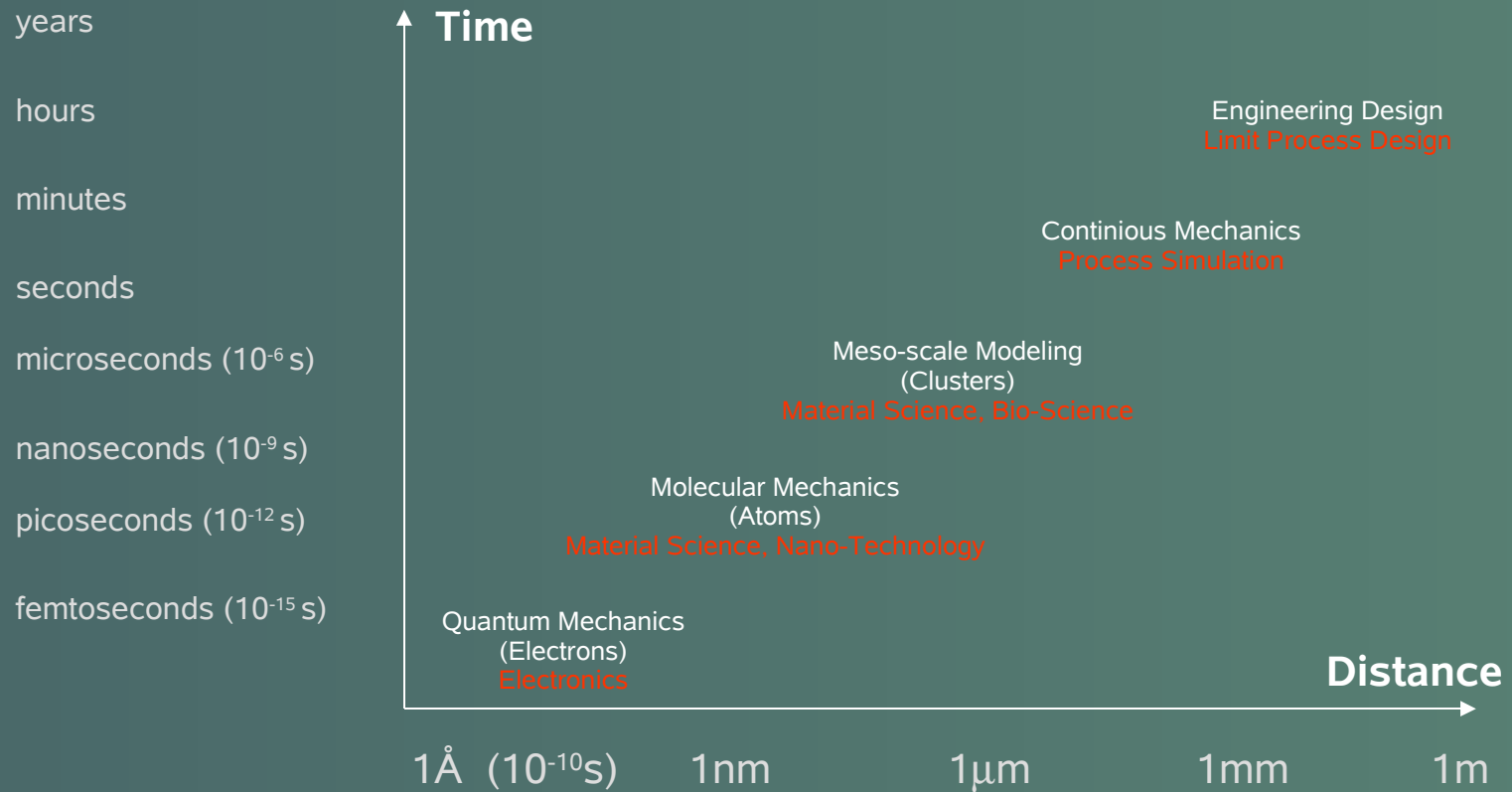


# Modeling of Micro-Fluidics by a Dissipative Particle Dynamics Method

Justyna Czerwinska

# Scales and Physical Models



# Micro- and Nanoscale effects

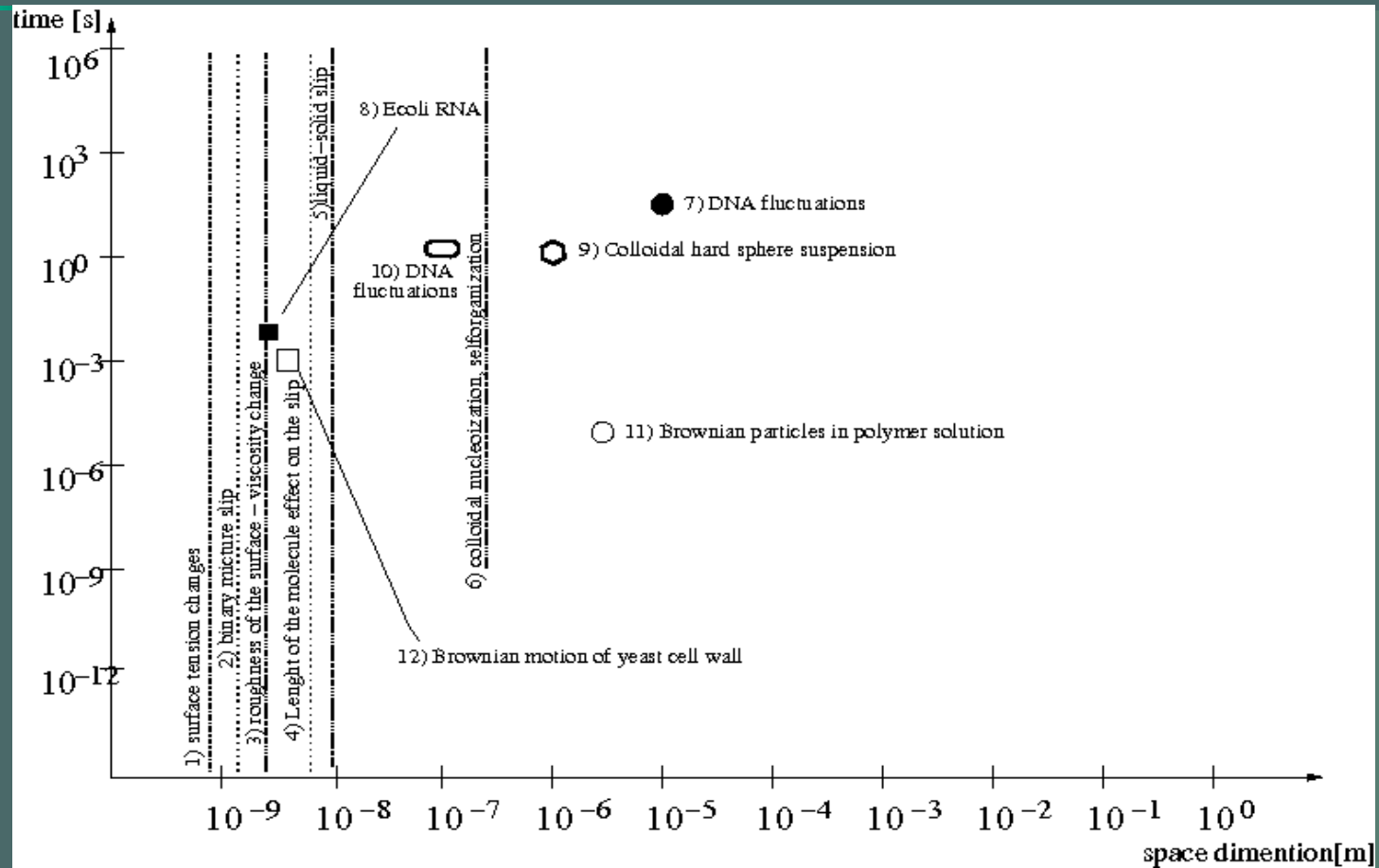


length	surface	volume
1m	1m <sup>2</sup>	1m <sup>3</sup>
1μm	1μm <sup>2</sup>	1μm <sup>3</sup>
10 <sup>-6</sup> m	10 <sup>-12</sup> m	10 <sup>-18</sup> m
S <sup>1</sup>	S <sup>2</sup>	S <sup>3</sup>

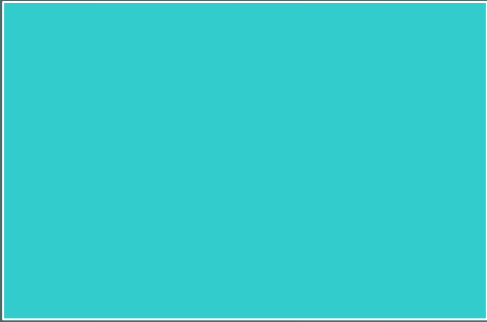
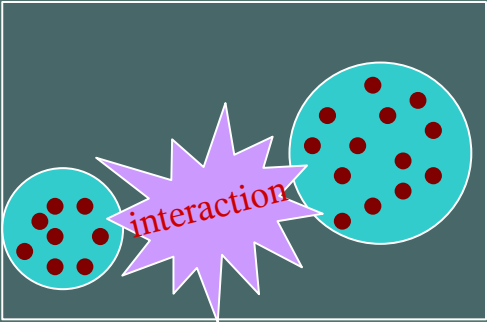
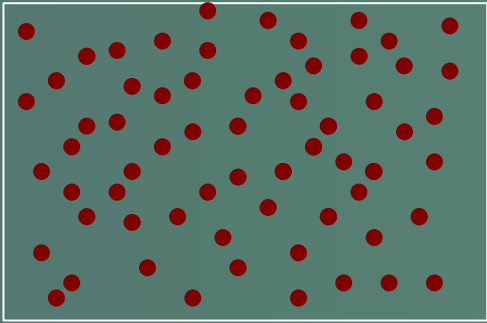
- increasing importance of surface to volume effects (surface tension important, gravity unimportant)
- slip, no-slip boundary conditions
- gas, liquid differences (1 μ m<sup>3</sup> -25 millions of air molecules, 34 billion water molecules)
- non-linear effects, thermodynamical nonequilibrium

**1mm - MACRO - MICRO PHYSICS BARRIER**

# Nano- and Micro- Effects

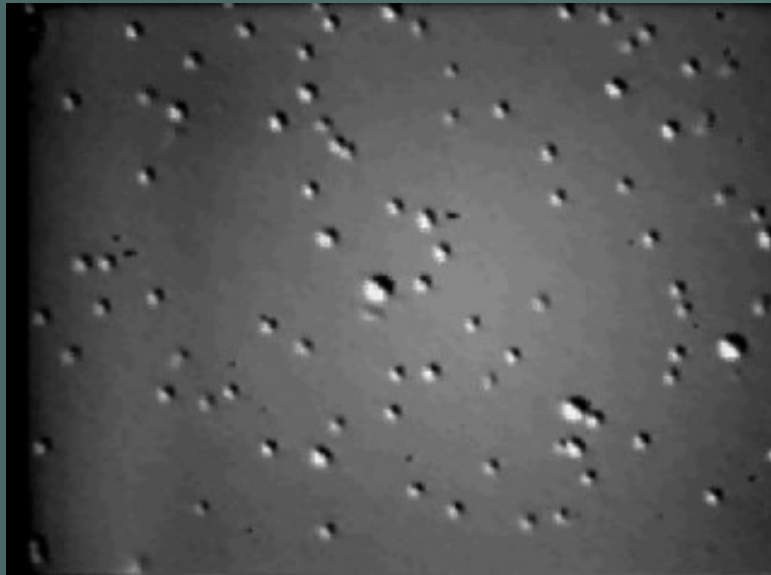


# Models of fluids

Continuum	Meso-scale	Microscopic
