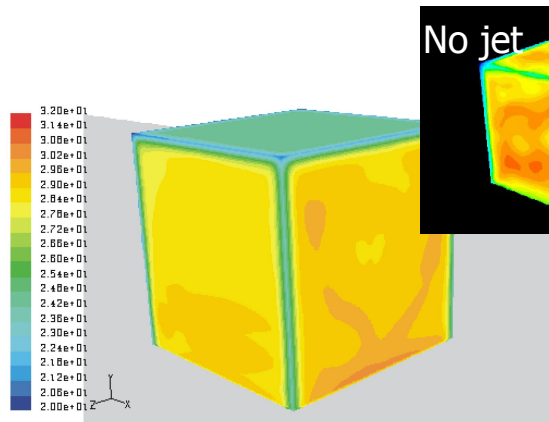
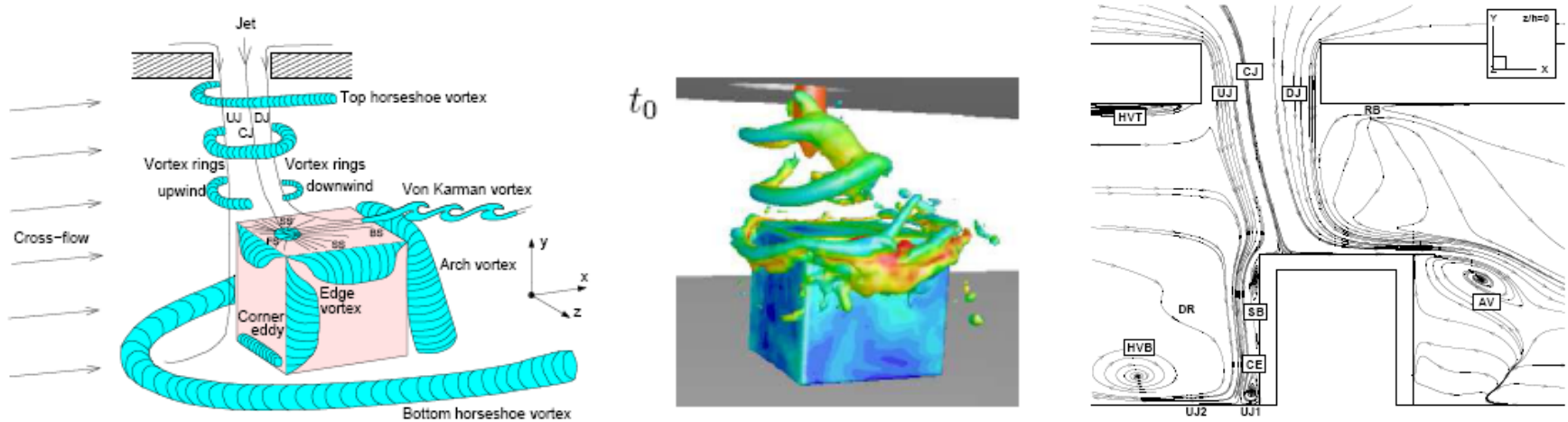