<ul style="list-style-type: none"><li>■ Global parameters: density, velocity, energy, temperature</li></ul> 	<ul style="list-style-type: none"><li>■ averaged properties, Brownian mechanics, stochastic equations of motion</li></ul> 	<ul style="list-style-type: none"><li>■ local description, molecules kinetic energy and intermolecular interaction potential</li></ul> 

# Meso-scale physics

Brownian motion of coiled DNA



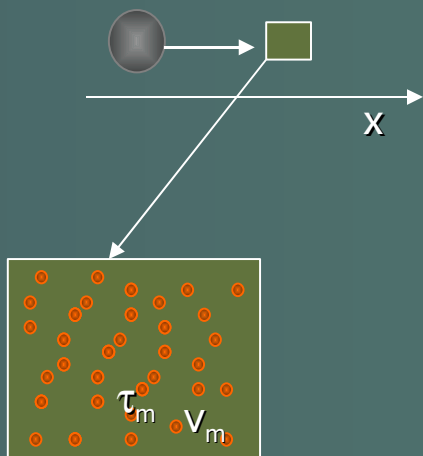
Micro-biological motor (1.50  $\mu\text{m/s}$ )



# Meso-scale physics

Langevin description of Brownian motion (1908)

$$M \frac{d\mathbf{V}}{dt} = -\xi M \mathbf{V} + F(\mathbf{R}) + \Theta(t)$$



- $\mathbf{V}$  – velocity
- $\mathbf{R}$  position;  $d\mathbf{R}/dt = \mathbf{V}$
- $\xi M \mathbf{V}$  – systematic contribution of forces
- $\xi M$  – friction coefficient
- $F(\mathbf{R})$  – external forces
- $\Theta(t)$  – random forces
- $1/\xi$  – relaxation time, without random forces

**Assumption:** Typical time scale on which collisions take place is very small compared to the evolution of the average velocity

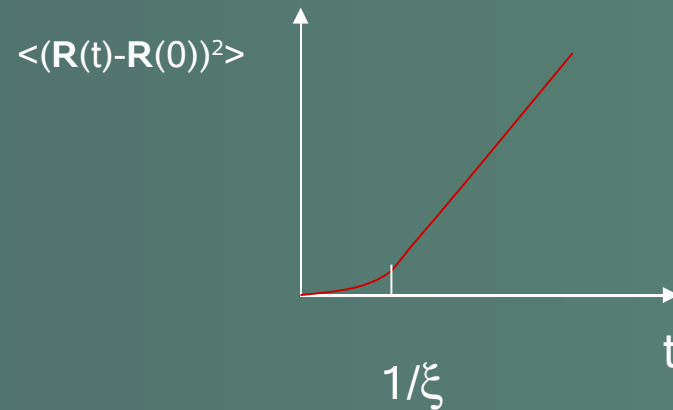
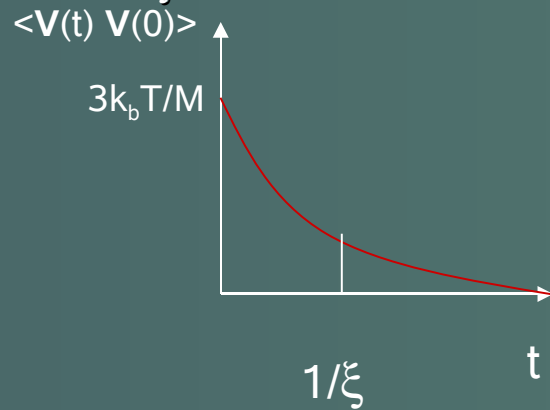
$$\langle \Theta_\alpha(t) \rangle = 0 \quad \langle \Theta_\alpha(t) \Theta_\beta(t') \rangle = 2\Theta_0 \delta_{\alpha\beta} \delta(t-t')$$

uncorrelated character of collisions

# Meso-scale physics

Langevin description of Brownian motion (1908)

- Langevin equation as an evolution equation for PDF of velocity
- $1/\xi$  – relaxation time scale



Evolution from ballistic to the diffusive motion

$$\text{Diffusion } D = k_b T / M \xi$$

Fokker-Plank equation; Kramers Smoluchowski: SDE  $\rightarrow$  PDE for PDF

$$\frac{\partial P(\mathbf{V}, t)}{\partial t} = \frac{\partial}{\partial \mathbf{V}} \left( \xi \mathbf{V} P(\mathbf{V}, t) + \frac{\partial}{\partial \mathbf{V}} \left( \frac{\Theta_0}{M^2} P(\mathbf{V}, t) \right) \right)$$



# Fluctuation-Dissipation Theorem

General solution of the Langevin equation

$$V(t) = V(0) \exp(-\xi t) + \frac{1}{M} \int_0^t ds \exp(-\xi(t-s)) \Theta(s)$$

$$\lim_{t \rightarrow \infty} \langle v(t)^2 \rangle = \frac{3\Theta_0}{M^2 \xi}$$

from equipartition theorem

$$\lim_{t \rightarrow \infty} \langle v(t)^2 \rangle = \frac{3k_b T}{M}$$

Fluctuation-Dissipation theorem ensures for equilibrium that PDF is equal to Maxwell distribution

$$M \xi k_b T = \Theta_0$$

# Numerical methods

Continuum fluid mechanics	Meso-scale mechanics	Microscopic mechanics
<ul style="list-style-type: none"><li>■ Global parameters: density, velocity, energy, temperature</li></ul>	<ul style="list-style-type: none"><li>■ <u>lattice methods</u> sacrifice detail in potential model, simpler interactions and motion rules; LBM</li><li>■ <u>particle methods</u> particles as mesoscopic object; replace small particle with random forces; DPD, SPH</li></ul>	<ul style="list-style-type: none"><li>■ <u>Molecular Dynamics</u> computation of forces, based on the interaction potential; particle move according to Newton's equation of motion</li><li>■ <u>Monte Carlo methods</u> Set a configuration, make a trial move; acceptance/rejection procedure, and accumulation of averages</li></ul>

# Dissipative Particle Dynamics

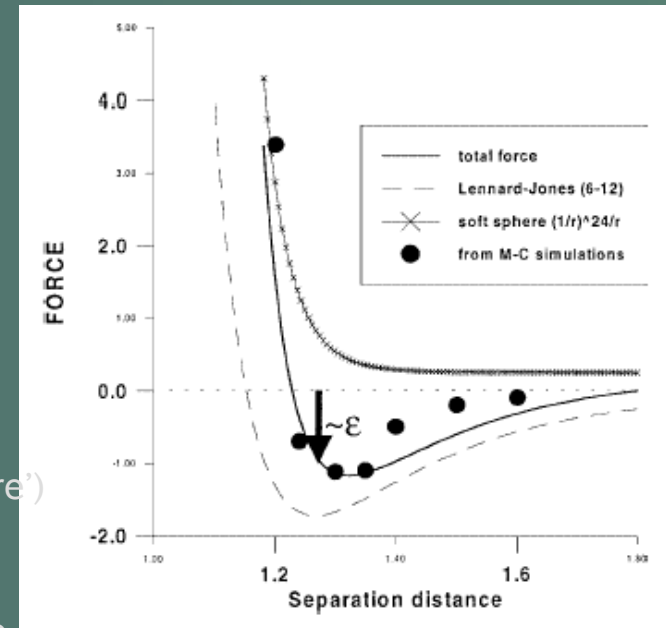
## Spherical formulation

two particle interaction



## Three type of forces

- **conservative** (purely repulsive, represents ‘pressure’)
- **dissipative** (reducing velocity difference between particles, ‘friction forces’)
- **stochastic** (‘degree of freedom’ removed by coarse-graining procedure)