Impinged cube in a cross-flow PIV measurements

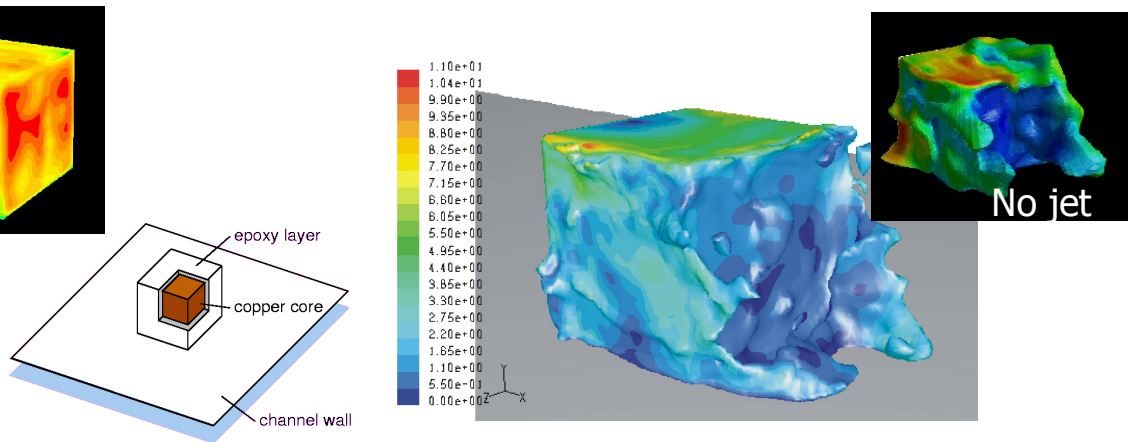


Impingement cooling of a wall-mounted cube in a cross-flow

(Conjugate LES+heat conduction in surface coating, 4.6×10^6 grid cells)



Popovac and Hanjalic 2006

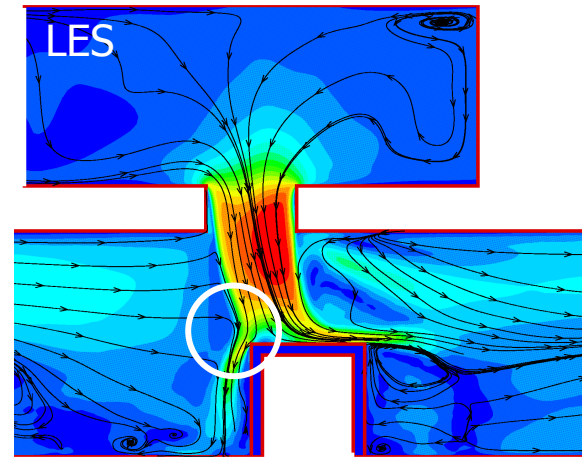
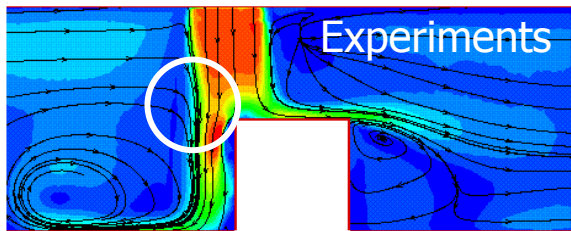


Niceno and Hanjalic, 2001,2002

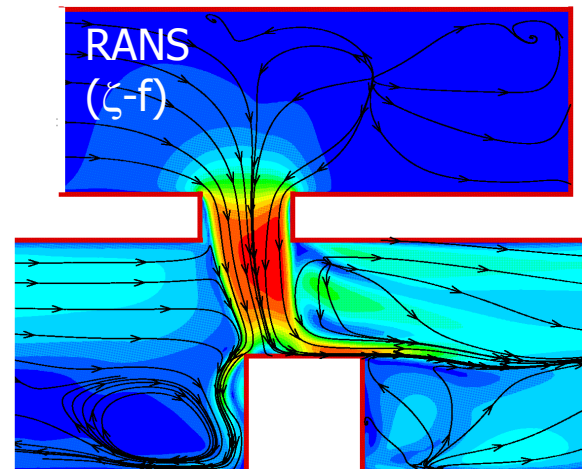
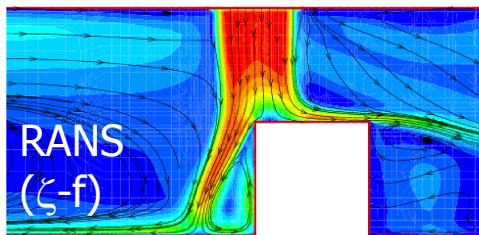
Impinged cube in a cross-flow

Challenges for modelling and simulations

- Proper imitation of experimental inflow and boundary conditions
- Grid resolution and distribution