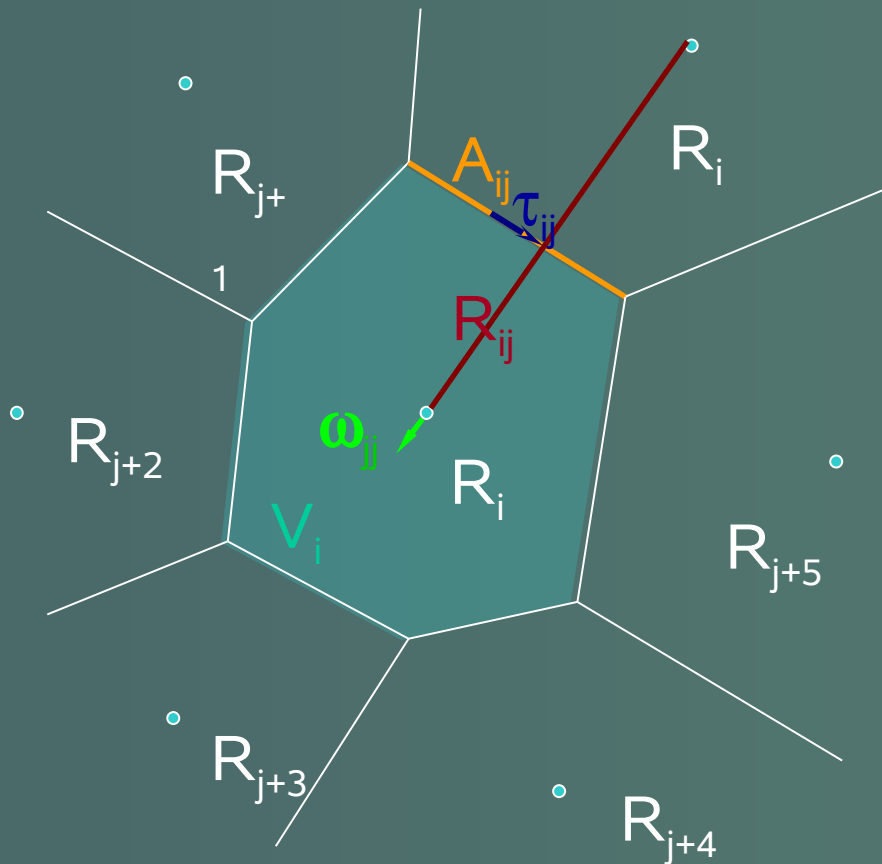


$$F_C = \pi \omega(r_{ij}) e_{ij};$$

$$F_D = \gamma M \omega(r_{ij}) (e_{ij} \circ v_{ij}) e_{ij};$$

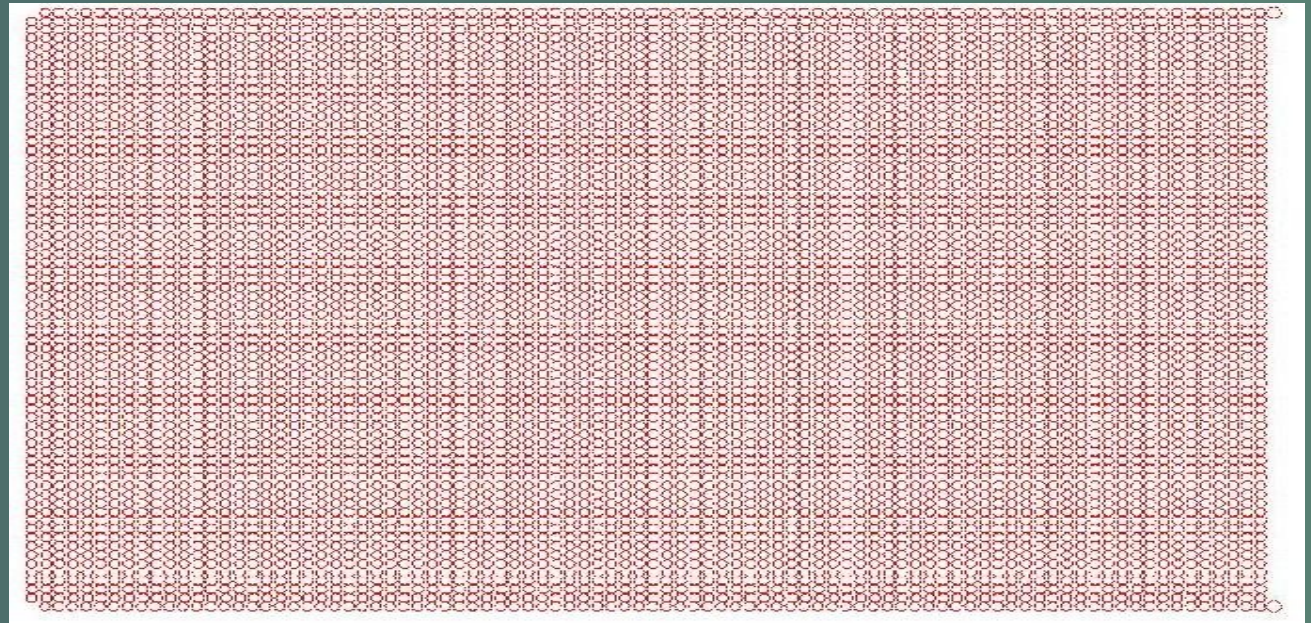
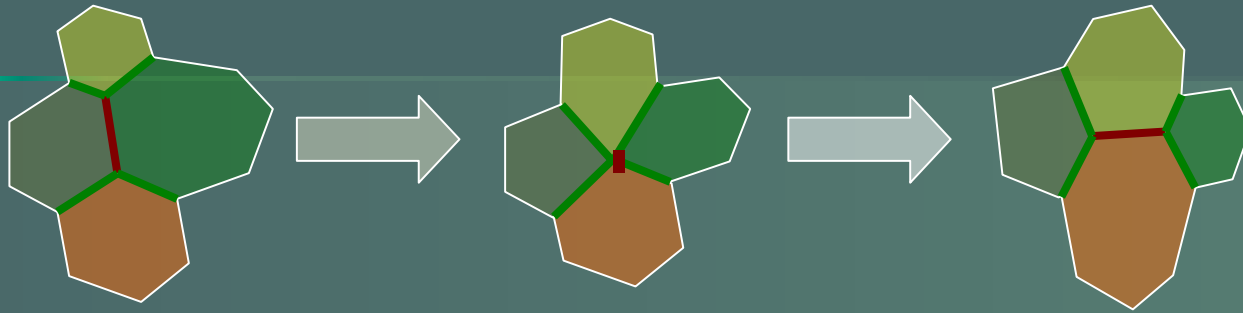
$$F_B = \frac{\delta \theta_{ij}}{\sqrt{\Delta t}} \omega(r_{ij}) e_{ij};$$

# Dissipative Particle Dynamics



- $R_{ij}$  – distance between centers of Voronoi  $i$  and  $j$ ;
- $\omega_{ij}$  – unit vector normal to the face  $ij$ , originated from the Voronoi center  $i$ ;
- $A_{ij}$  – area of the contact surface: length in 2D;
- $r_{ij}$  – vector indicating the mass center of the contact surface  $ij$  originated from the center of the surface  $ij$ : edge in 2D ;
- $V_i$  – is Voronoi volume: surface in 2D.

# Topological Properties



# Dissipative Particle Dynamics

## Classical DPD

$$\frac{d\mathbf{r}_i}{dt} = \mathbf{v}_i(t) \quad \frac{d\mathbf{p}_i}{dt} = f_i(t) \equiv \sum_{j \neq i} \left[ \mathbf{F}_{ij}^C(\mathbf{r}_{ij}) + \mathbf{F}_{ij}^D(\mathbf{r}_{ij}, \mathbf{v}_{ij}) + \mathbf{F}_{ij}^R(\mathbf{r}_{ij}) \right]$$

## DPDE – energy conservation

$$\frac{de_{ij}}{dt} = \sum_{j \neq i} \left[ \frac{dq_{ij}}{dt} + \frac{dq_{ij}^D}{dt} + \frac{dq_{ij}^R}{dt} \right]$$

## Voronoi DPD

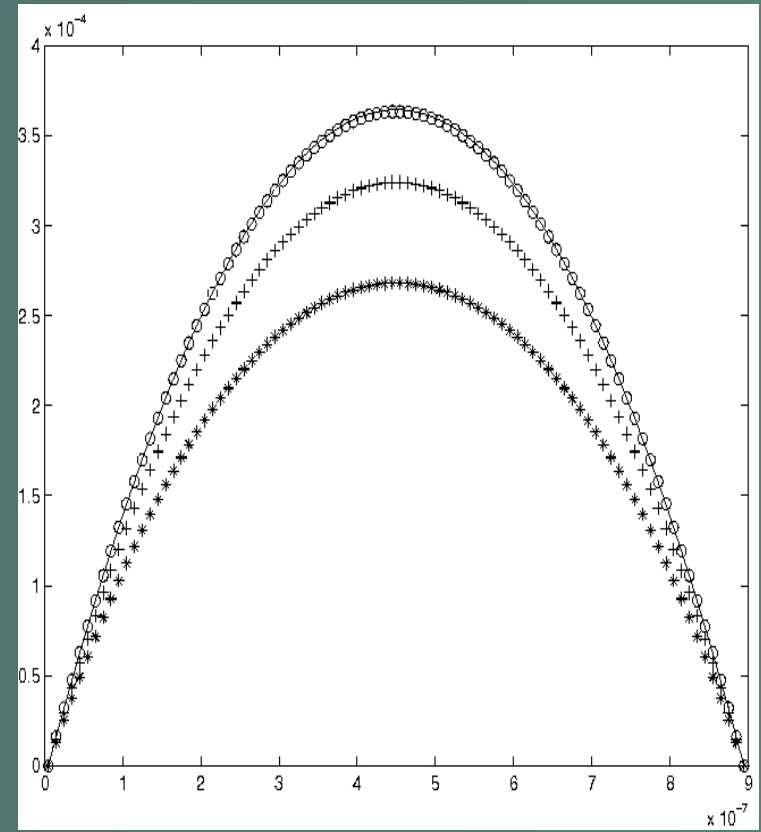
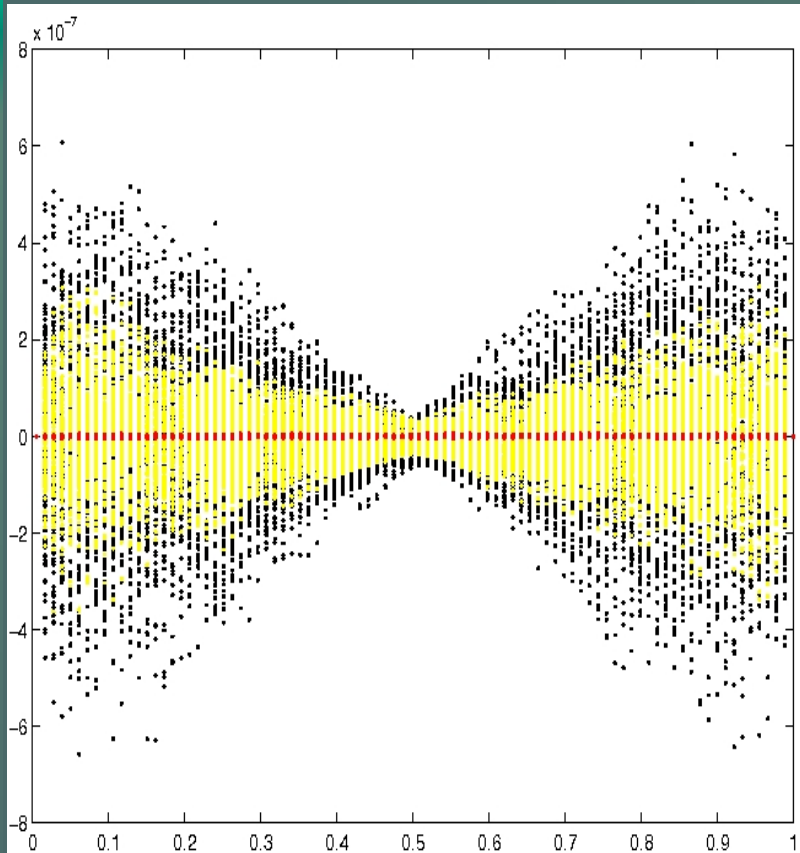
reversible part

$$\begin{aligned} \dot{\mathbf{R}}_i &= \mathbf{v}_i \\ \dot{M}_i &= \sum_j \frac{A_{ij}}{R_{ij}} \frac{\rho_i + \rho_j}{2} \boldsymbol{\tau}_{ij} \cdot (\mathbf{v}_i - \mathbf{v}_j) \\ \dot{\mathbf{P}}_i &= \sum_j A_{ij} \boldsymbol{\omega}_{ij} (p_i - p_j) / 2 + \sum_j \frac{A_{ij}}{\hat{R}_{ij}} \frac{\rho_i + \rho_j}{2} \frac{\mathbf{v}_i + \mathbf{v}_j}{2} \times \boldsymbol{\tau}_{ij} \cdot (\mathbf{v}_i - \mathbf{v}_j) + \\ &+ \sum_j \frac{A_{ij}}{R_{ij}} \boldsymbol{\tau}_{ij} \left[ (p_i - p_j) - \frac{\rho_i + \rho_j}{2} (\mu_i - \mu_j) + -\frac{s_i + s_j}{2} (T_i - T_j) \right] \\ \dot{S}_i &= \sum_j \frac{A_{ij}}{R_{ij}} \frac{s_i + s_j}{2} \boldsymbol{\tau}_{ij} \cdot (\mathbf{v}_i - \mathbf{v}_j) \end{aligned}$$

irreversible part

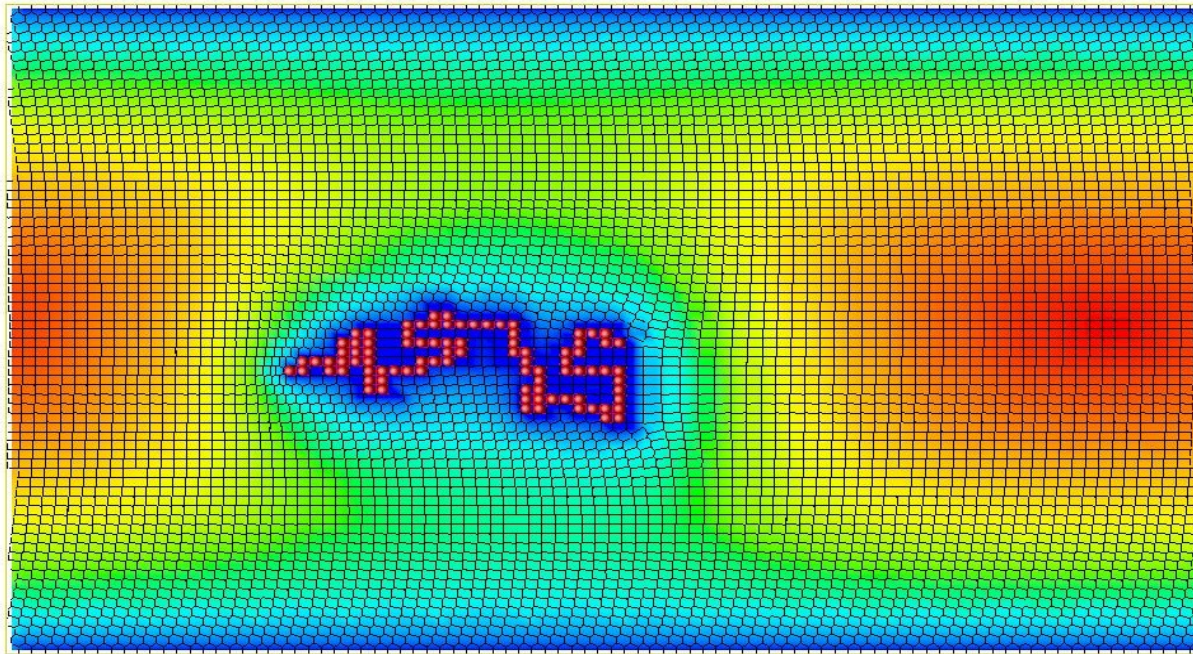
$$\begin{aligned} d\mathbf{P}_i^{irr} &= \sum_j \boldsymbol{\sigma}_{ij} (\boldsymbol{\Pi}_j + \boldsymbol{\Pi}_i \mathbf{1}) dt + d\mathbf{P}_i \\ T_i dS_i^{irr} &= \left( 1 - \frac{k_B}{\hat{C}_i} \right) \left[ \frac{2\eta_i}{\hat{V}_i} \bar{\boldsymbol{\Omega}}_i \otimes \bar{\boldsymbol{\Omega}}_i + \frac{\xi_i}{\hat{V}_i} \Psi_i^2 \right] d\hat{t} + \sum_j \boldsymbol{\sigma}_{ij} \cdot \boldsymbol{\Phi}_j dt \\ &- \frac{k_B}{T_i C_i} \sum_j \boldsymbol{\sigma}_{ij}^2 \frac{\kappa_j}{V_j} T_j^2 dt - T_i \left( \frac{2\eta_i}{\hat{V}_i} + \frac{\xi_i}{\hat{V}_i} \right) \sum_j \frac{\boldsymbol{\sigma}_{ij}}{M_j} dt + T_i d\hat{S}_i \\ d\hat{S}_i &= \frac{1}{T_i} \sum_j \boldsymbol{\sigma}_{ij} \cdot d\hat{\boldsymbol{\Phi}}_j - \frac{1}{T_i} d\hat{\boldsymbol{\theta}} \otimes \sum_j \boldsymbol{\sigma}_{ji} \mathbf{v}_j^T \end{aligned}$$