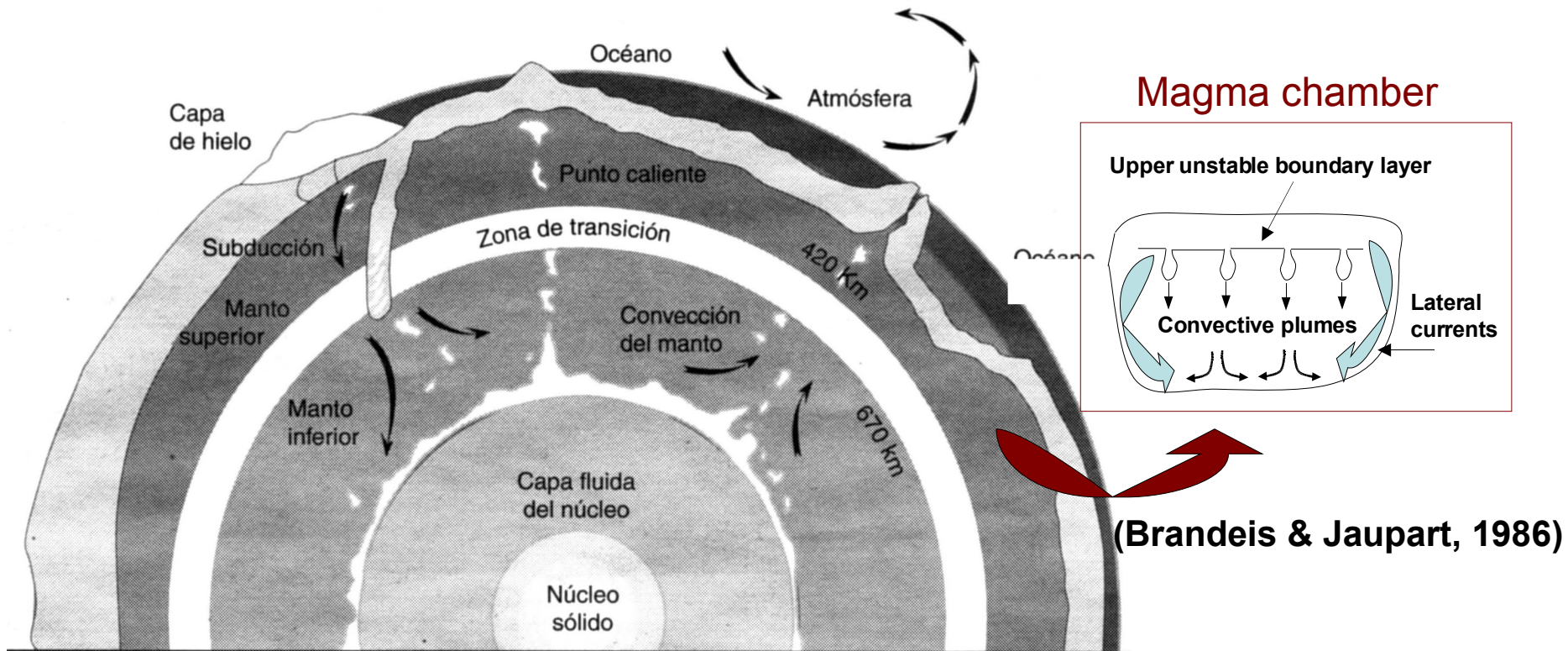
- Solution domain
- Appropriate RANS and sgs model



1. Fluid magnetic dynamo

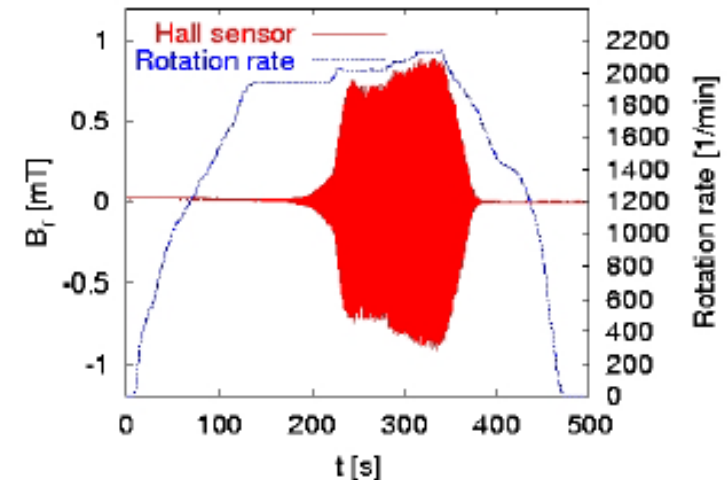
Fluid-Magnetic Dynamo (FMD)