# Channel Flow



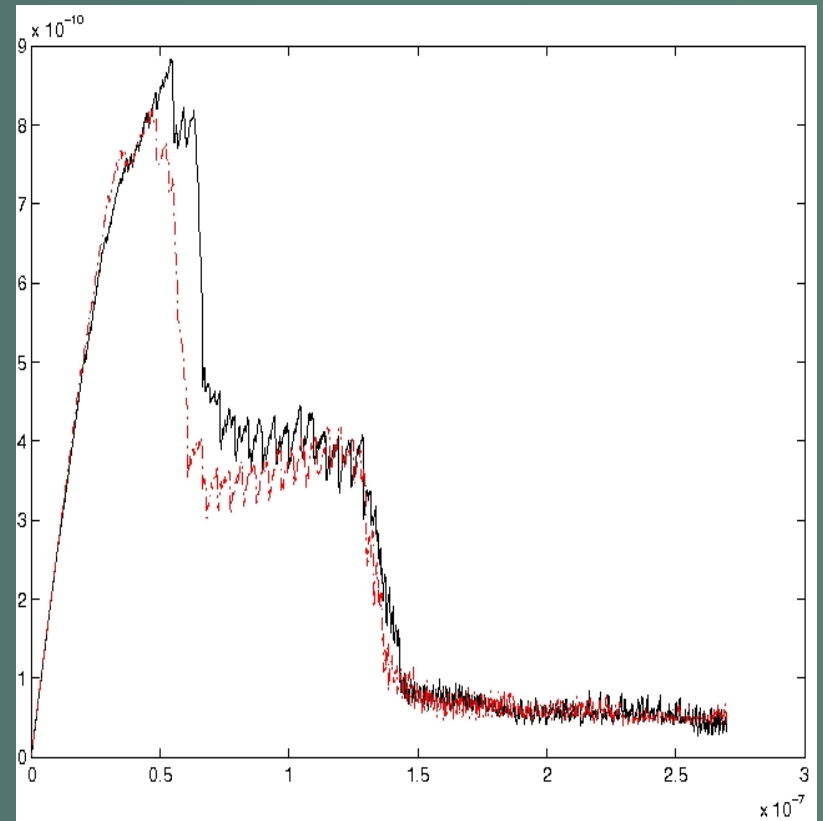
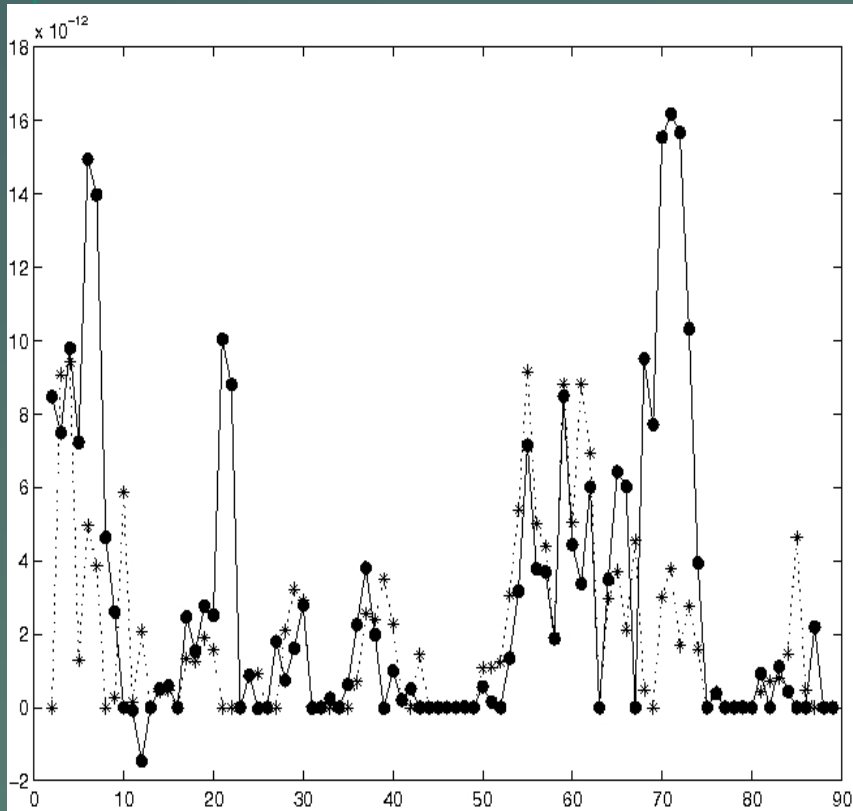


# Channel Flow

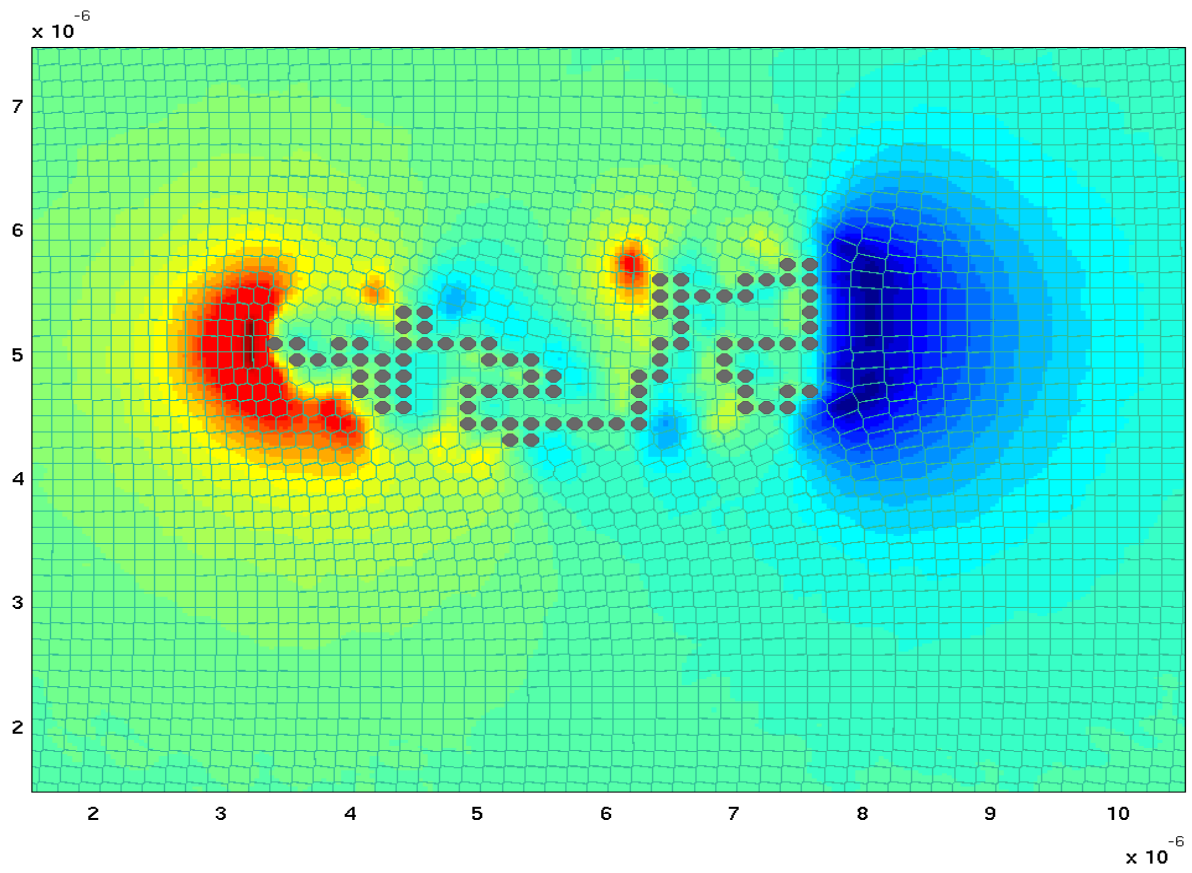




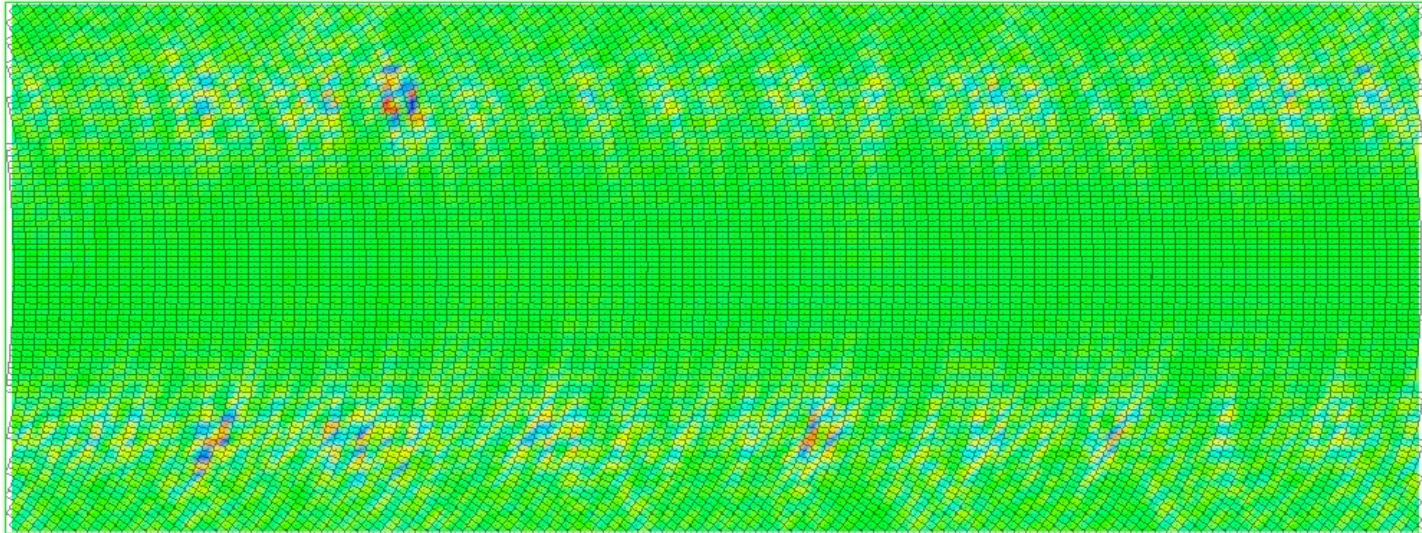
# DNA in the Flow



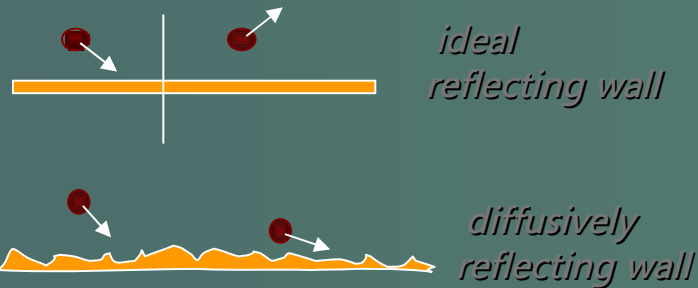
# DNA in the Flow



# Numerical Issues



# Fluctuation-Relaxation Model



Slip velocity  
Maxwell

$$v_{gas} - v_{wall} = \frac{2 - \sigma_v}{\sigma_v} \lambda \left. \frac{\partial v}{\partial y} \right|_{wall}$$

Slip flow and temperature jump  
Smoluchowski

$$v_{gas} - v_{wall} = \frac{2 - \sigma_v}{\sigma_v} \lambda \left. \frac{\partial v}{\partial y} \right|_{wall} + \frac{3}{4} \frac{\mu}{\rho T_{gas}} \left. \frac{\partial T}{\partial y} \right|_{wall}$$