Sketch of the Convective Motions in the Earth: Magma Chambers and Magma Eruption



Fluid-Magnetic Dynamo (FMD)

- FMD is believed to be the origin of all magnetic fields in Earth and most celestial bodies
- The basic mechanism of the field self-excitation and sustenance:
 - Thermal convection + Earth rotation drive liquid metallic core from its interior out to the mantle
 - This motion through the already existing magnetic field induces electric current, which amplifies the original field, preventing its decay with time.
- This “model”, established in twenties, was not proved until first successful experiment in 1999 in Riga (Latvia) and afterwards in Karlsruhe (Germany)



Fluid-Magnetic Dynamo (FMD)

Riga Experiments (sodium)

(*Gailitis et al. 1999*)

Major challenge for experiments:
Achieving critical $Re_m = UL/\eta = 10-10^3$,
where $\eta = 1/\mu\sigma$ (μ =magnetic permeability,
 σ =electric conductivity)

Note:

$\eta \geq 0.1 \text{ m}^2/\text{s}$ hence $UL \sim 1-10$
(difficult for liquid metals)

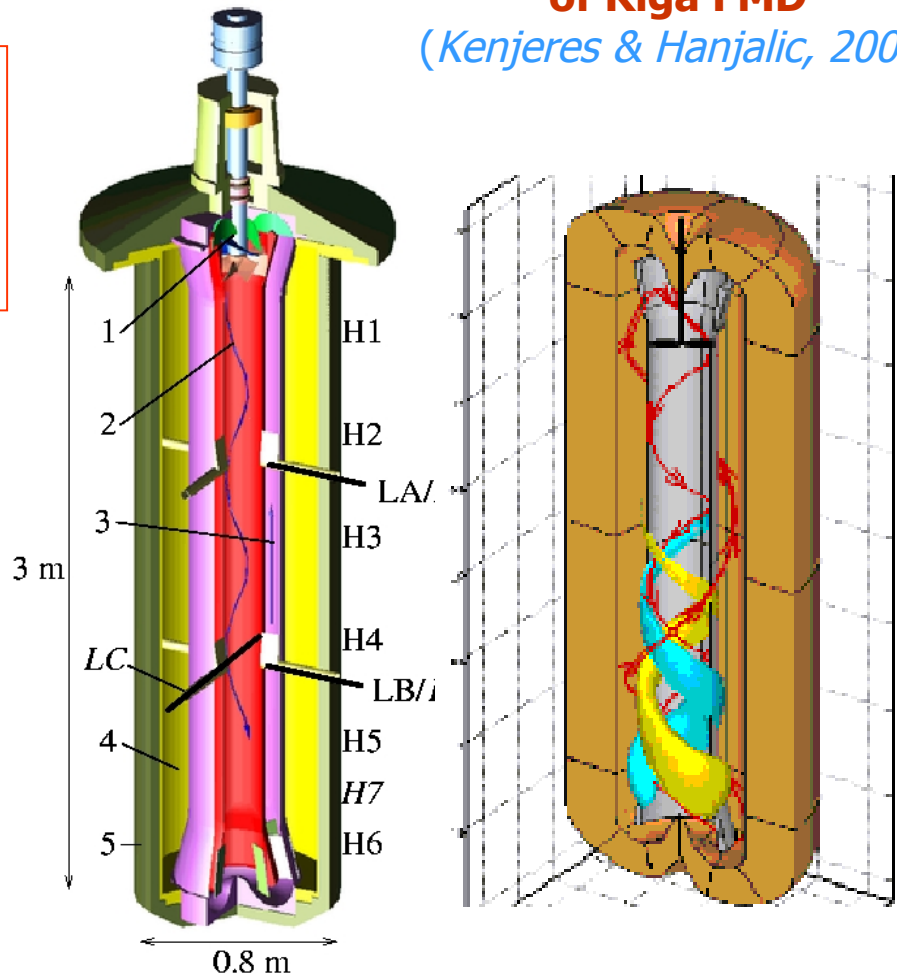
In Riga experiment:

$Re_m \sim 20$, $Re \sim 10^6$

Computer simulations: DNS of
magnetic field + URANS of velocity
field (mutually coupled)!