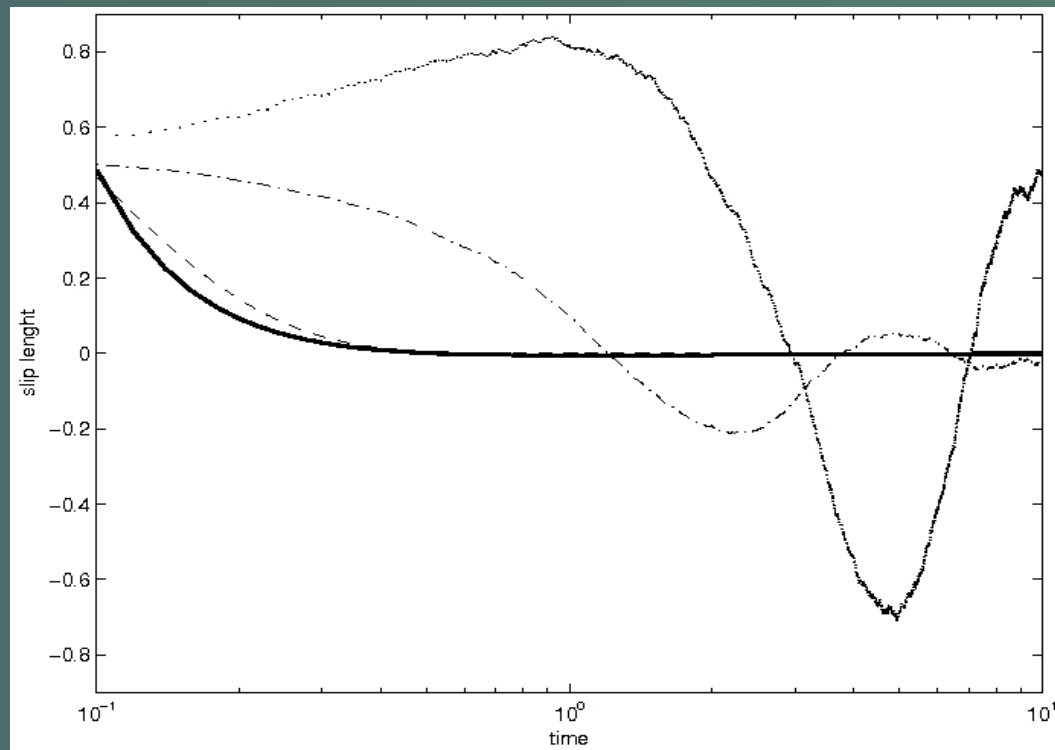
$$T_{gas} - T_{wall} = \frac{2 - \sigma_T}{\sigma_T} \left[ \frac{2\gamma}{\gamma + 1} \right] \frac{\lambda}{Pr} \left. \frac{\partial T}{\partial y} \right|_{wall}$$

# Fluctuation-Relaxation Model

- Relaxation time determines if velocity slip or thermal jump occurs; it defines scale for which meso-scopic (non-equilibrium) and continuum (equilibrium) boundary condition occurs
- Fluctuation Theorem (FDT) compensates molecular level interactions, which has been lost in coarse-grained procedure
- Gas-solid, liquid-solid interface need to be treated similarly; fluid-solid boundary interaction is represented as a non-equilibrium rheological process, the difference between liquid and gas would be mainly in the relaxation time of this phenomena.

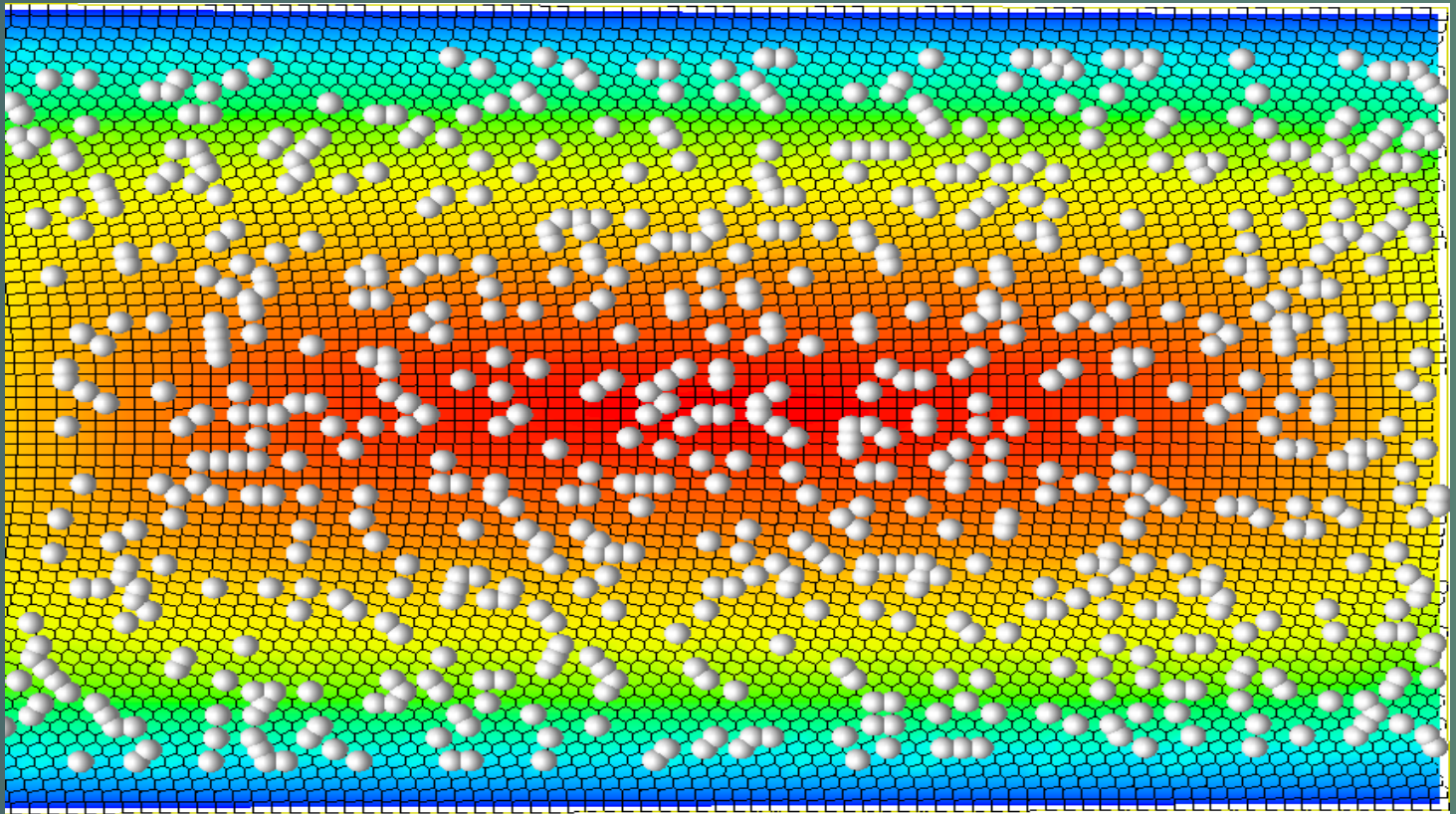
$$\dot{\sigma} = - \int_0^t \eta(t-\tau) \dot{\gamma} d\tau \quad \eta(t-\tau) = \Gamma e^{t/\tau} + \Omega$$

# Slip : No-Slip Phenomena





# Self-organization



# Conclusion

- Mesoscopic description of fluid;
- Time scale and relaxation phenomena – slip, self-organization;
- Mesoscopic solid-fluid interaction – the most important scale for nano- and micro fluidics