Computer Simulations of Riga FMD

(*Kenjeres & Hanjalic, 2005*)



T-RANS-“DNS” of the Riga Fluid-Magnetic Dynamo (FMD)

T-RANS model for hydrodynamic field

$$\frac{\partial U_i}{\partial t} + U_j \frac{\partial U_i}{\partial x_j} = \frac{\partial}{\partial x_j} \left[\nu_t \left(\frac{\partial U_i}{\partial x_j} + \frac{\partial U_j}{\partial x_i} \right) \right]$$

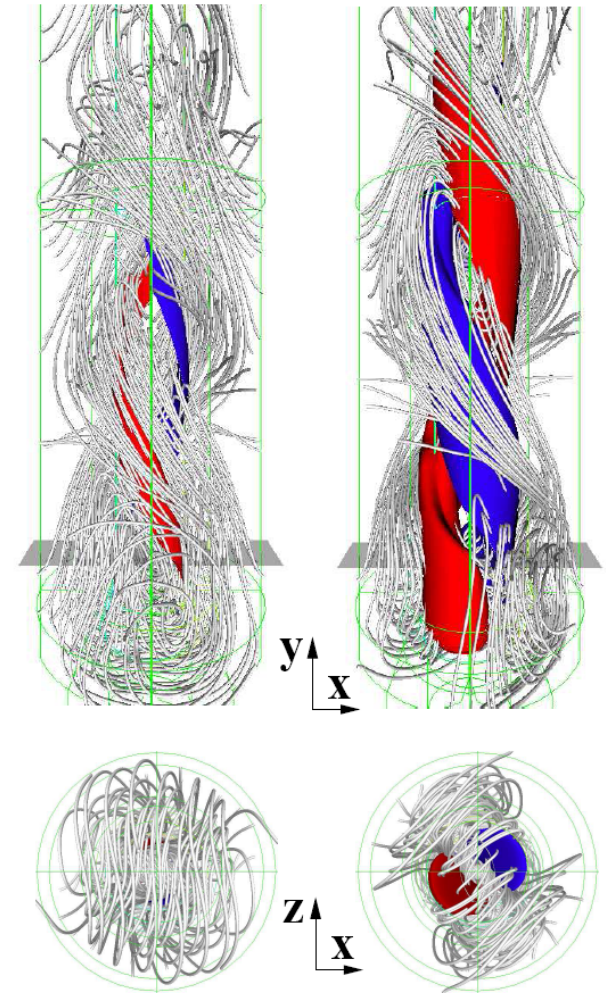
$$- \frac{1}{\rho} \frac{\partial p}{\partial x_i} - \underbrace{\frac{1}{\rho \mu_0} \left(B_k \frac{\partial B_i}{\partial x_k} - B_k \frac{\partial B_k}{\partial x_i} \right)}_{F^L = 1/\rho \mu_0 (\nabla \times B) \times B}$$

Closed with k-ε with magnetic source terms

$$S_k^M = -\frac{\sigma}{\rho} B_0^2 k \exp\left(-C_1^M \frac{\sigma}{\rho} B_0^2 \frac{k}{\varepsilon}\right); \quad S_\varepsilon^M = S_k^M \frac{\varepsilon}{k}$$

Magnetic induction equation (“DNS”)

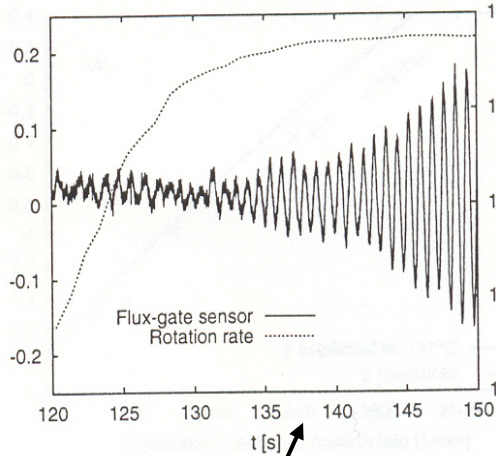
$$\frac{\partial B_i}{\partial t} + U_j \frac{\partial B_i}{\partial x_j} = \frac{\partial}{\partial x_j} \left(\frac{1}{\mu_0 \sigma} \frac{\partial B_i}{\partial x_j} \right) + B_j \frac{\partial U_i}{\partial x_j}$$



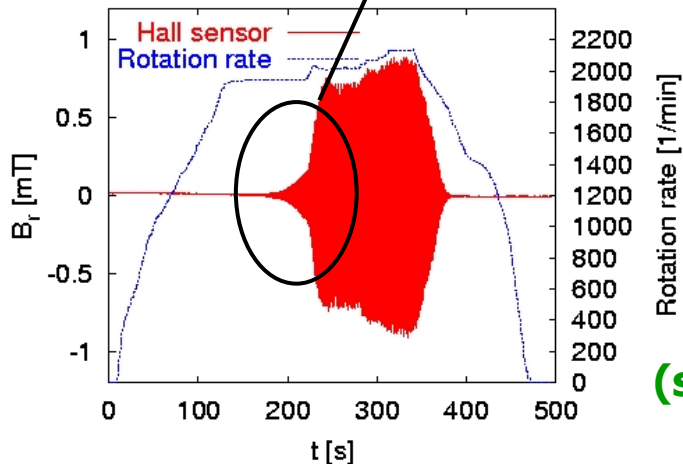
RIGA FMD: Numerical confirmation

Riga Experiments

(Gailitis et al. 1999)



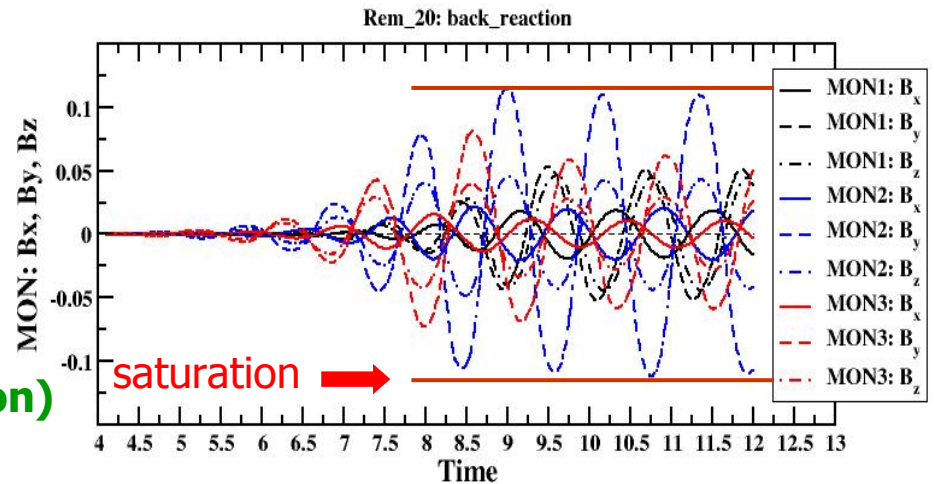
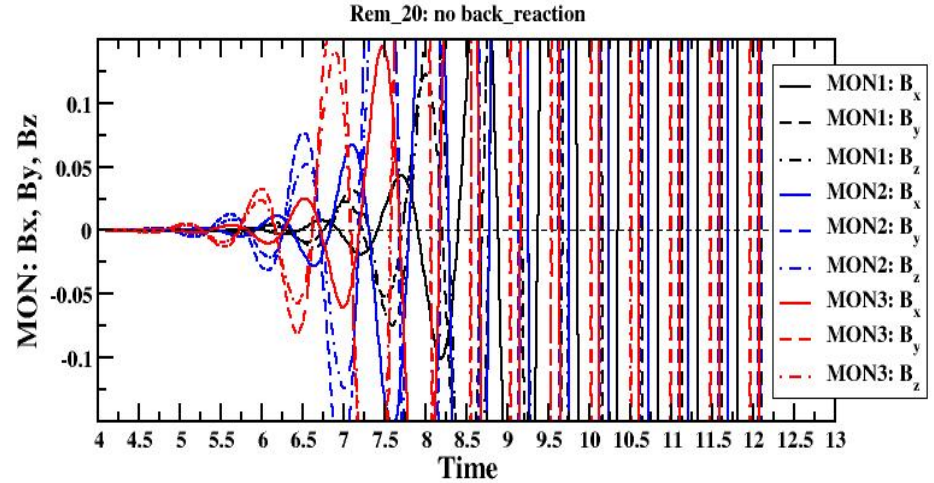
Kinematic regime

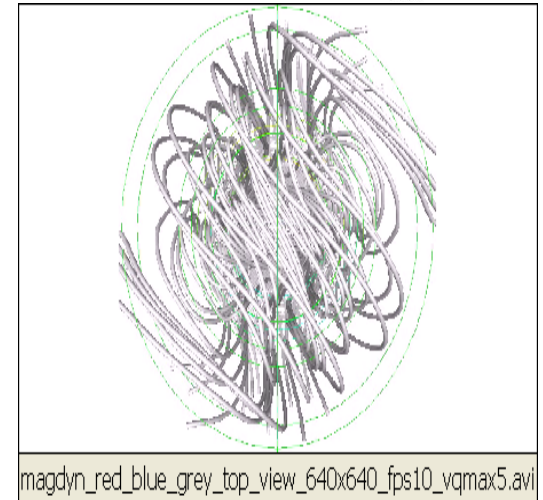
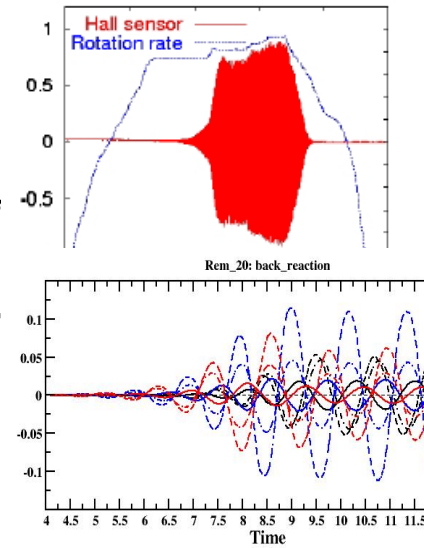
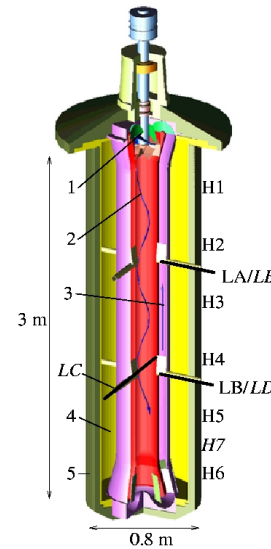
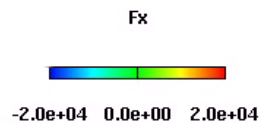
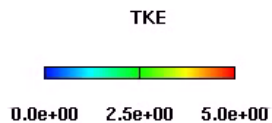
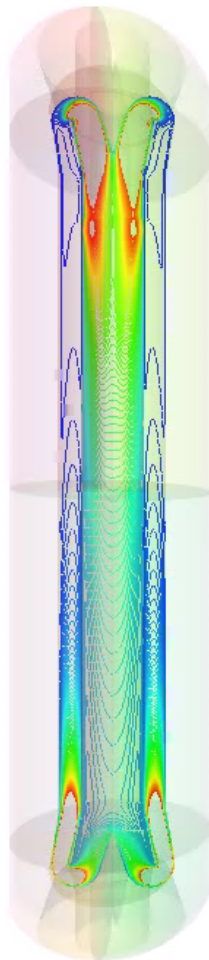
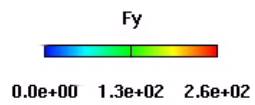
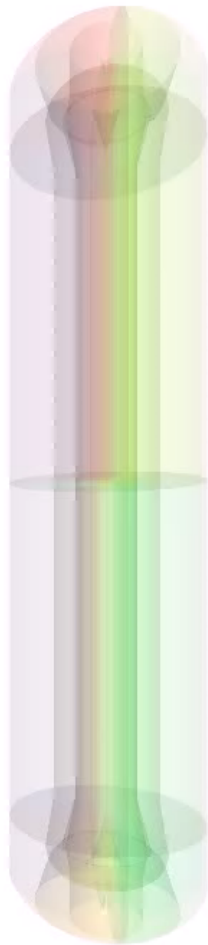
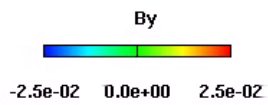
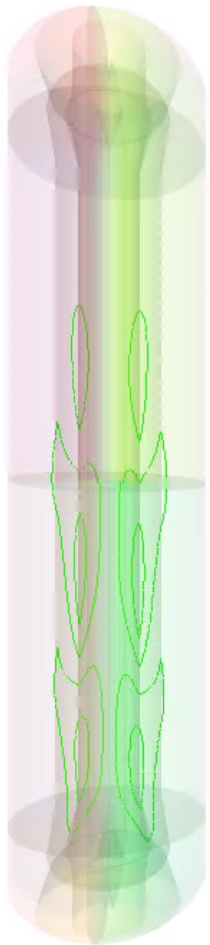


True (saturation) regime

Computer Simulations

(Kenjeres & Hanjalic, 2005)





Kenjeres & Hanjalic, 2006

Concluding Remarks

- Three sets of examples illustrate potentials and limitations of Experiments and Computer Simulations/Modelling, but also their complementarity and synergy potential;
- The robustness and repeatability will keep experiments irreplaceable in *detecting* new physics, gathering new information and databases
- Computer simulations (DNS, LES): indispensable tool for collecting high-resolution 4D information (a true research tool for *explaining* new physics), but limited to small Re and Ra Nos
- Semi-empirical models and mixed approaches (URANS, VLES, hybrid RANS/LES) complement and extrapolate DNS, LES and Experiments, though will hardly ever be accepted as a trustworthy research instrument!
- Judiciously combined, they can generate invaluable synergy!

Concluding Remarks, cont.

- Computer visualization and animations, pioneered by experimentalist, but reached full blossom with computer simulations, is growing into its own branch of science:
 - they can reveal events, phenomena, structures etc., which may be just too complex for abstract imaging in ones mind.
- This all has been made possible primarily by Computer Simulation, but the abundance of information is creating **new** problems:

“Having terabytes of data at your disposal greatly increases the chances that you can find the answers to even the toughest questions
– if you do not mind searching for a needle in a giant haystack”

(G. Ehrenman, Mechanical Engineering (ASME), February 2005).

Deficiency of the Basic EDM for Buoyant Flows

- Isotropic eddy-diffusivity model (EDM) for heat flux ("Simple Gradient Diffusion Hypothesis", SSGD) :

$$\overline{\theta u_j} = -\frac{v_t}{\sigma_T^t} \frac{\partial T}{\partial x_j} \equiv -\frac{v_t}{\sigma_T^t} \nabla T$$

- Consider two generic situations:

1. A fluid layer heated from below, $g_i \parallel \nabla T$

Outside the thin layers, $\nabla T \approx 0$ (or = 0!),

yet, the vertical heat transport

2. Vertical heated walls, $g_i \perp \nabla T$

Buoyancy source of k (and ε)

yet, the vertical $\nabla T \approx 0$!

$$q_i = -\overline{\theta u_i} \neq 0!$$

$$G = \beta g_i \overline{\theta u_i} \neq 0,$$

